

Studying hot and dense nuclear matter in heavy-ion collisions

Agnieszka Sorensen

University of Washington



INSTITUTE for NUCLEAR THEORY

July 19th, 2023

Studying hot and dense nuclear matter in heavy-ion collisions (HICs)

What is “dense”?

Number density in the cores of heavy nuclei: $n_0 \approx 0.160 \text{ fm}^{-3} \Rightarrow \rho_0 = 2.5 \times 10^{17} \frac{\text{kg}}{\text{m}^3}$

(250 trillion times the density of water)

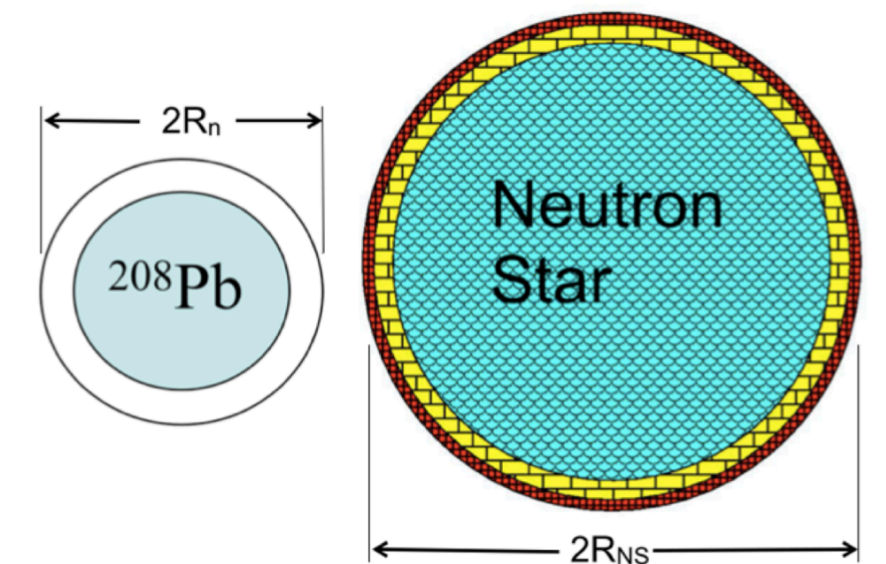
Back-of-the-envelope calculation:

proton radius: $r_p \approx 0.84 \text{ fm}$

volume occupied by a nucleon: $V_N \approx 2.5 \text{ fm}^3 \Rightarrow n_N \equiv \frac{1}{V_N} \approx 0.4 \text{ fm}^{-3}$

percentage of all volume occupied by hard spheres $\approx 75\% \Rightarrow n_{\text{max}} \equiv 0.75 \times n_N = 0.3 \text{ fm}^{-3} \approx 2n_0$

\Rightarrow there's room to squeeze nuclei more! **nucleons are not hard spheres etc.**



How can we squeeze nuclear matter?

- 1) Gravitational force: density in neutron stars likely reaches several times n_0 (**asymmetric matter**)
- 2) Collide heavy nuclei (gold, lead, etc.) at relativistic speeds = relativistic HICs: up to several times n_0

Studying hot and dense nuclear matter in heavy-ion collisions (HICs)

What is “hot”? Always depends on what you’re comparing to.

(surface of the Sun: $\sim 5,800$ K)

“Cold” neutron stars on average have a temperature of about $\approx 2 \times 10^6$ K ≈ 170 eV $\approx 2 \times 10^{-4}$ MeV

Neutron mass: $m_N \approx 938$ MeV Kinetic energy at n_0 : $E \Big|_{n_B=n_0} \approx 994$ MeV \Rightarrow $E_{\text{kin}} \approx 56$ MeV

\Rightarrow the available thermal energy is comparatively small \Rightarrow one can just as well say that $T = 0$

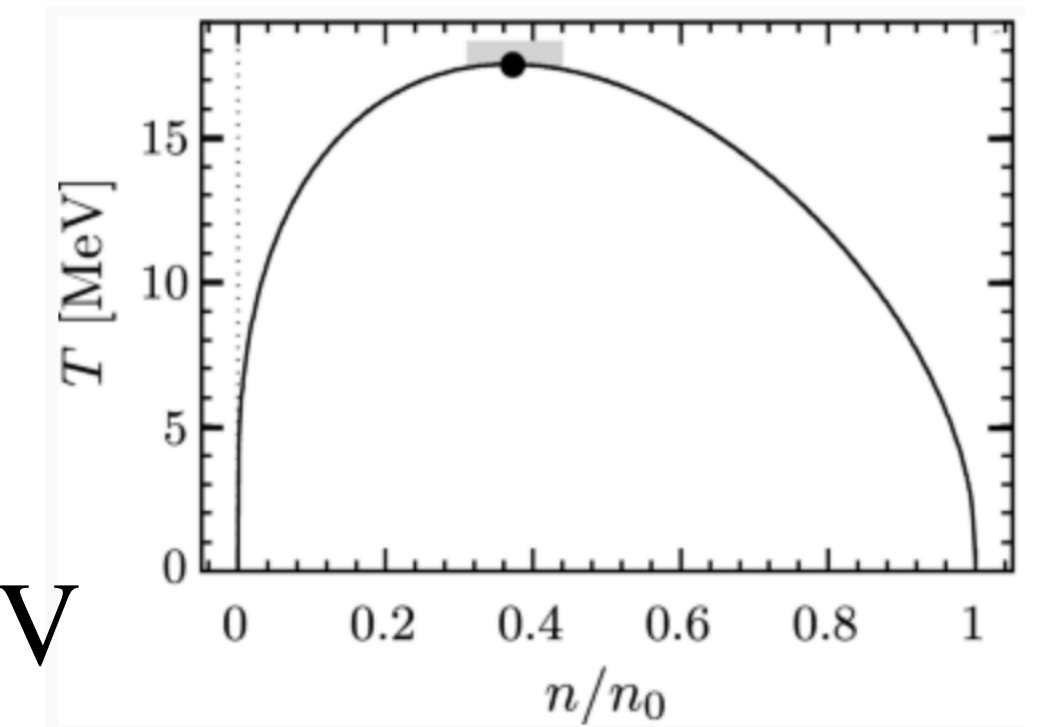
Binding energy of nuclear matter at saturation density: $B_0 \approx -16$ MeV

At $T = 0$, coexistence with vacuum \sim a nuclear liquid drop in empty space

At higher temperatures, some of the nuclear matter coexists with nucleon gas

Critical temperature of the nuclear liquid-gas phase transition: $T_c \in (15, 20)$ MeV

\Rightarrow thermal energy is comparable to relevant energies \Rightarrow for nuclear matter, tens of MeV matter



How can we heat nuclear matter? Squeeze it!

1) Neutron star mergers reach up to ≈ 50 to 100 MeV (**asymmetric matter**)

2) Collide heavy nuclei (gold, lead, etc.) at relativistic speeds = relativistic HICs: up to *hundreds* of MeV!

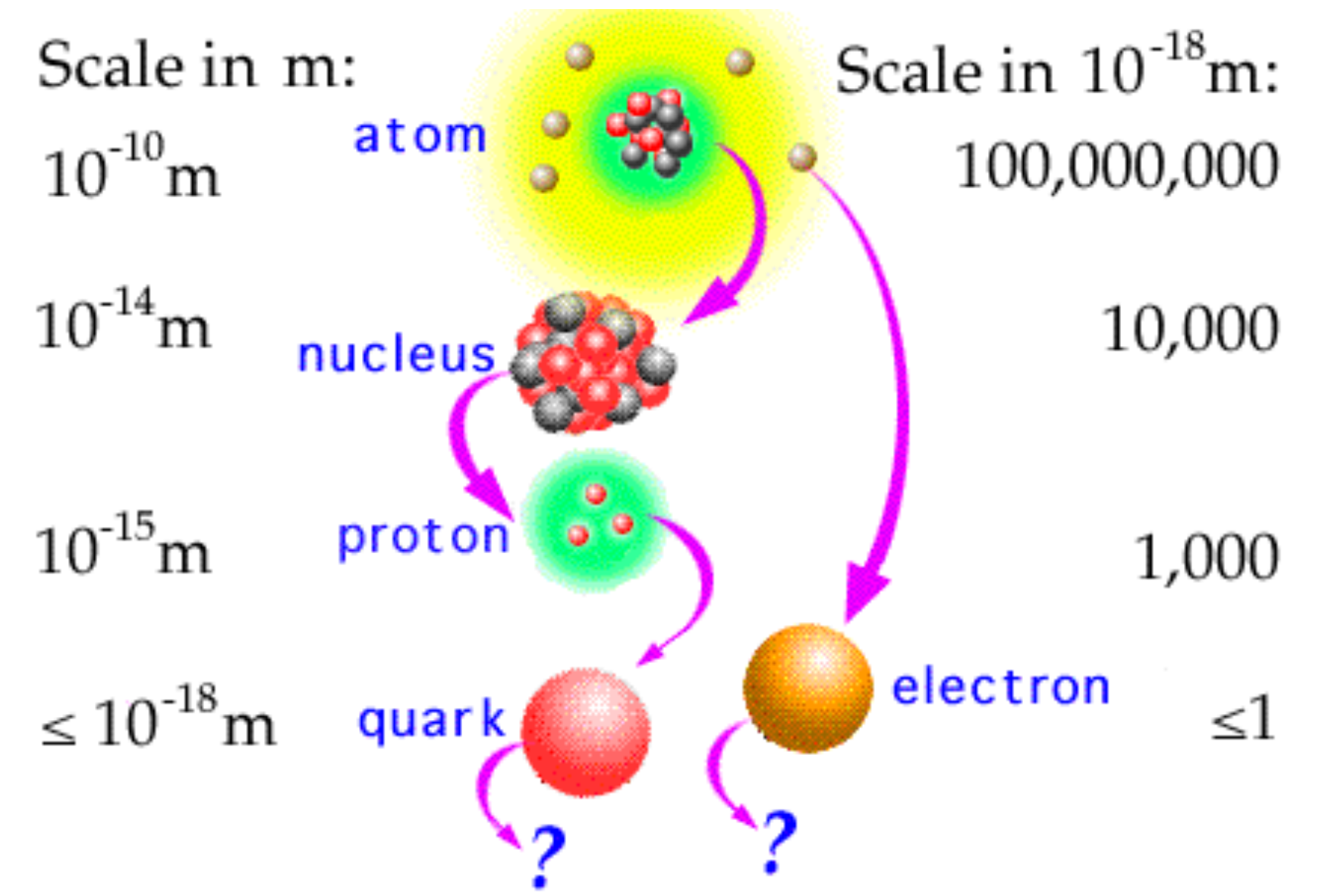
The key mission of RHIC (HICs from the early 2000s): QGP

What happens to nuclei when they are heated up to *hundreds* of MeV?

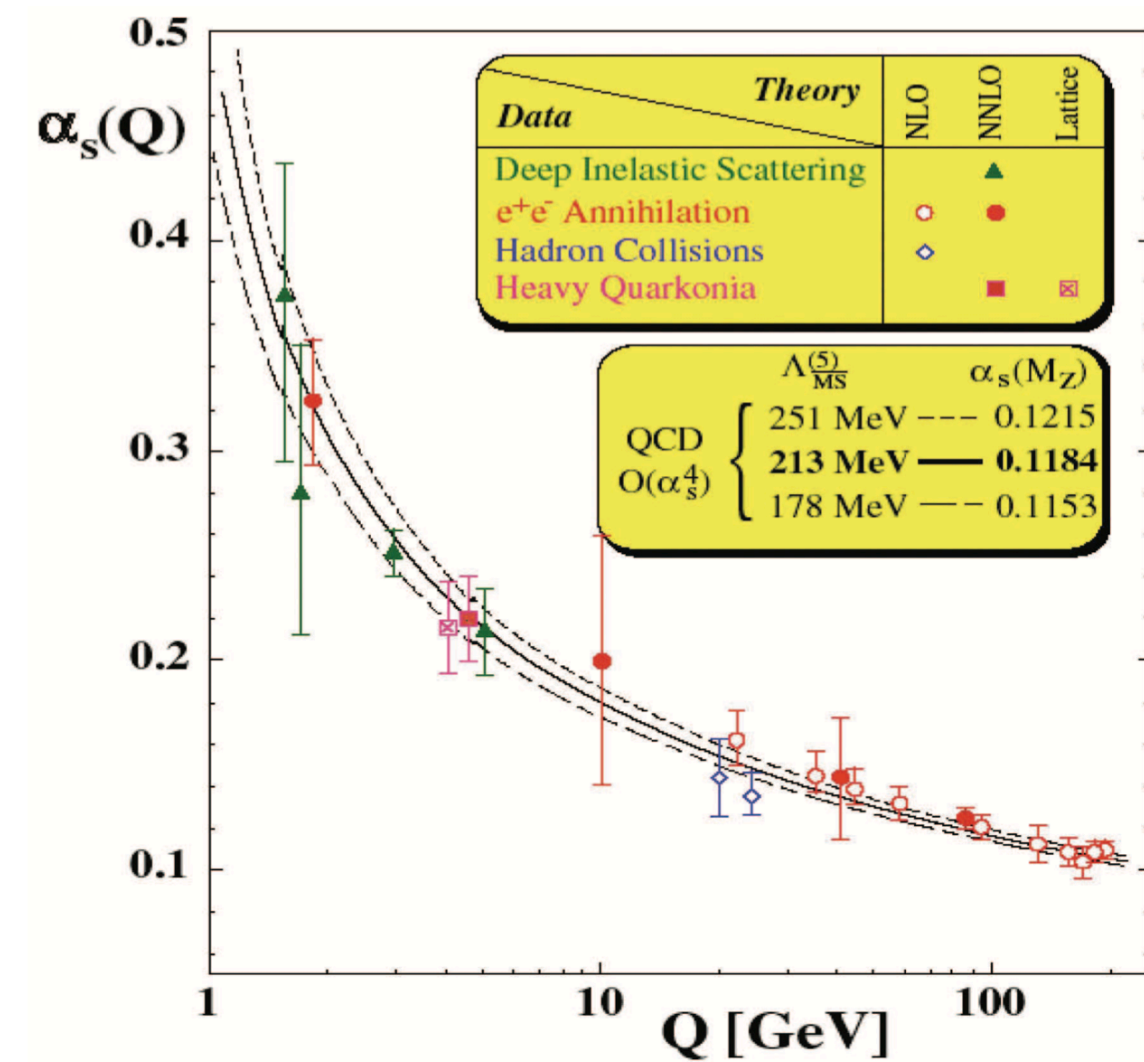
We know that at $T \gtrsim 20$ MeV, nucleons are a homogeneous fluid.

We also know that nucleons are made out of quarks and gluons.

At what temperature = energy scale does it start to matter that nucleons are made out of quarks and gluons?



The QCD coupling constant changes with energy:



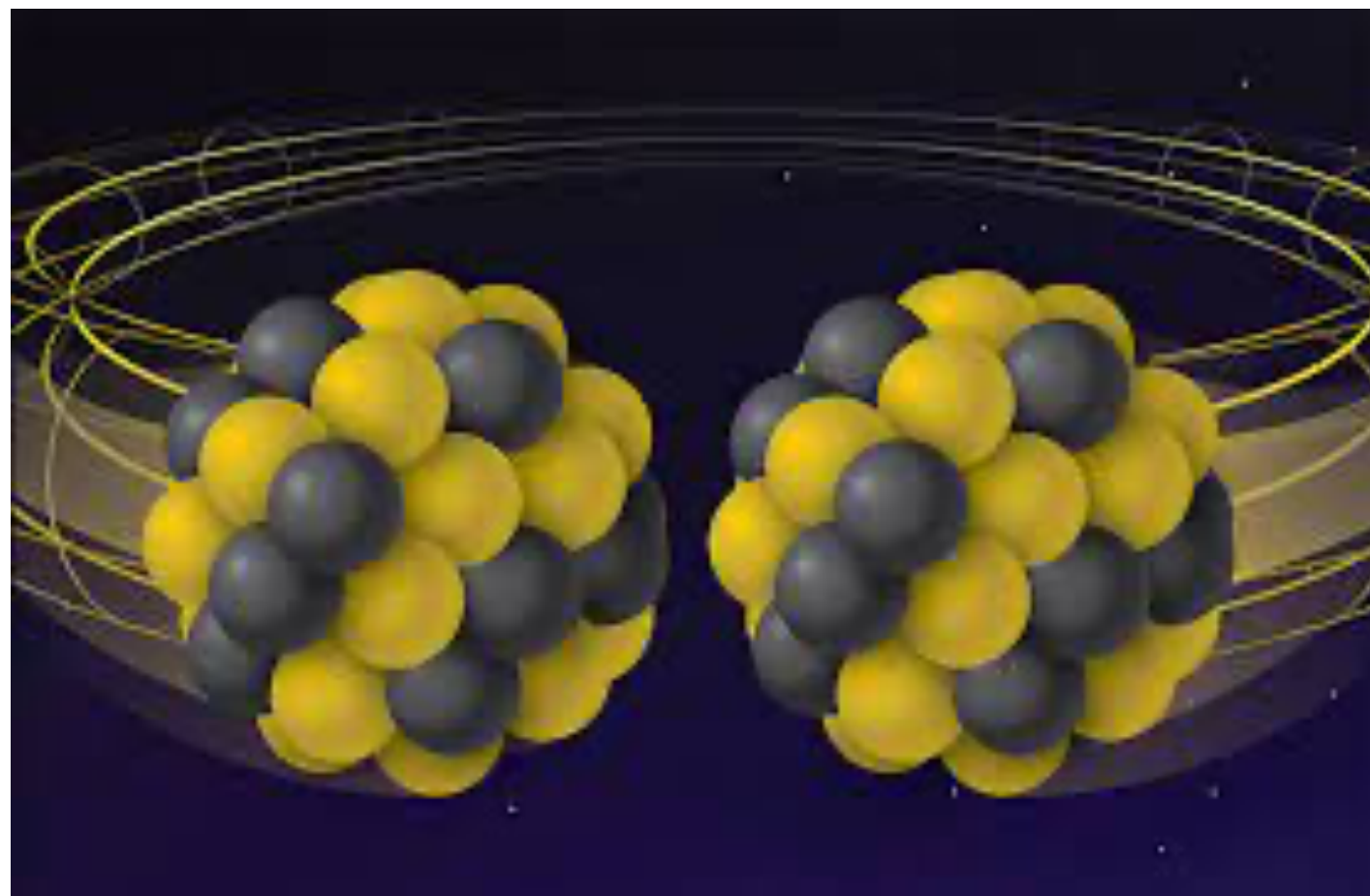
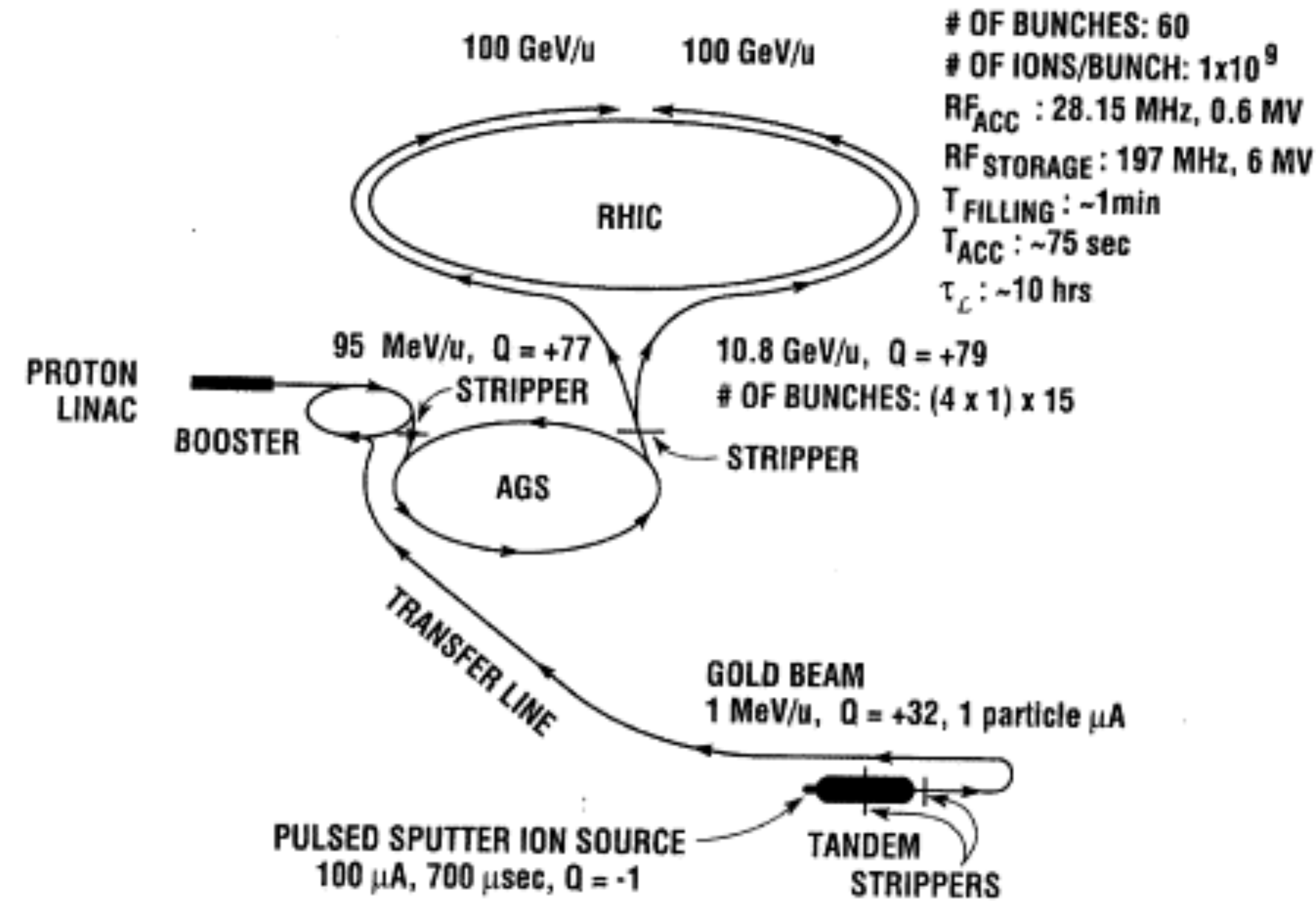
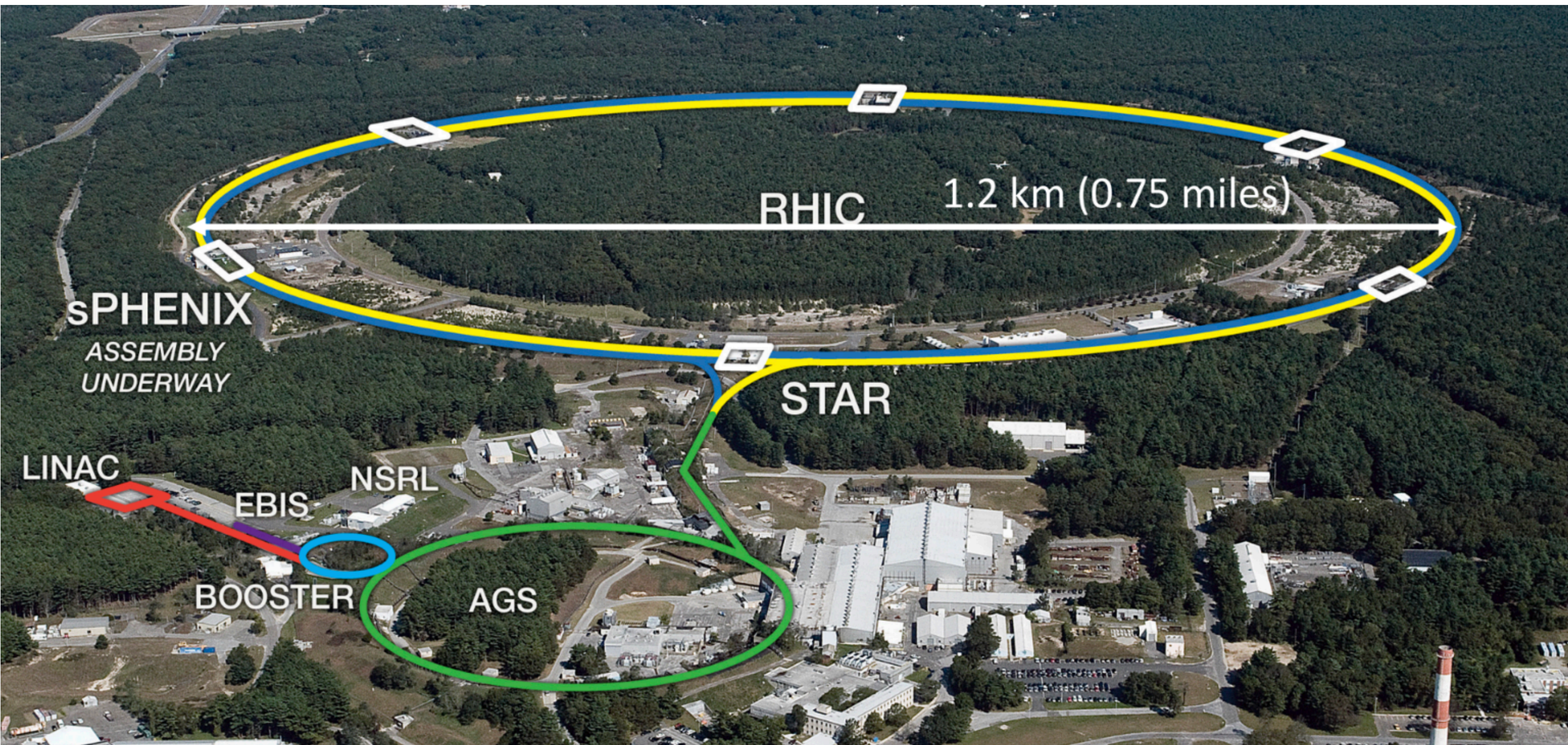
Energy where non-perturbative effects (i.e., confinement) start/stop being important:

$$\Lambda_{\text{QCD}} \approx 200 \text{ MeV}$$

Since HICs lead to high temperatures, can one probe free quarks and gluons in HICs?

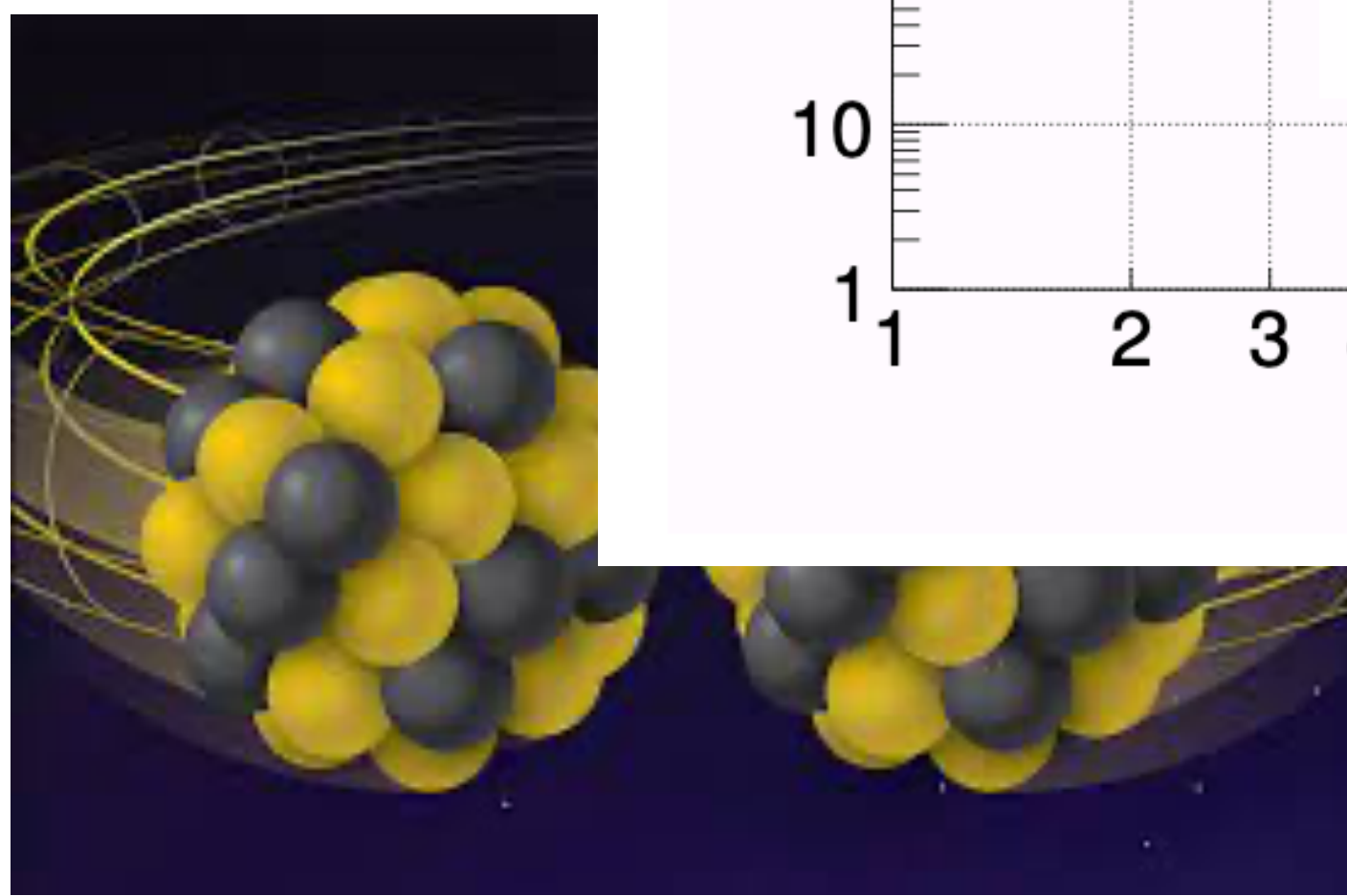
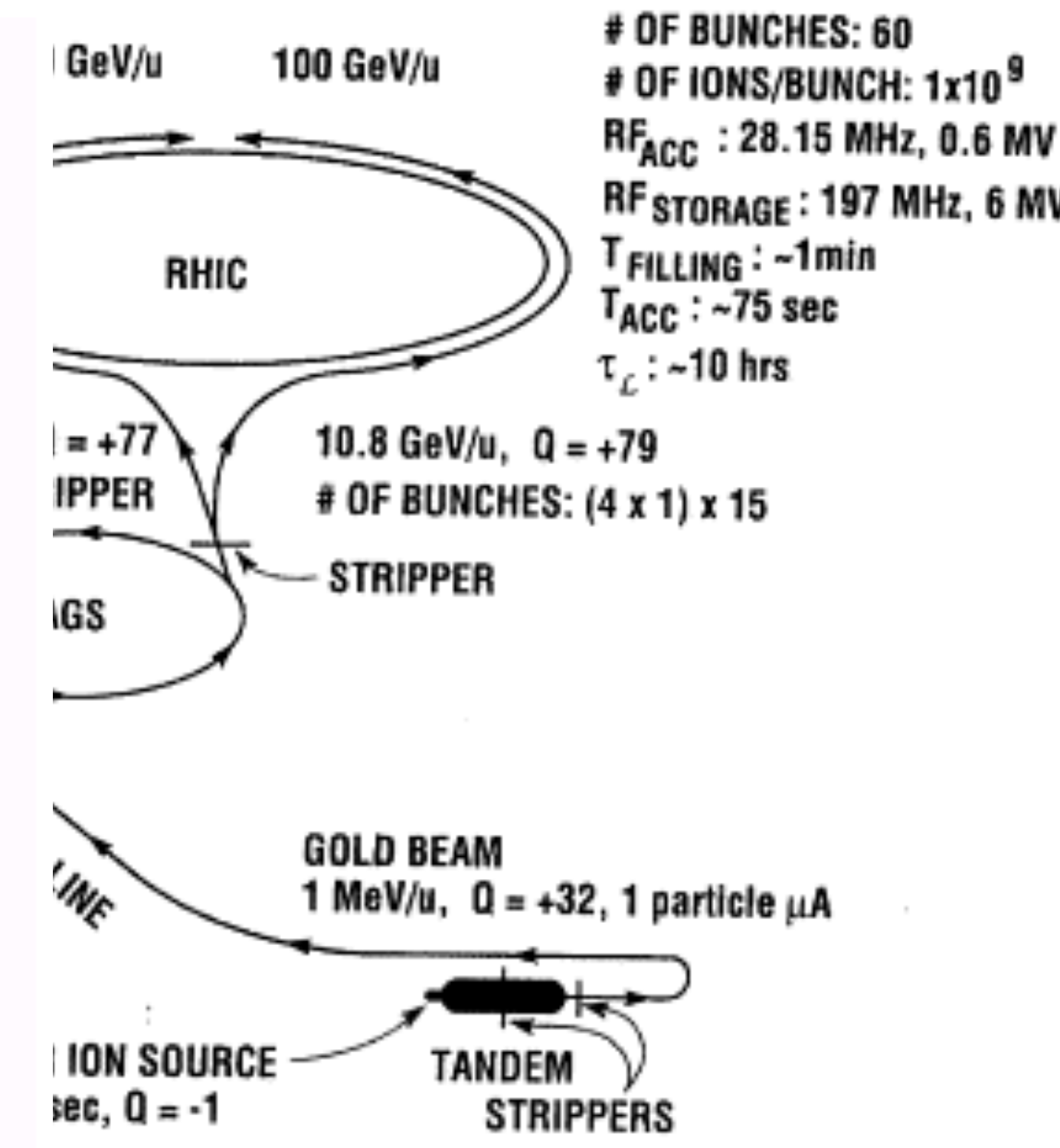
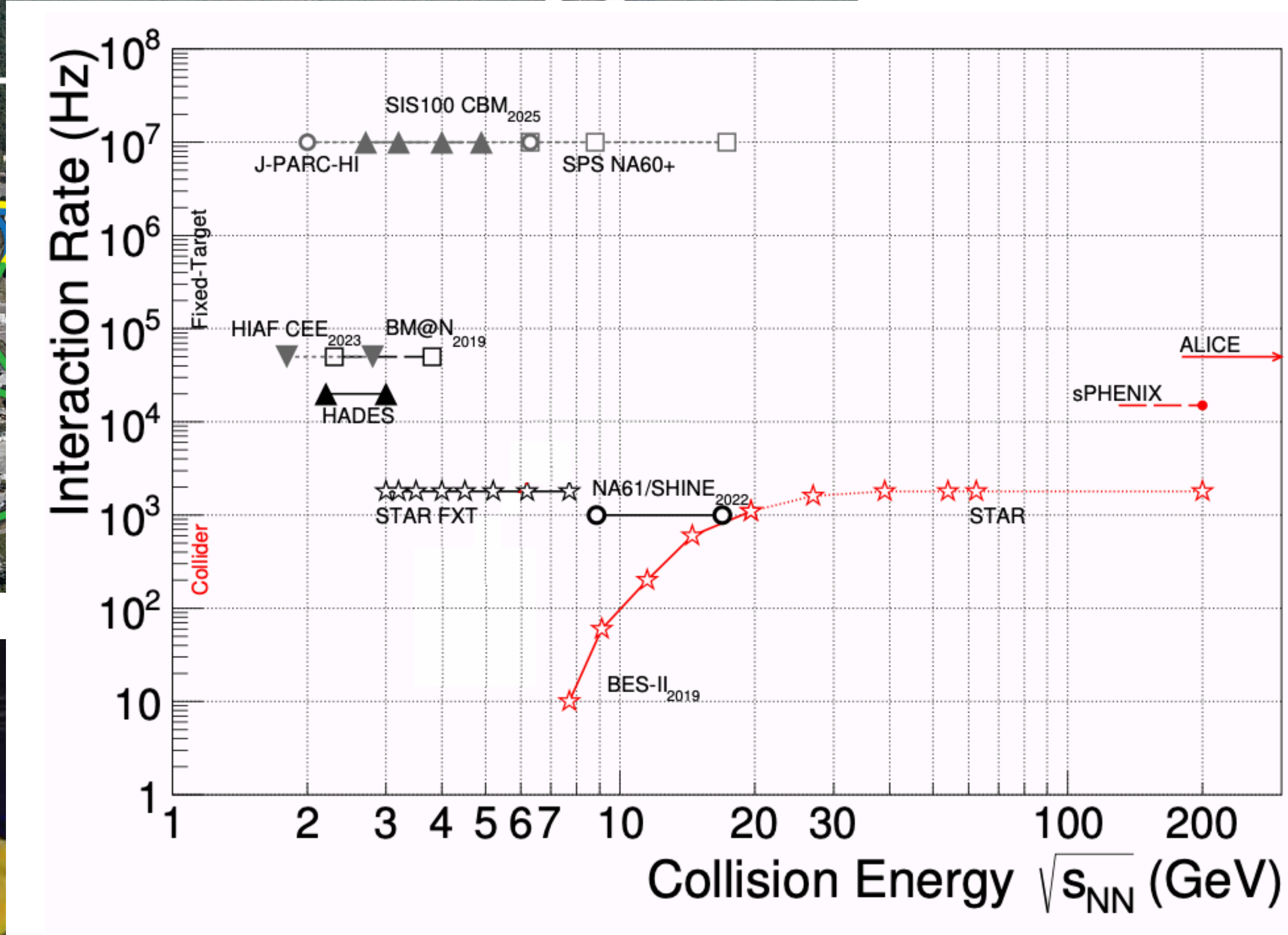
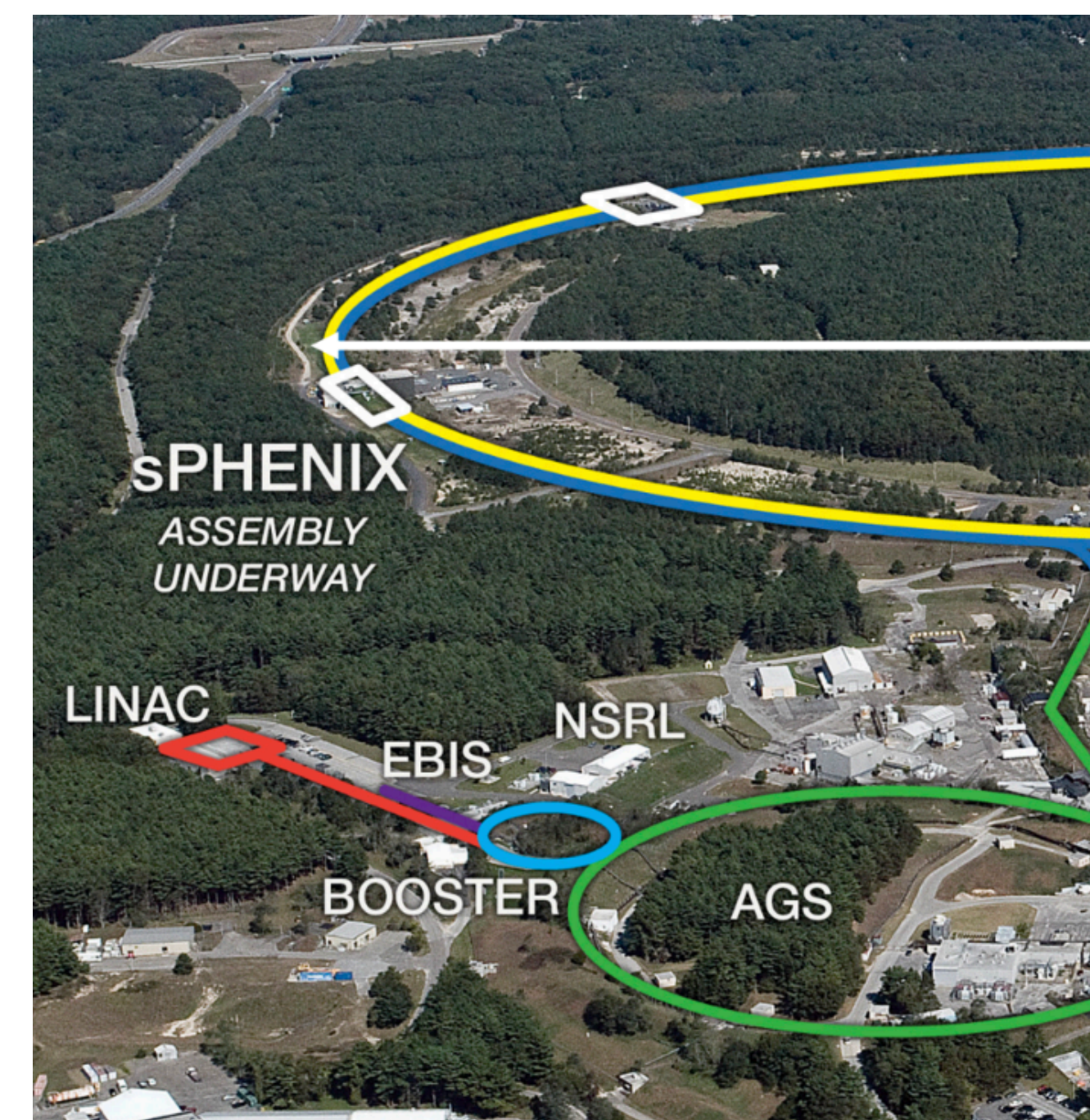
F. Wilczek, "Asymptotic Freedom: From Paradox to Paradigm", https://frankwilczek.com/Wilczek_Easy_Pieces/373_Asymptotic_Freedom.pdf

RHIC = Relativistic Heavy Ion Collider, Brookhaven National Lab, NY



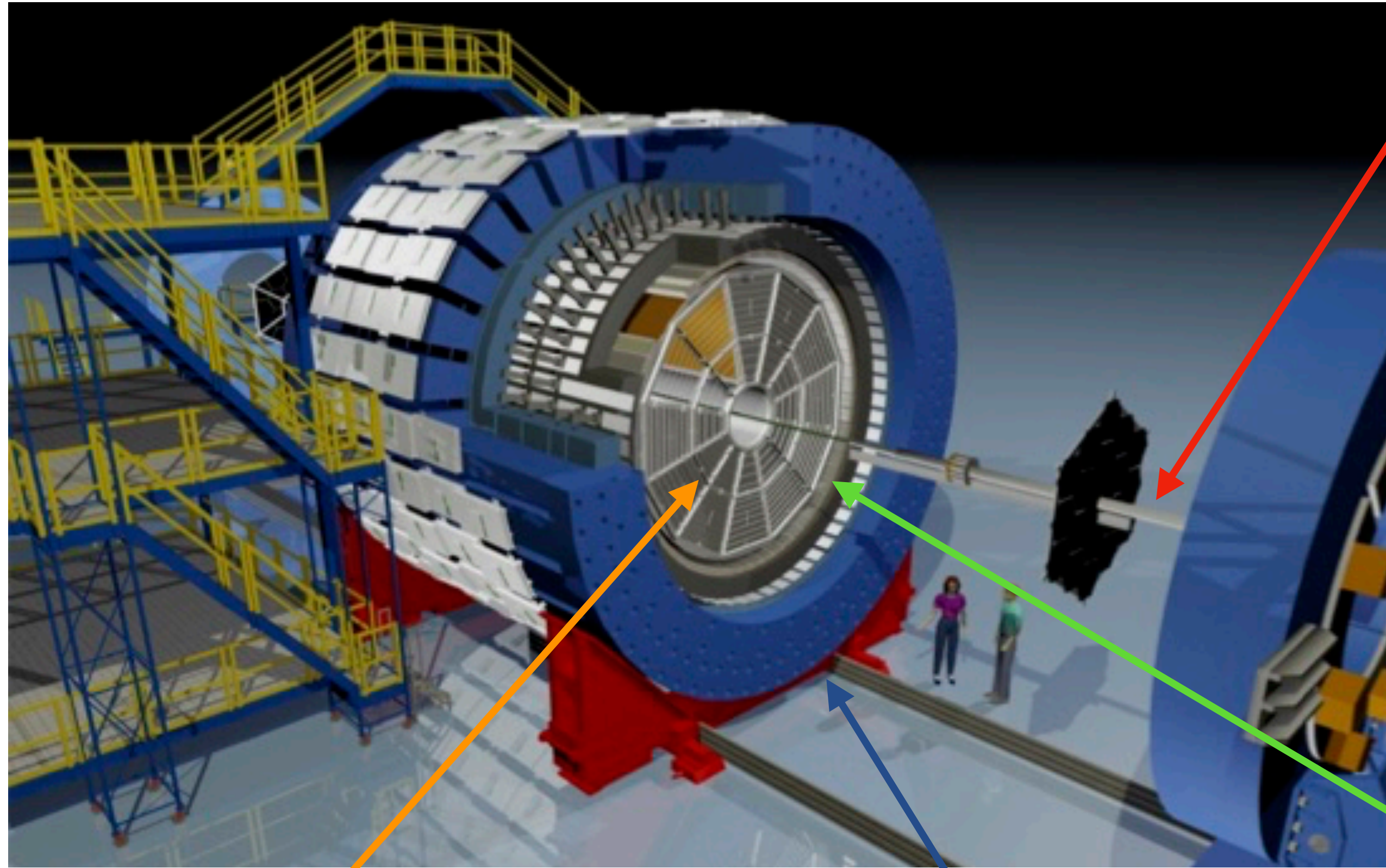
$$\sqrt{s_{NN}} = 200 \text{ GeV} \quad \Rightarrow \quad E_{\text{lab}} \approx 21,000 \text{ GeV}$$

RHIC = Relativistic Heavy Ion Collider, Brookhaven National Lab, NY

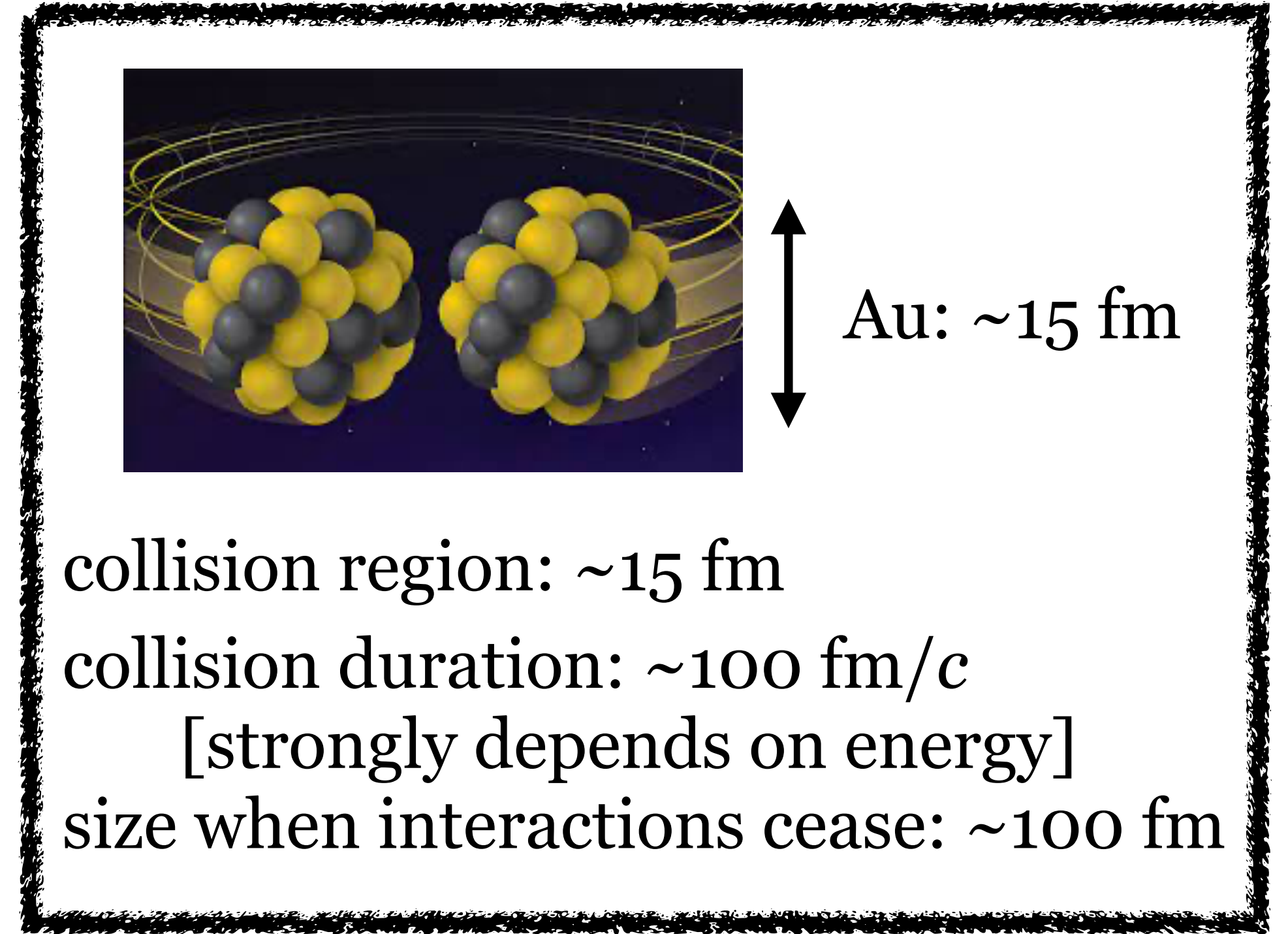


$$\sqrt{s_{NN}} = 200 \text{ GeV} \Rightarrow E_{\text{lab}} \approx 21,000 \text{ GeV}$$

STAR detector



beam pipe: radius ~ 3 cm (diameter of a Coke can)



time projection chamber (TPC):
measures paths of particles
inner radius ~ 0.5 m
outer radius ~ 1.9 m
length ~ 4 m

magnet: curves paths
of charged particles

time of flight (TOF):
measures velocities of particles

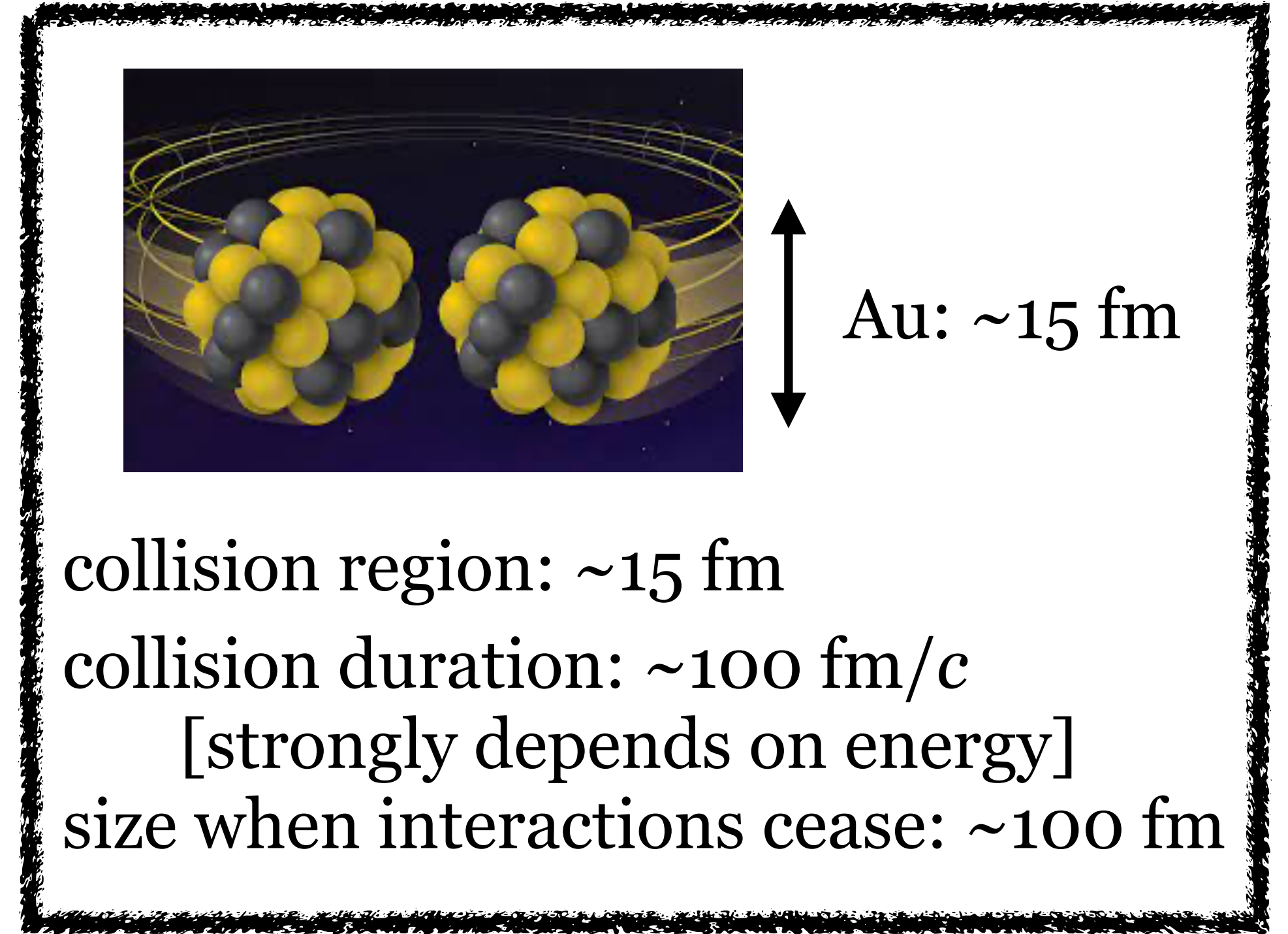
particle momenta

particle mass = particle species identification (ID)

STAR detector

beam pipe: radius ~ 3 cm (diameter of a Coke can)

Detectors only measure *final momenta*;
particle *positions* during collisions unknown



collision region: ~ 15 fm
collision duration: ~ 100 fm/ c
[strongly depends on energy]
size when interactions cease: ~ 100 fm

time projection chamber (TPC):
measures paths of particles
inner radius ~ 0.5 m
outer radius ~ 1.9 m
length ~ 4 m

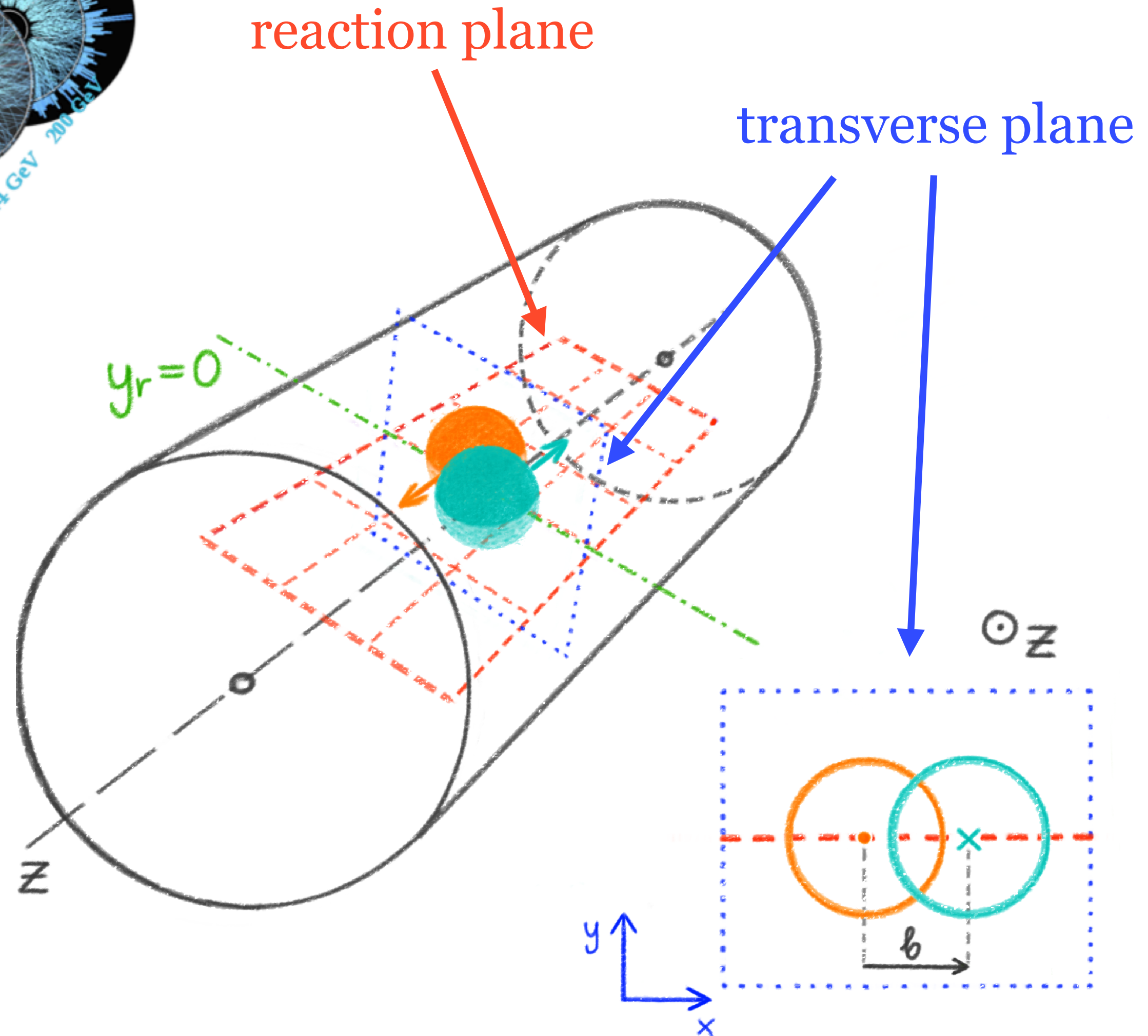
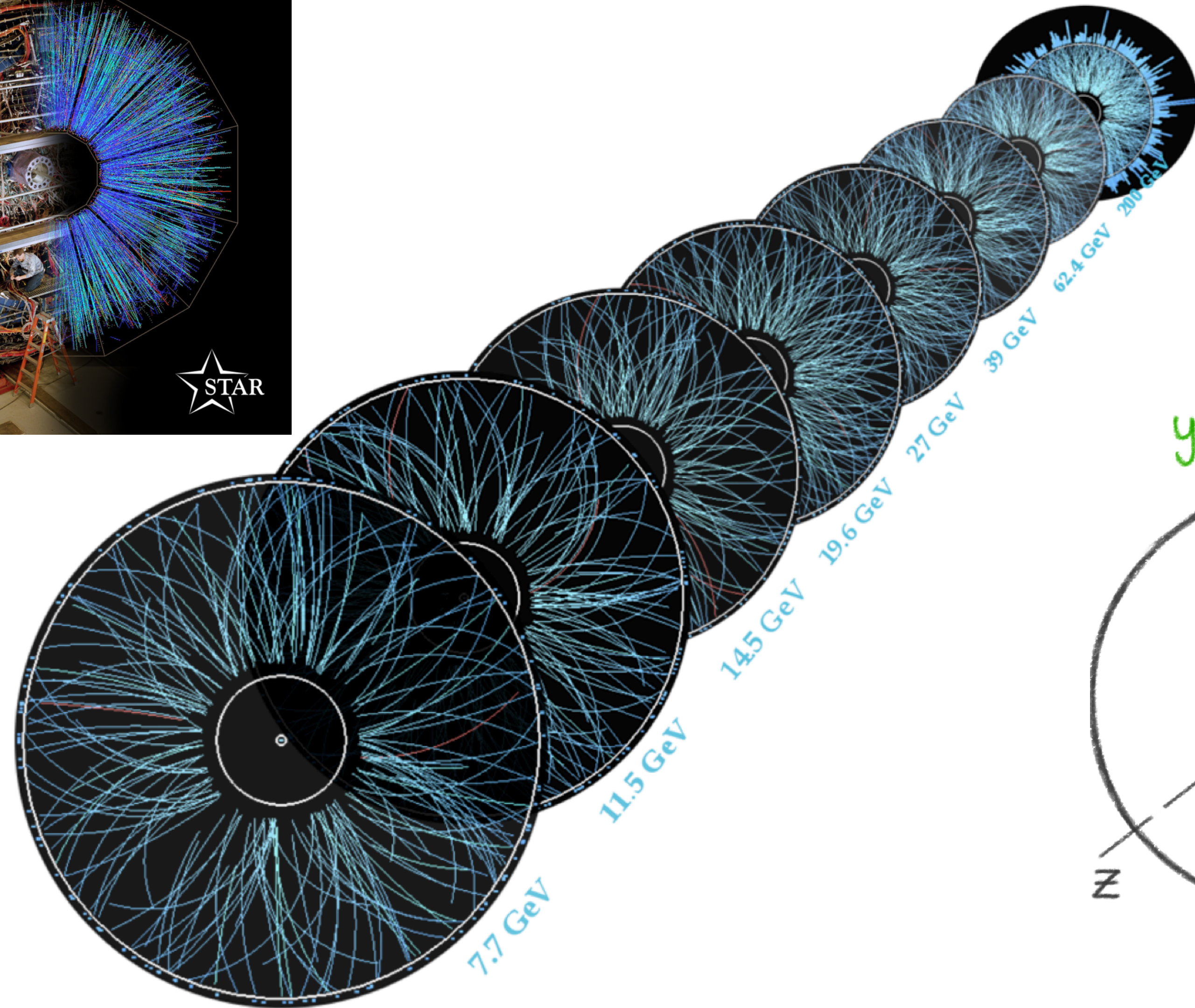
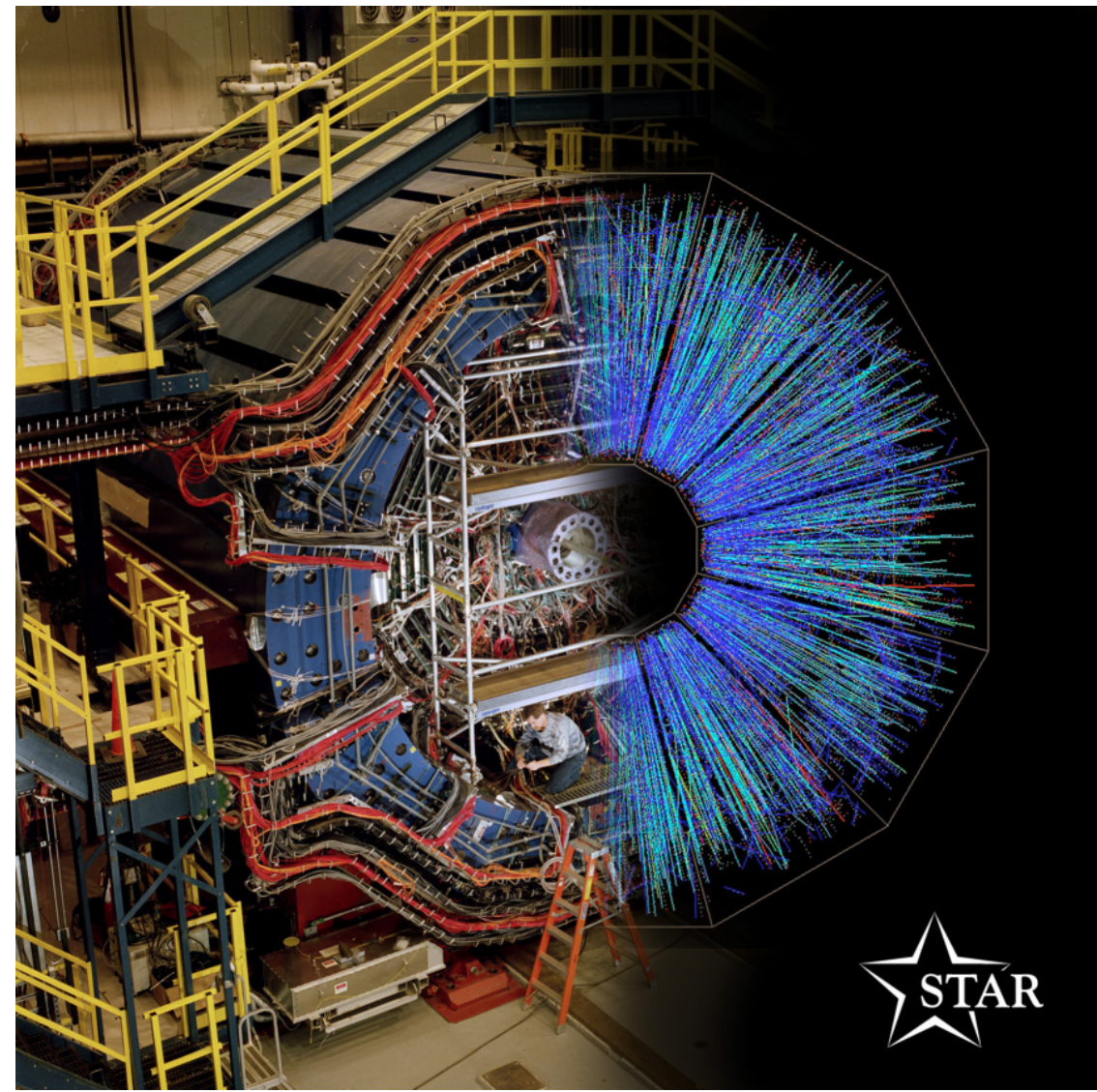
magnet: curves paths
of charged particles

time of flight (TOF):
measures velocities of particles

particle momenta

particle mass = particle species identification (ID)

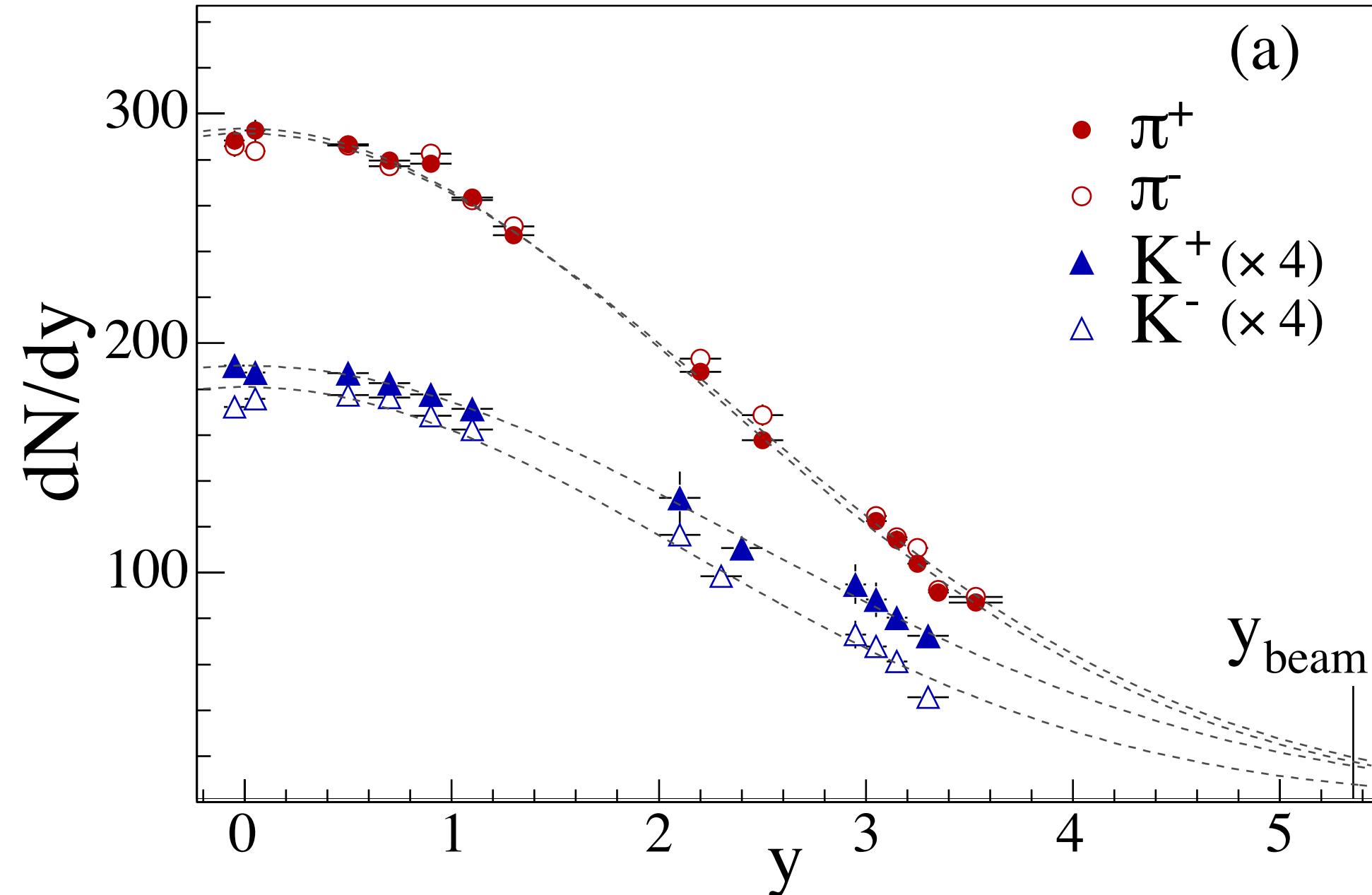
STAR event display



HIC observables

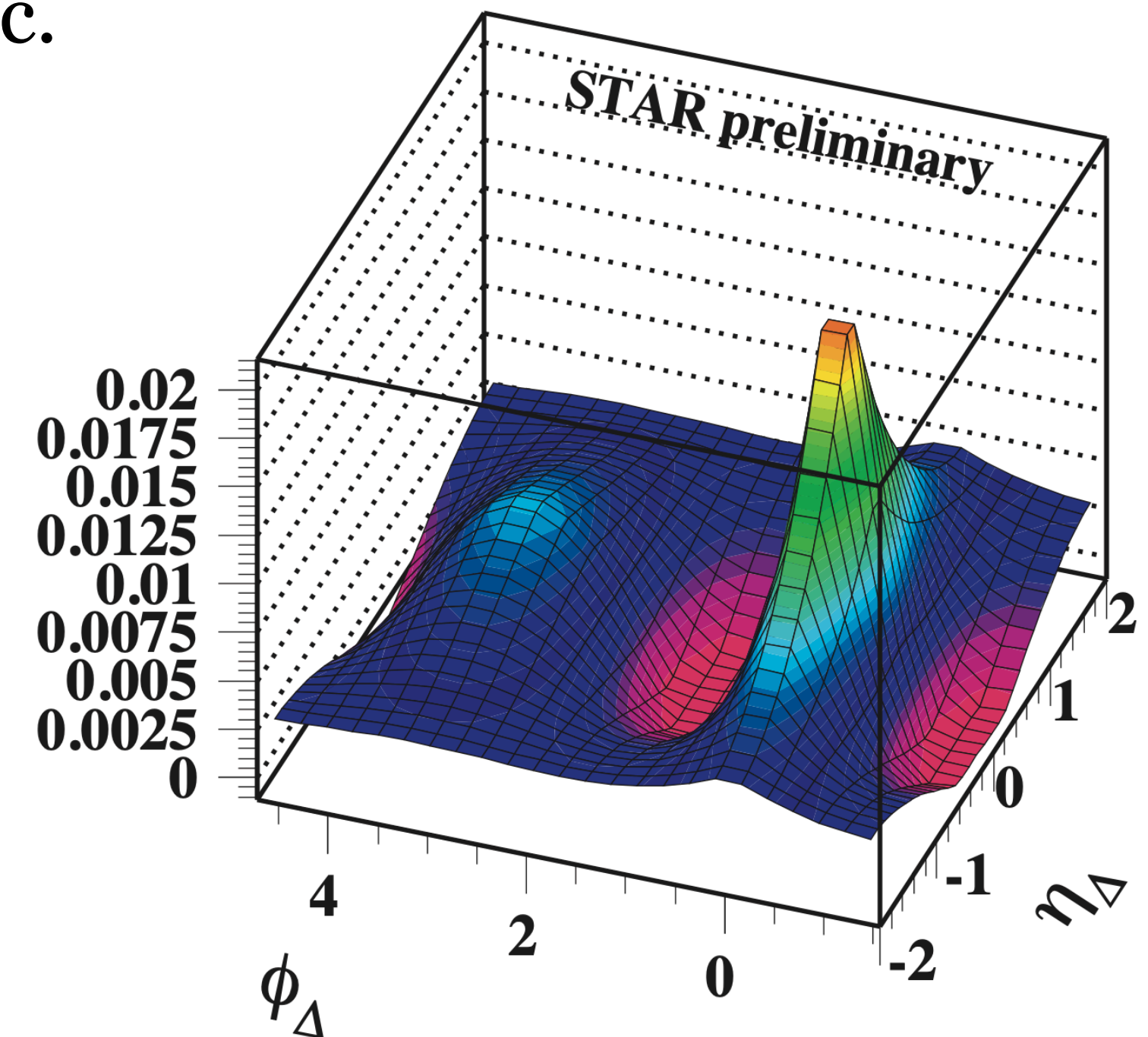
What can we measure? Particle species and momenta:

- single-particle observable: particle yields
 - as a function of rapidity y
 - as a function of transverse momentum p_T
 - etc.



BRAHMS Collaboration, "Charged meson rapidity distributions in central Au+Au collisions at $\sqrt{s(NN)}^{1/2} = 200\text{GeV}$ ", Phys. Rev. Lett. 94, 162301 (2005)
arXiv:nucl-ex/0403050

- many-particle observables:
 - e.g., two-particle correlations,
 - as a function of rapidity y
 - as a function of transverse momentum p_T
 - etc.

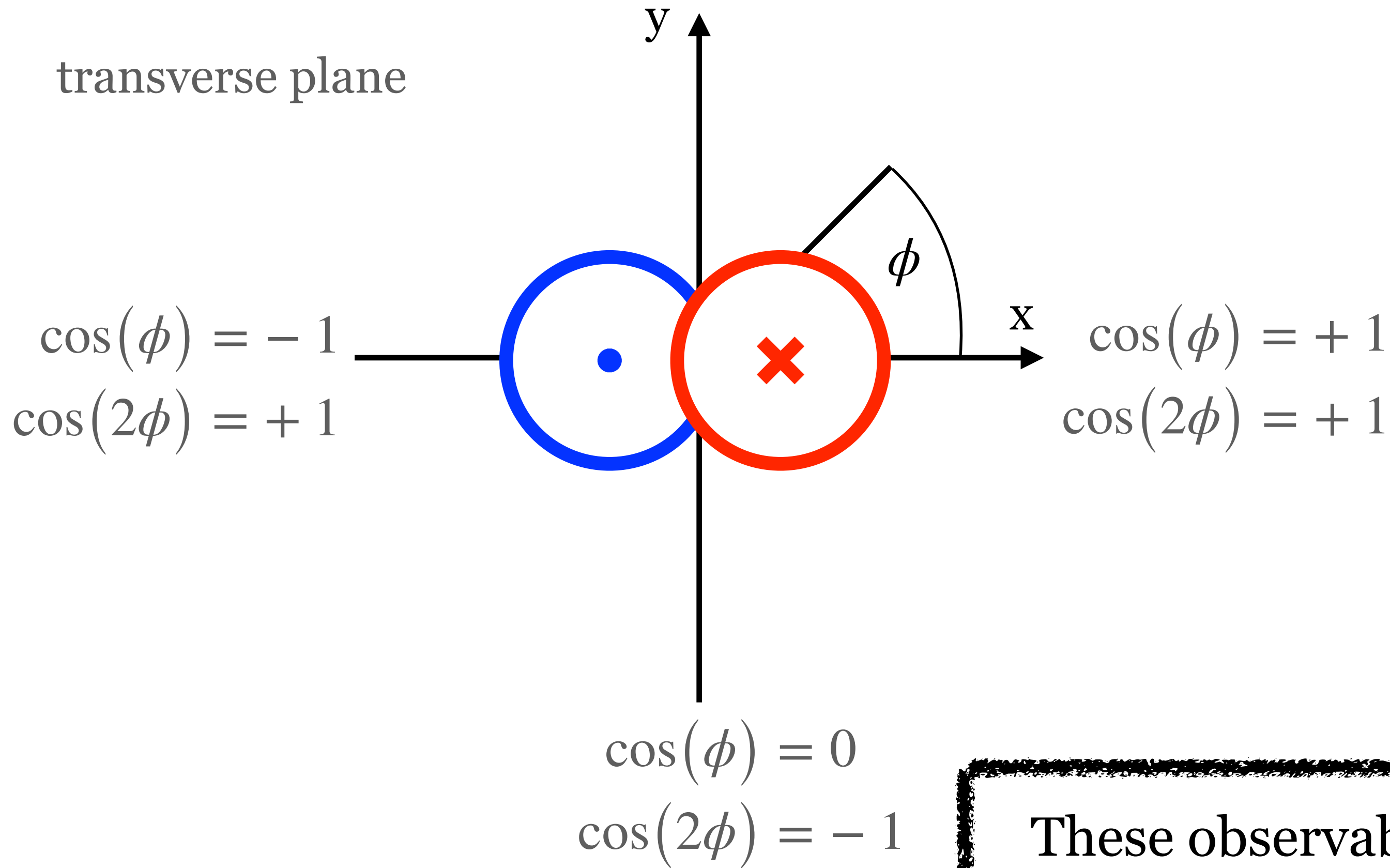


P. Sorensen, "Searching for Superhorizon Fluctuations in Heavy-Ion Collisions", 24th Winter Workshop on Nuclear Dynamics, arXiv:0808.0503

Flow observables in HICs

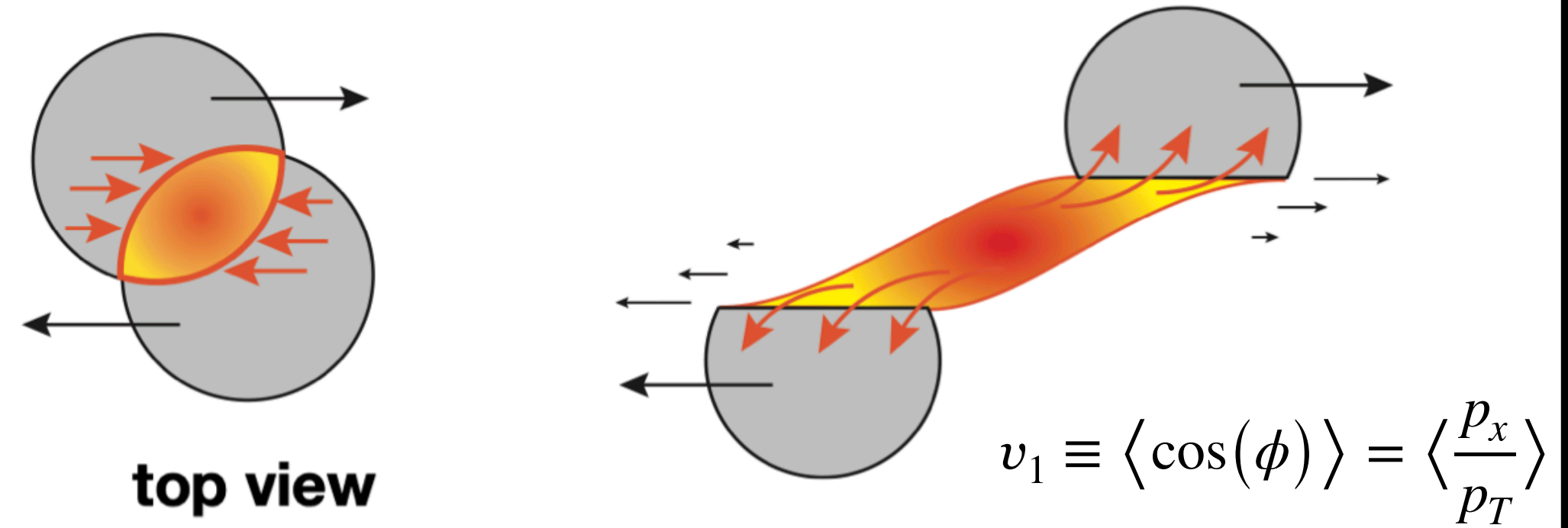
Flow $v_n \equiv \langle \cos(n\phi) \rangle$ $\cos(\phi) = 0$
 $\cos(2\phi) = -1$

transverse plane

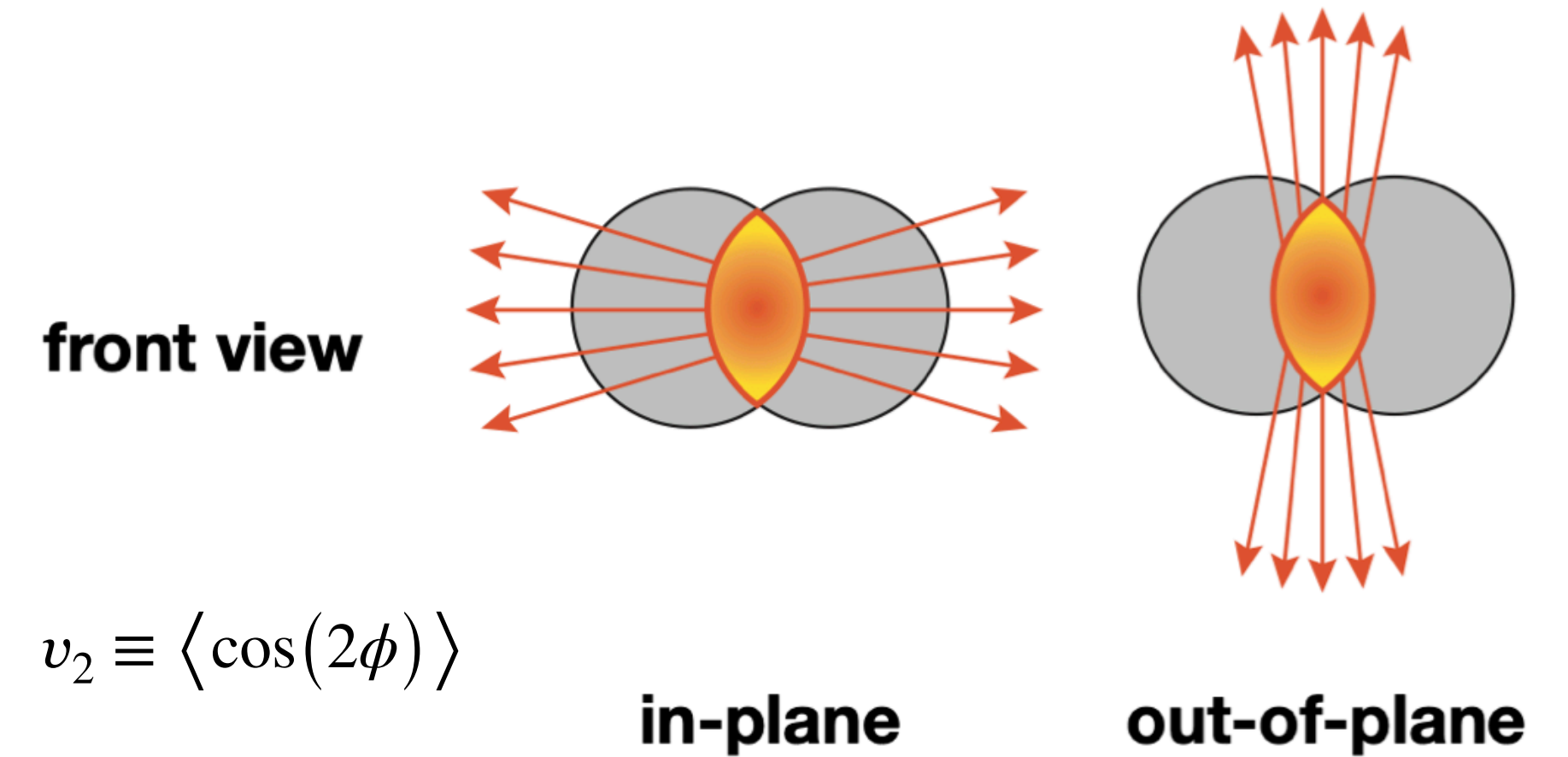


These observables
 are extremely
 sensitive to the EOS

directed flow v_1 ($dv_1/dy \sim$ longitudinal expansion)



elliptic flow v_2 ($v_2(y \approx 0) \sim$ midrapidity)



illustrations from a presentation
 by B. Kardan (HADES)

Evidence for production of QGP in ultra-relativistic HICs

If quarks and gluons are deconfined in high-energy HICs, how would we know?

We only measure final momenta of produced hadrons!

How are hadrons produced from quarks? Consider two proposed processes:

1. Fragmentation: as the system expands (becomes more dilute), distances between quarks increase = quark-quark potential increases \Rightarrow new quark pairs are produced \Rightarrow eventually the produced pairs stay intact

Note: final mesons and baryons have a *fraction* of the original quark momentum

2. Coalescence: two or three quarks very close in position and momentum space combine to form a meson or a baryon

Note: final mesons and baryons have a *sum* of the original quark momentum

Consequently: fragmentation leads to statistically lower momenta of measured particles, while coalescence leads to statistically higher momenta of measured particles

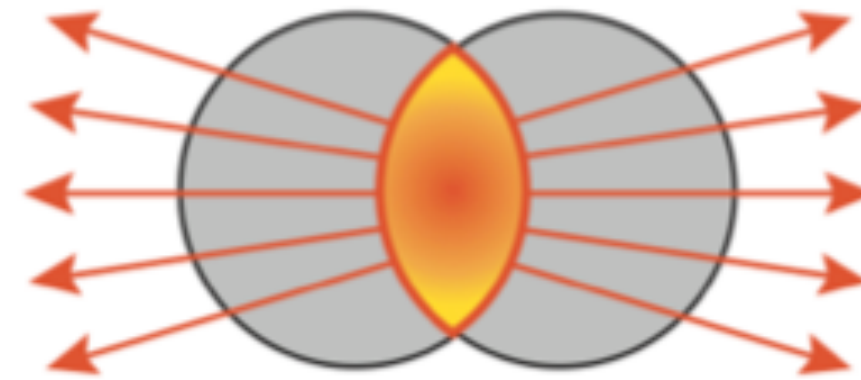
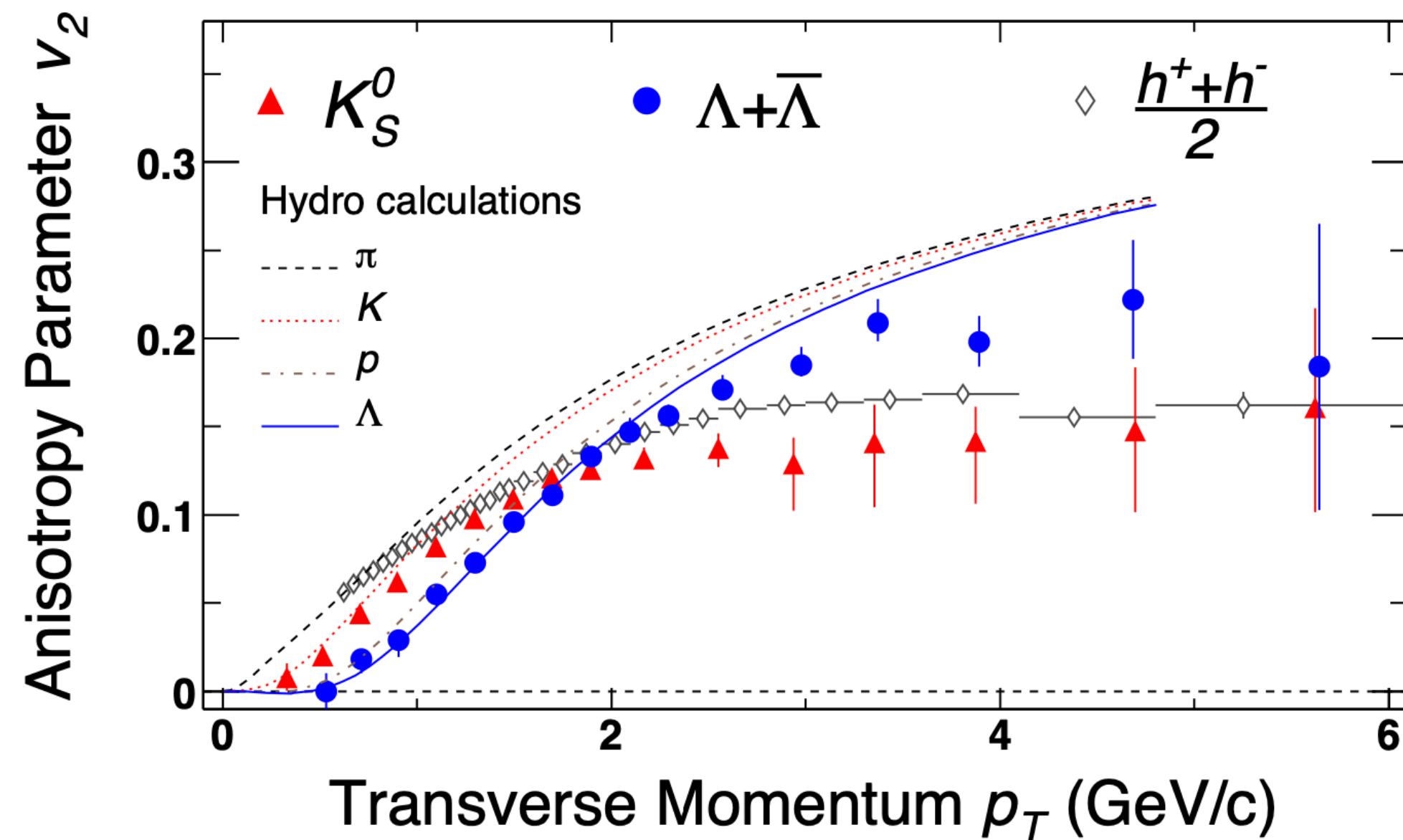
Fragmentation: probability to produce a particle \propto quark density

Coalescence: probability to produce a particle \propto (quark density)² [mesons] or (quark density)³ [baryons]

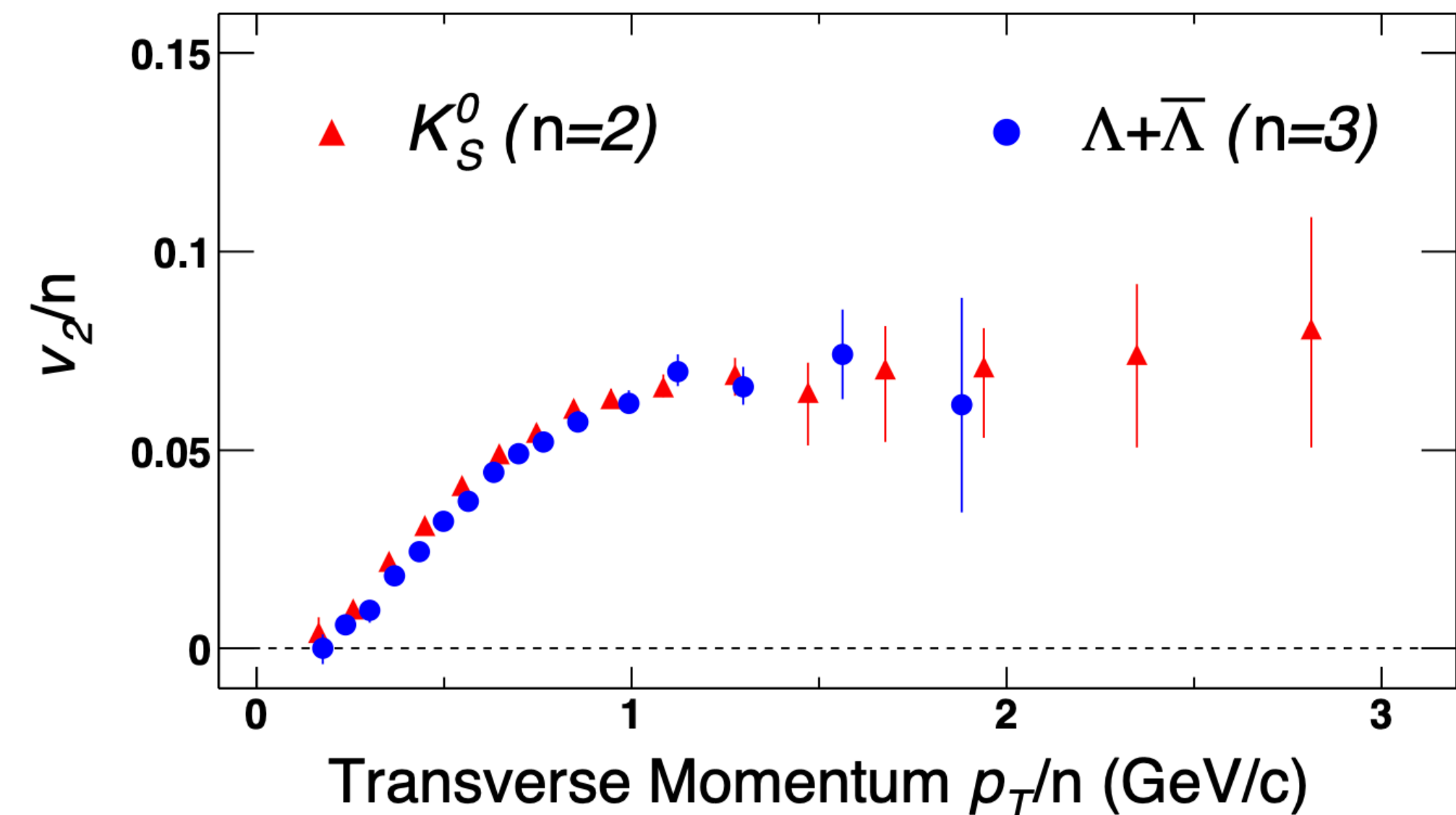
Evidence for production of QGP in ultra-relativistic HICs **from flow**

Fragmentation: statistically lower momenta of measured particles,
probability to produce a particle \propto quark density

Coalescence: statistically higher momenta of measured particles,
probability to produce a particle \propto (quark density)² [mesons] or (quark density)³ [baryons]



Flow of baryons from coalescence
enhanced w.r.t.
flow of mesons from coalescence



STAR Collaboration, "Particle type dependence of azimuthal anisotropy and nuclear modification of particle production in Au + Au collisions at $\sqrt{s(NN)} = 200$ GeV",
Phys. Rev. Lett. 92, 052302 (2004)
[arXiv:nucl-ex/0306007](https://arxiv.org/abs/nucl-ex/0306007)