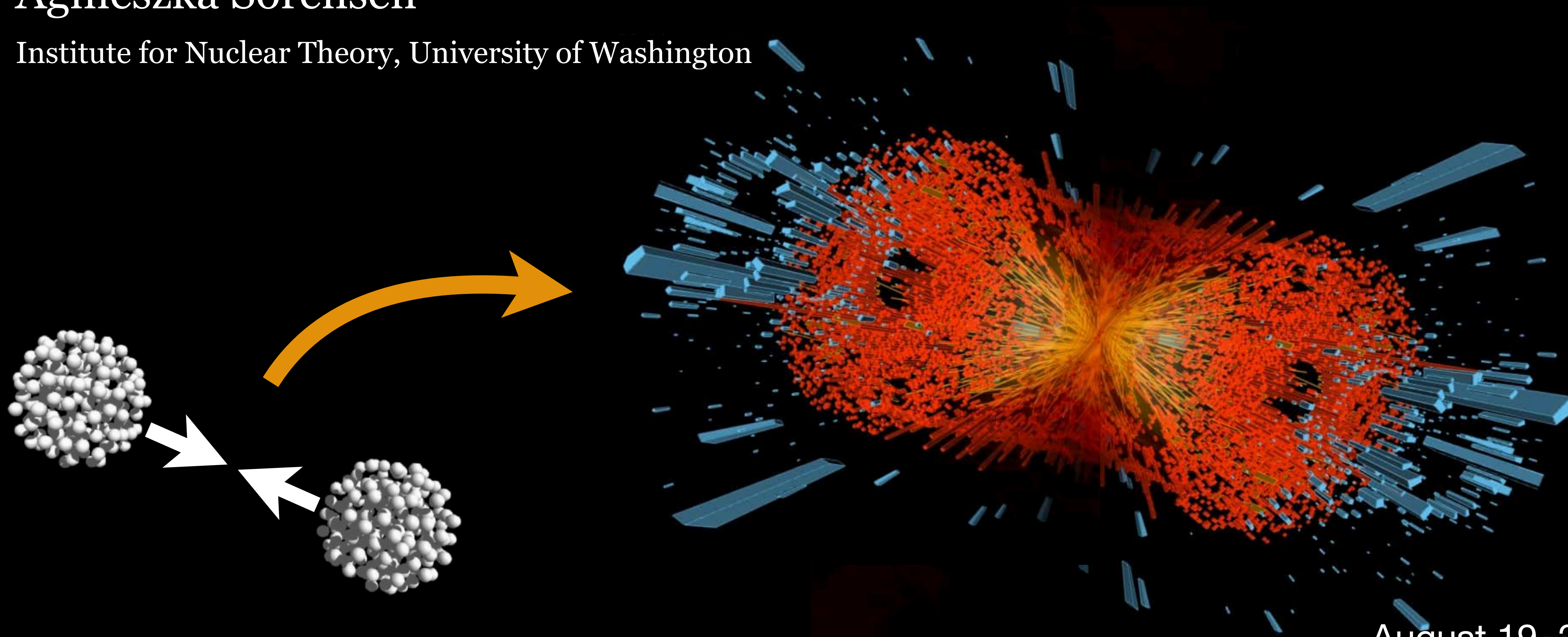


Extracting properties of dense nuclear matter from heavy-ion collisions

Agnieszka Sorensen

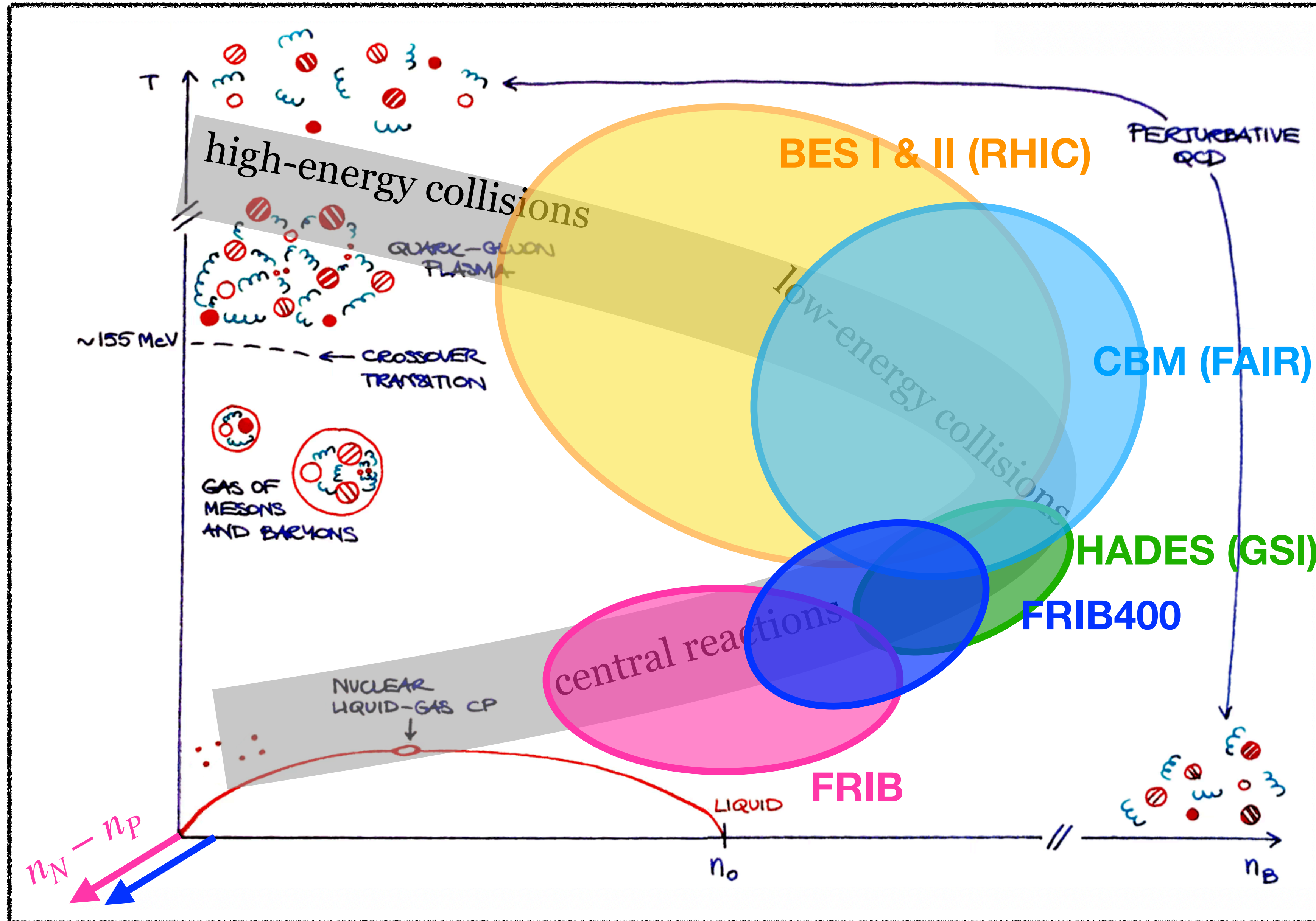
Institute for Nuclear Theory, University of Washington



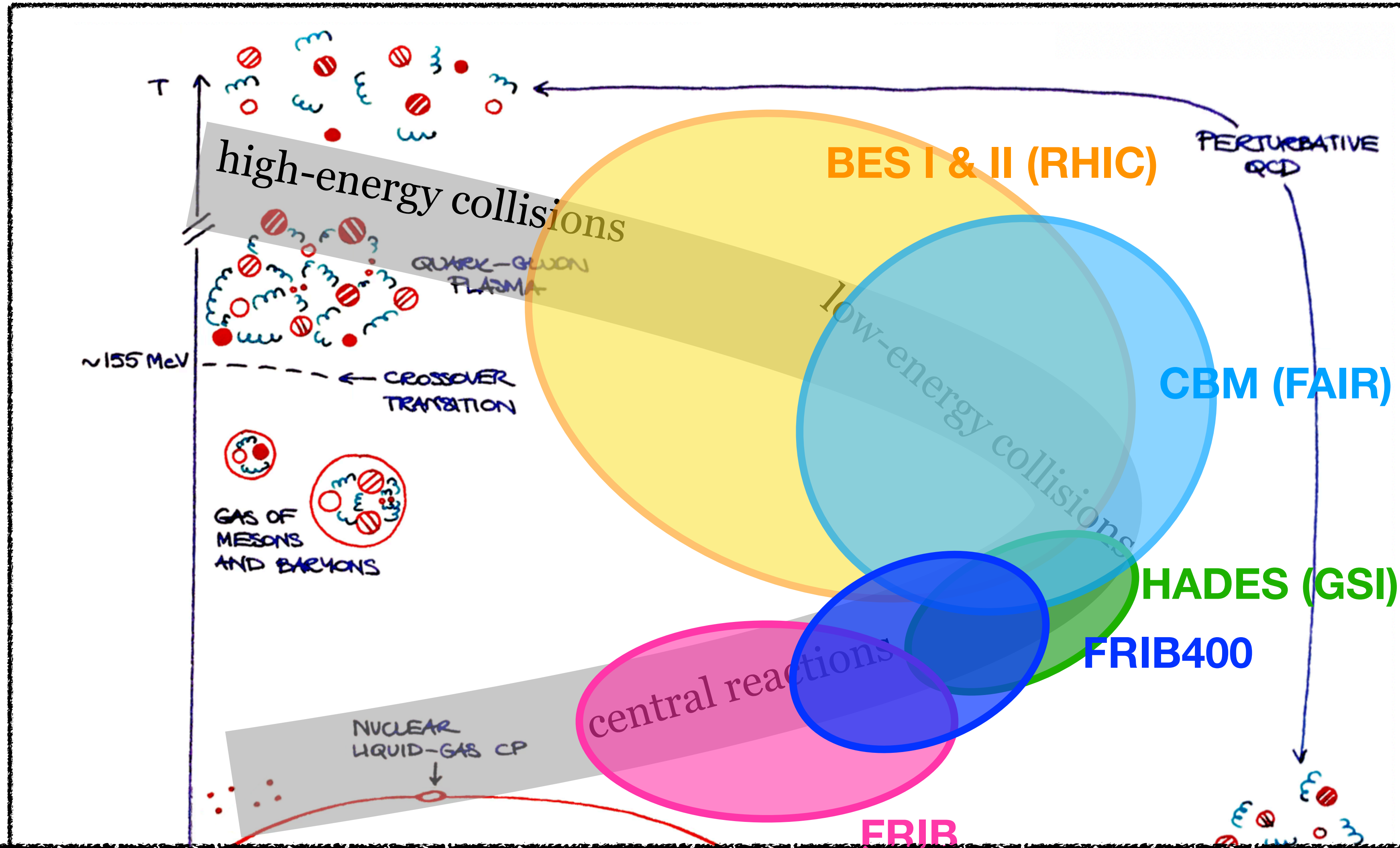
August 19, 2024

14th International Conference on Nucleus-Nucleus Collisions

Heavy-ion collisions = window on properties of dense nuclear matter



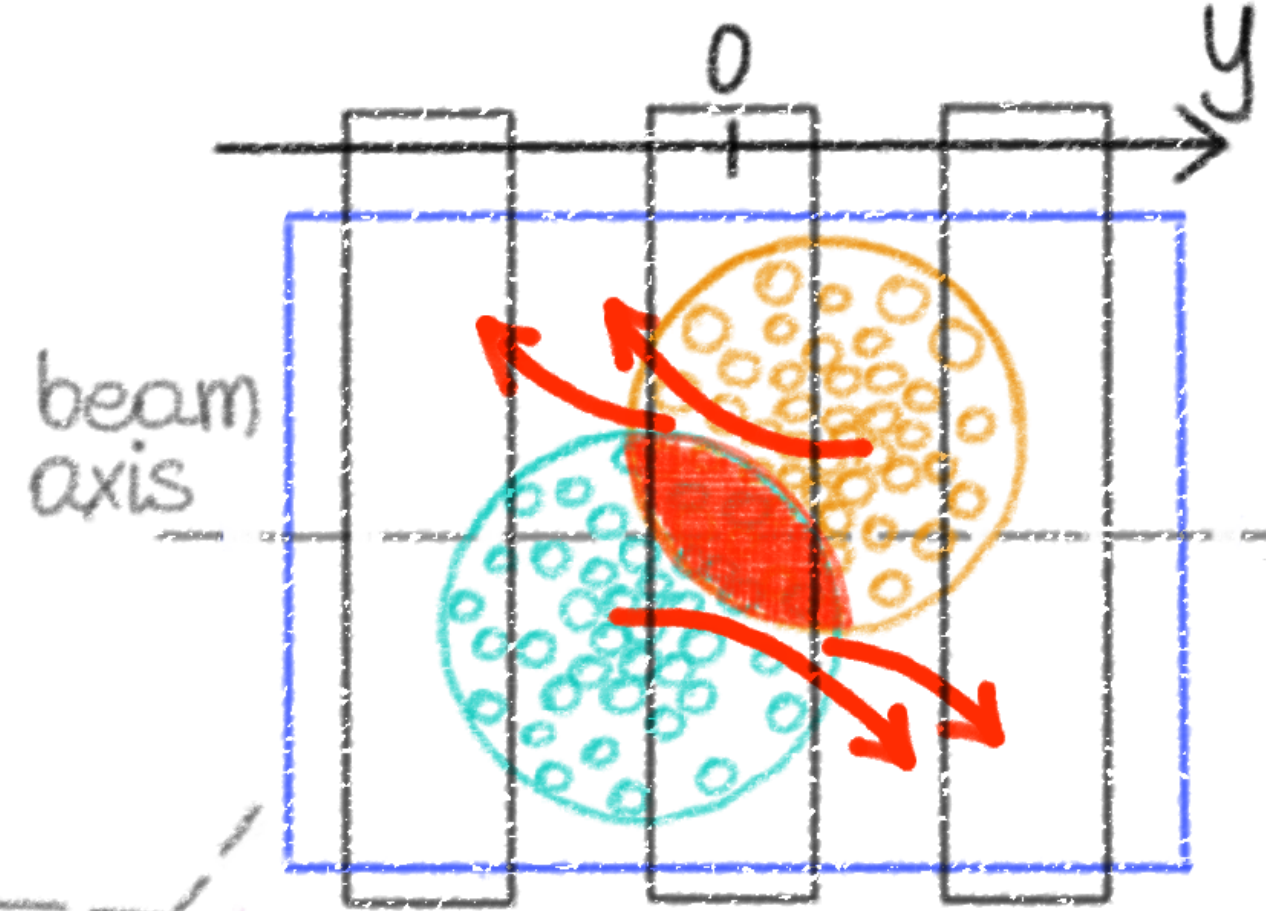
Heavy-ion collisions = window on properties of dense nuclear matter



- HICs = the **only** means to probe densities above n_0 in controlled terrestrial experiments

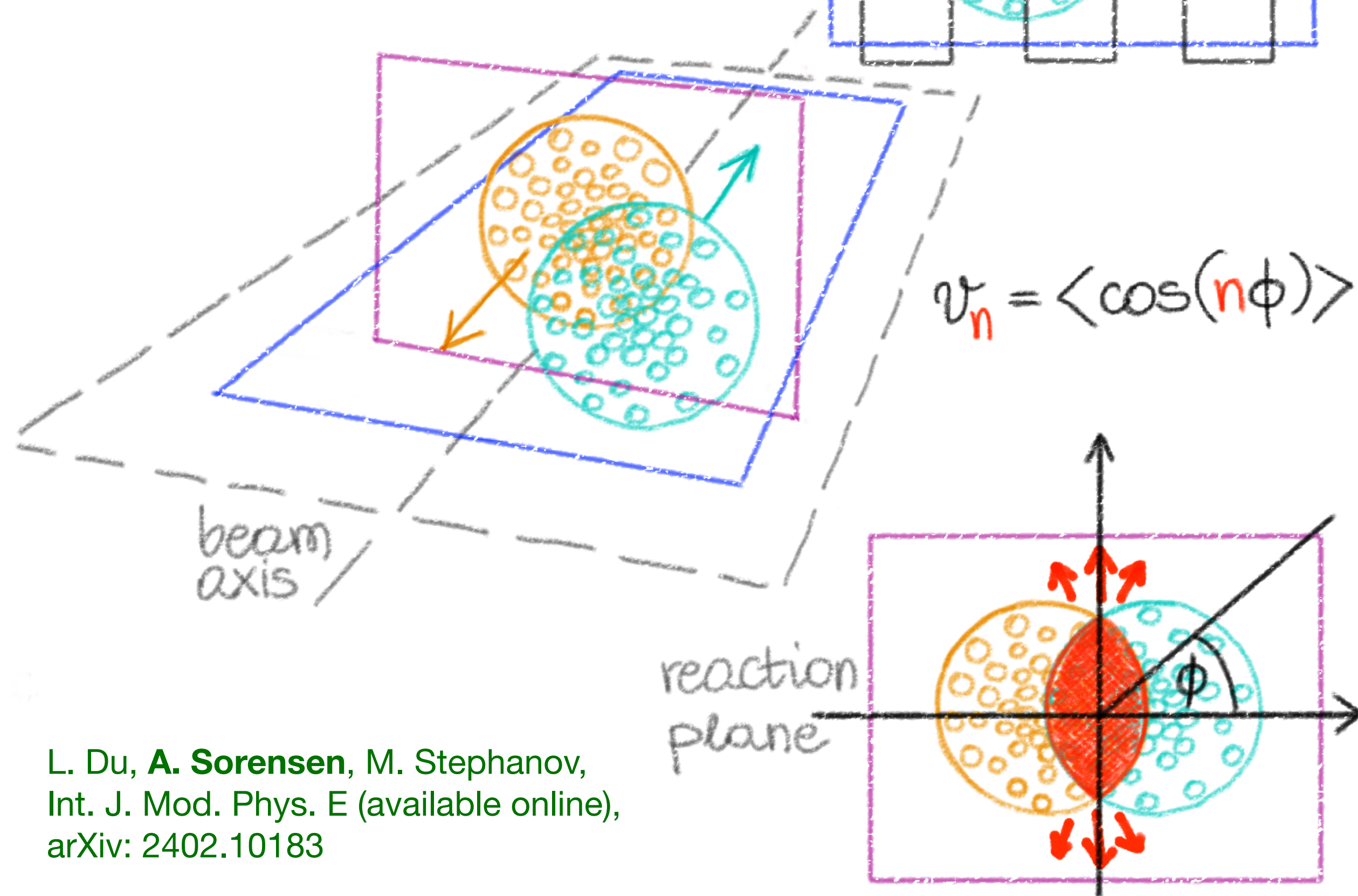
Sketch of a heavy-ion collision evolution and development of flow

* the sketch is informative but not highly realistic



J. Adamczewski-Musch *et al.* (HADES),
Eur.Phys.J.A 59 (2023) 4, 80,
arXiv:2208.02740

$v_1 =$ directed flow

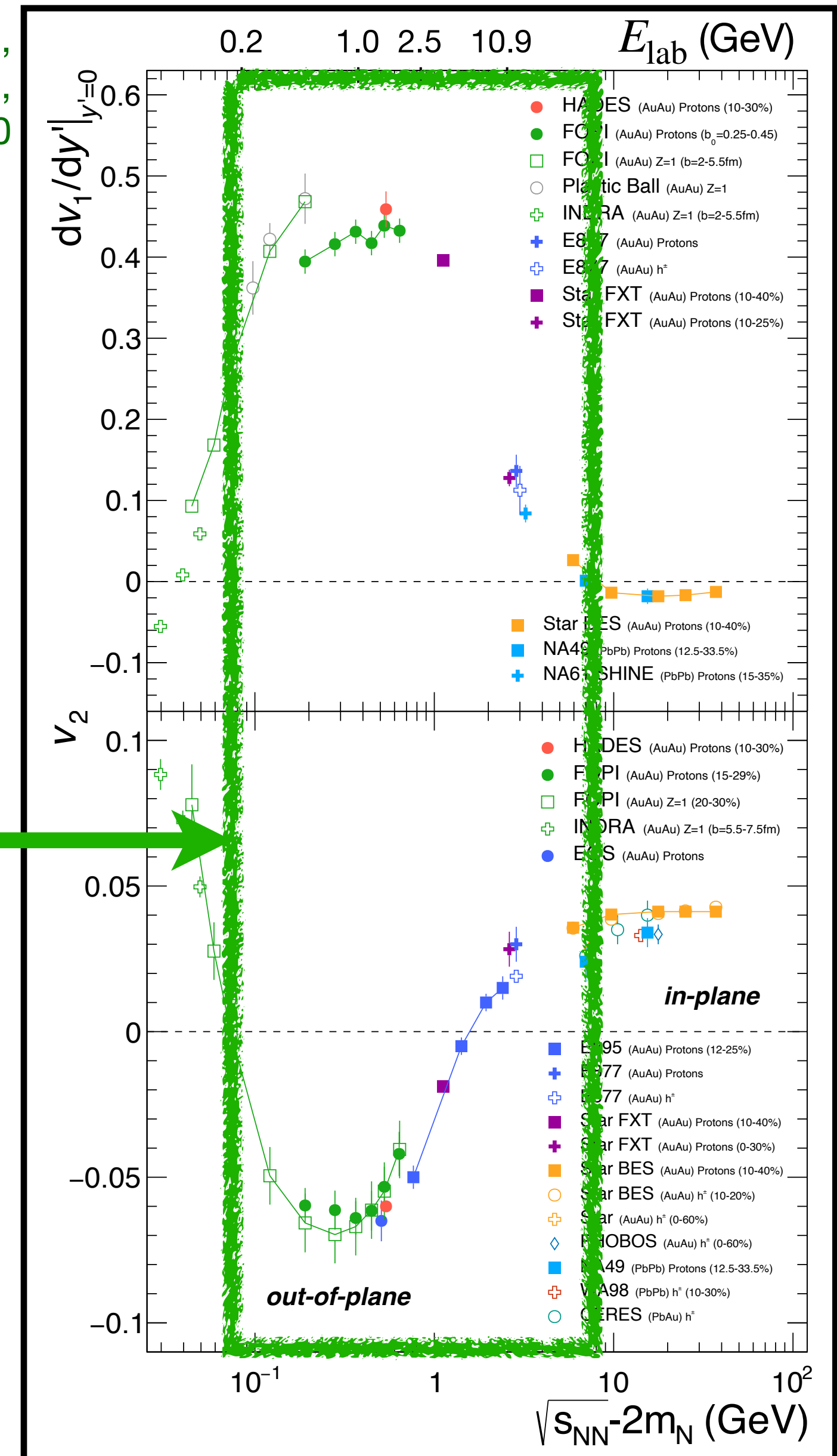


$$v_n = \langle \cos(n\phi) \rangle$$

Experiments:
FRIB & FRIB400,
BES FXT,
HADES, CBM, ...

$v_2 =$ elliptic flow

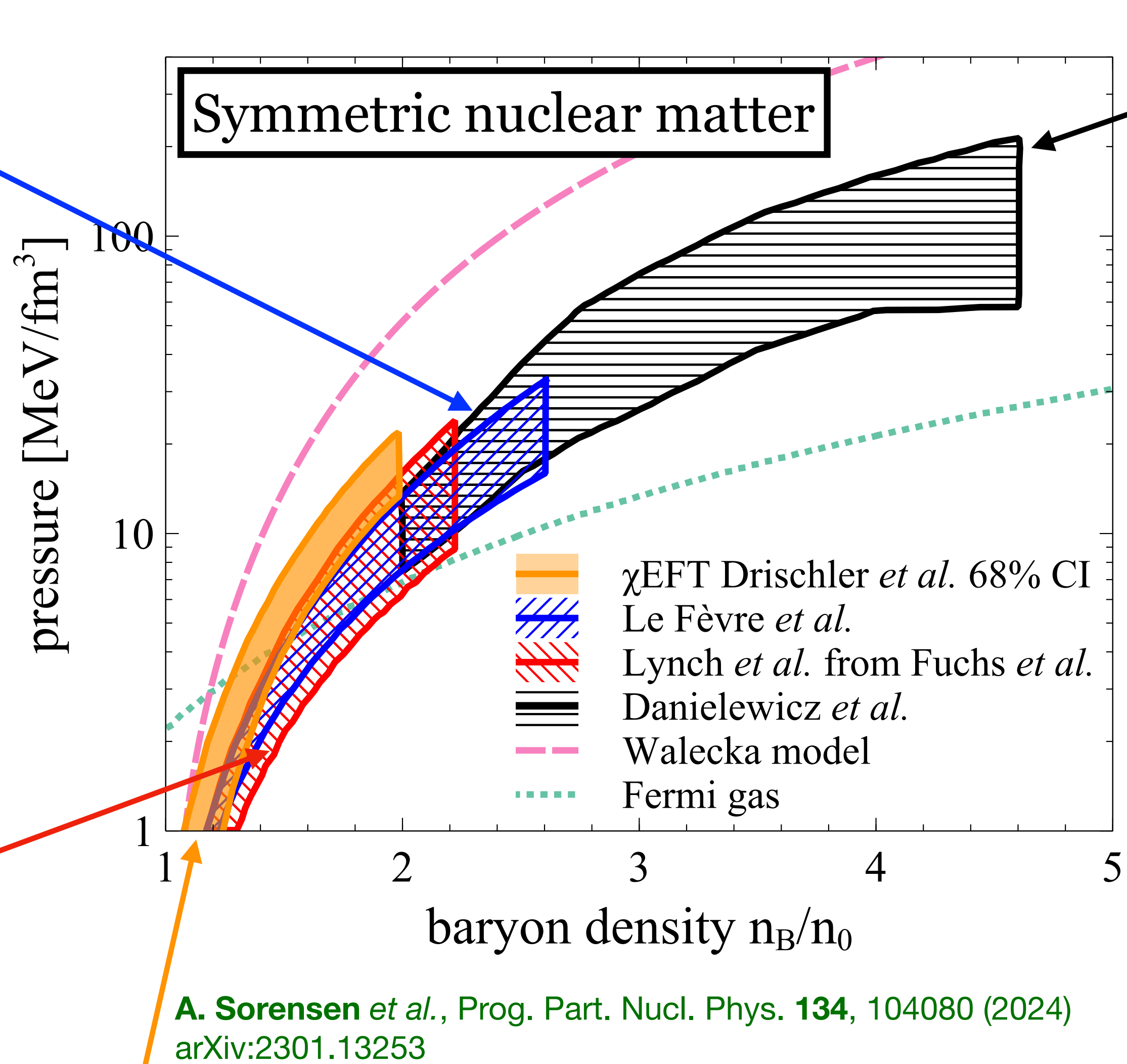
L. Du, A. Sorensen, M. Stephanov,
Int. J. Mod. Phys. E (available online),
arXiv: 2402.10183



Constraints on the EOS come from comparisons to transport models

197Au+197Au @ 0.4–1.5 GeV/u
 ($\sqrt{s_{NN}} = 2.07 - 2.52$ GeV)
 observables: proton flow (FOPI)
 model used: **isospin QMD (IQMD)** w/
 nucleons, Δ , $N^*(1440)$, deuterons, tritons;
 EOS parametrized by K_0 ;
 momentum dependence
 A. Le Fèvre, Y. Leifels, W. Reisdorf, J.
 Aichelin, C. Hartnack, Nucl. Phys. A 945,
 112 (2016), arXiv:1501.05246

197Au+197Au & 12C+12C @ < 1.5 GeV/u
 ($\sqrt{s_{NN}} < 2.5$ GeV)
 observables: subthreshold kaon production
 (KaoS)
 model used: **QMD** w/ nucleons, Δ ,
 $N^*(1440)$, pions, kaons;
 EOS parametrized by K_0 ;
 kaon potentials, momentum dependence
 C. Fuchs *et al.*, Prog. Part. Nucl. Phys. **53**,
 113–124 (2004) arXiv:nucl-th/0312052

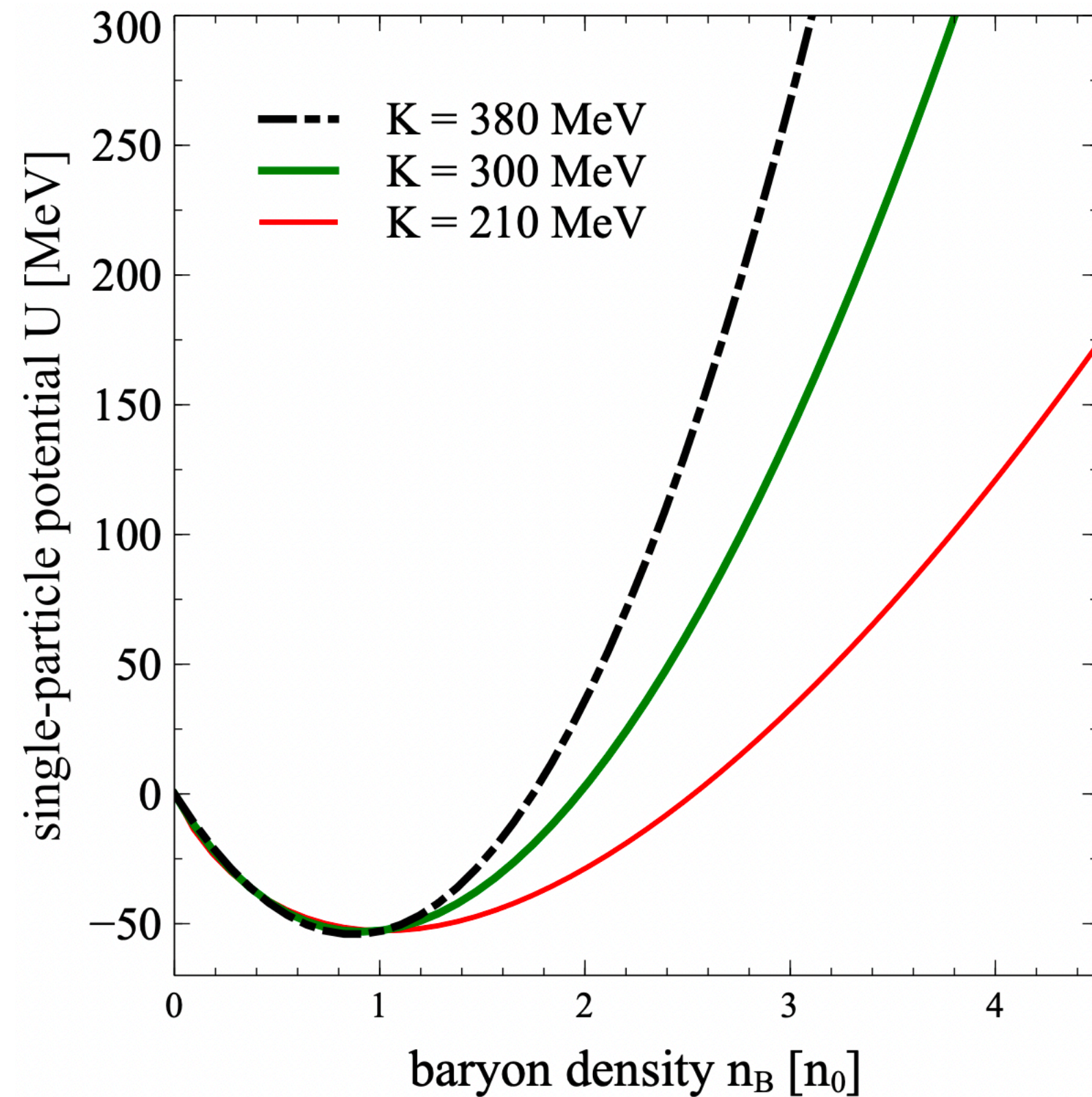


197Au+197Au @ 0.15–10 GeV/u
 ($\sqrt{s_{NN}} = 1.95 - 4.72$ GeV)
 observables: proton flow
 (Plastic Ball, EOS, E877, E895)
 model used: **pBUU** w/ nucleons, Δ ,
 $N^*(1440)$, pions;
 EOS parametrized by K_0 ;
 momentum dependence
 P. Danielewicz, R. Lacey, W. G. Lynch,
 Science **298**, 1592–1596 (2002)

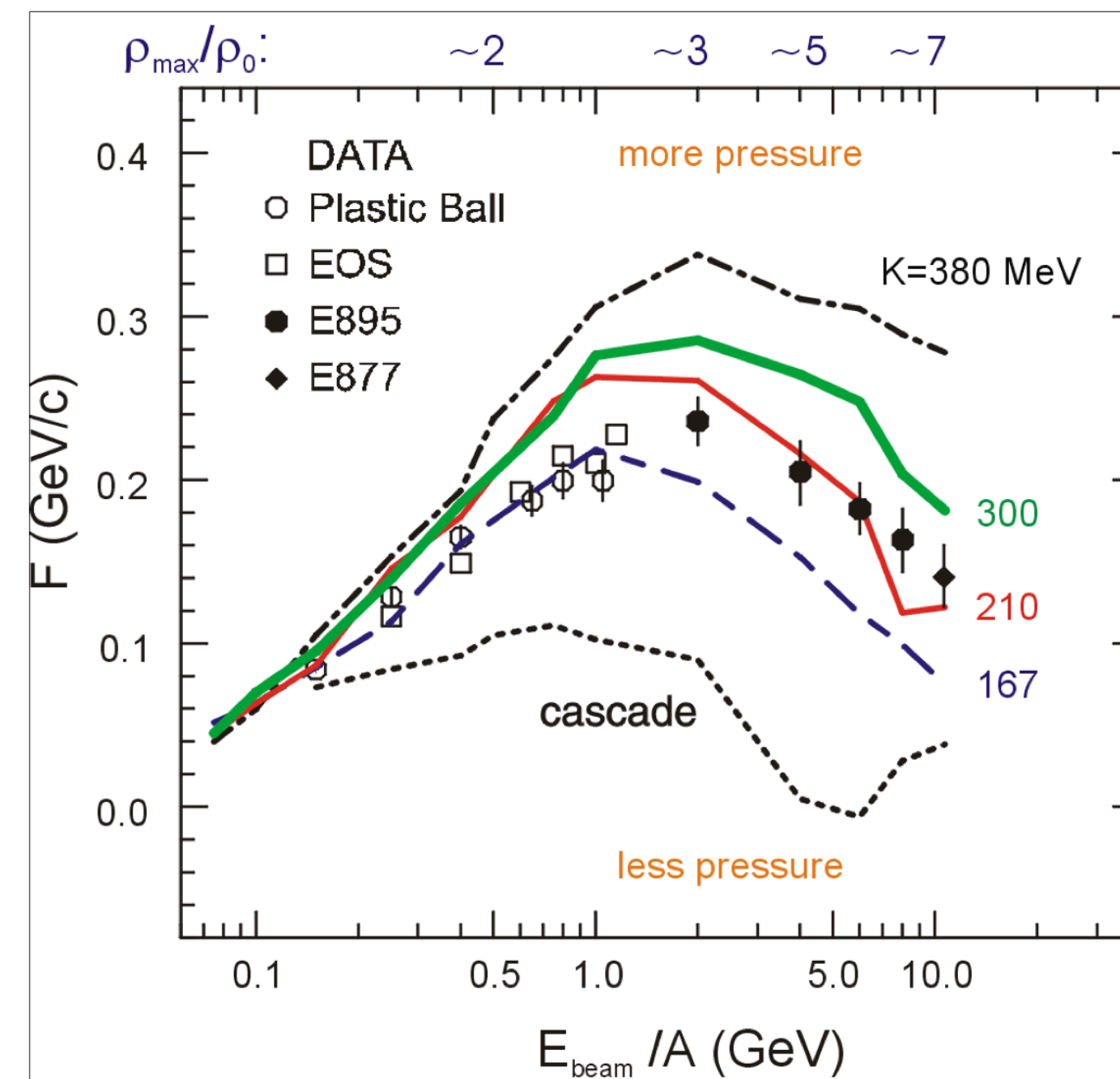
χ EFT
 C. Drischler *et al.*, Phys. Rev. C **102** 5, 054315 (2020)
 arXiv:2004.07805

Standard way of modeling the EOS: Skyrme potential

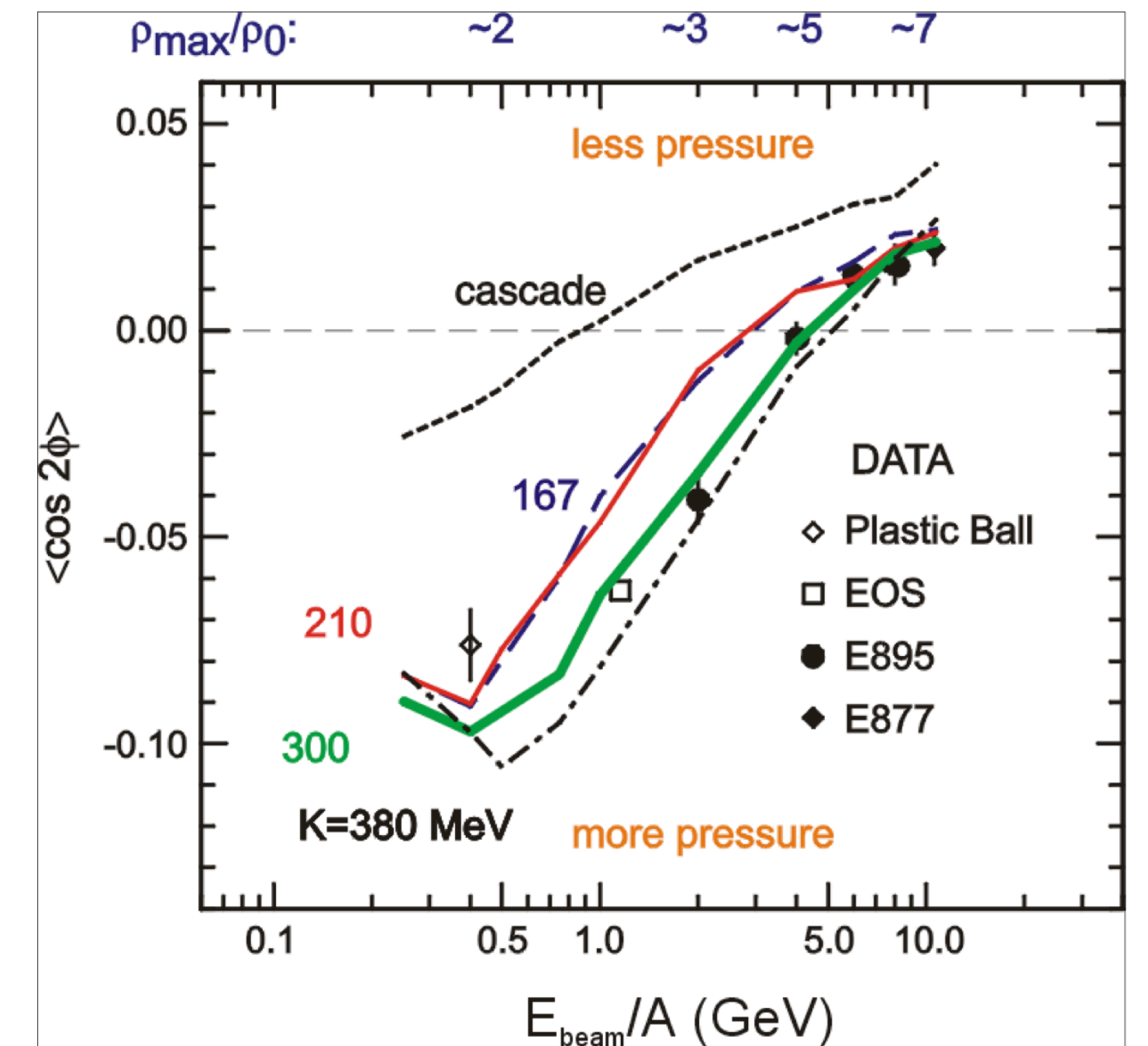
The most common form of the EOS is the “Skyrme potential”: $U(n_B) = A \left(\frac{n_B}{n_0} \right) + B \left(\frac{n_B}{n_0} \right)^\tau$



$$F = \left. \frac{d\langle p_x/A \rangle}{d(y/y_{cm})} \right|_{v/v_1 = 1} \sim \frac{dv_1}{dy}$$



v_2

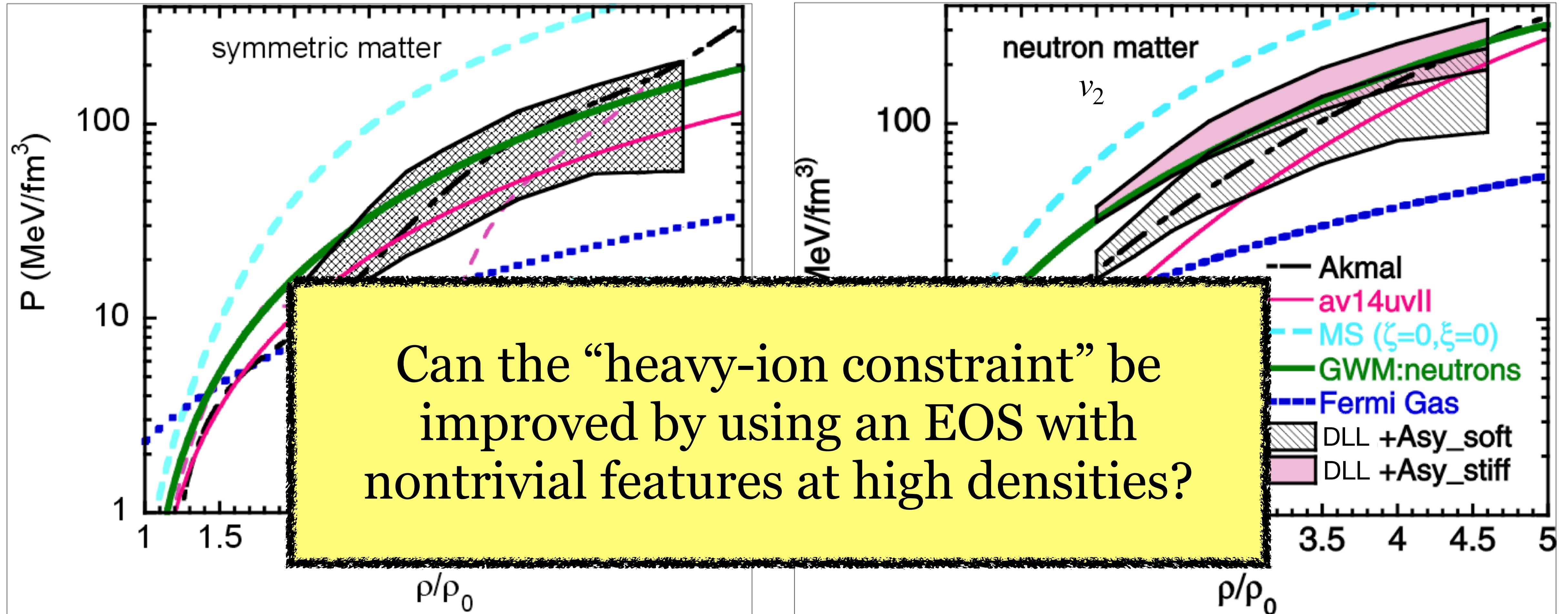


P. Danielewicz, R. Lacey, W. G. Lynch,
 Science **298**, 1592–1596 (2002), arXiv:nucl-th/0208016

Standard way of modeling the EOS: Skyrme potential

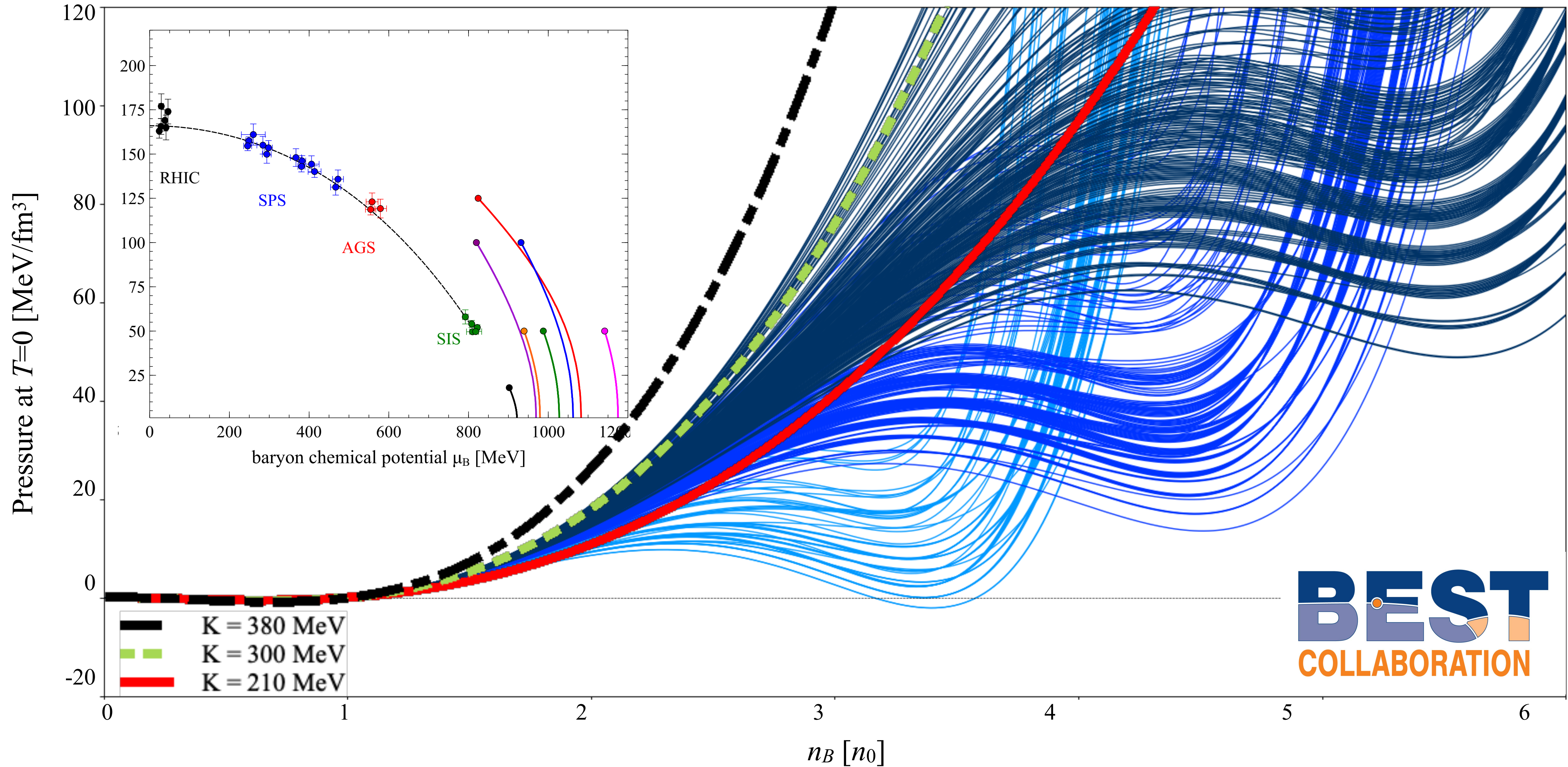
P. Danielewicz, R. Lacey, W. G. Lynch,
Science **298**, 1592–1596 (2002), arXiv:nucl-th/0208016

“the heavy-ion constraint”



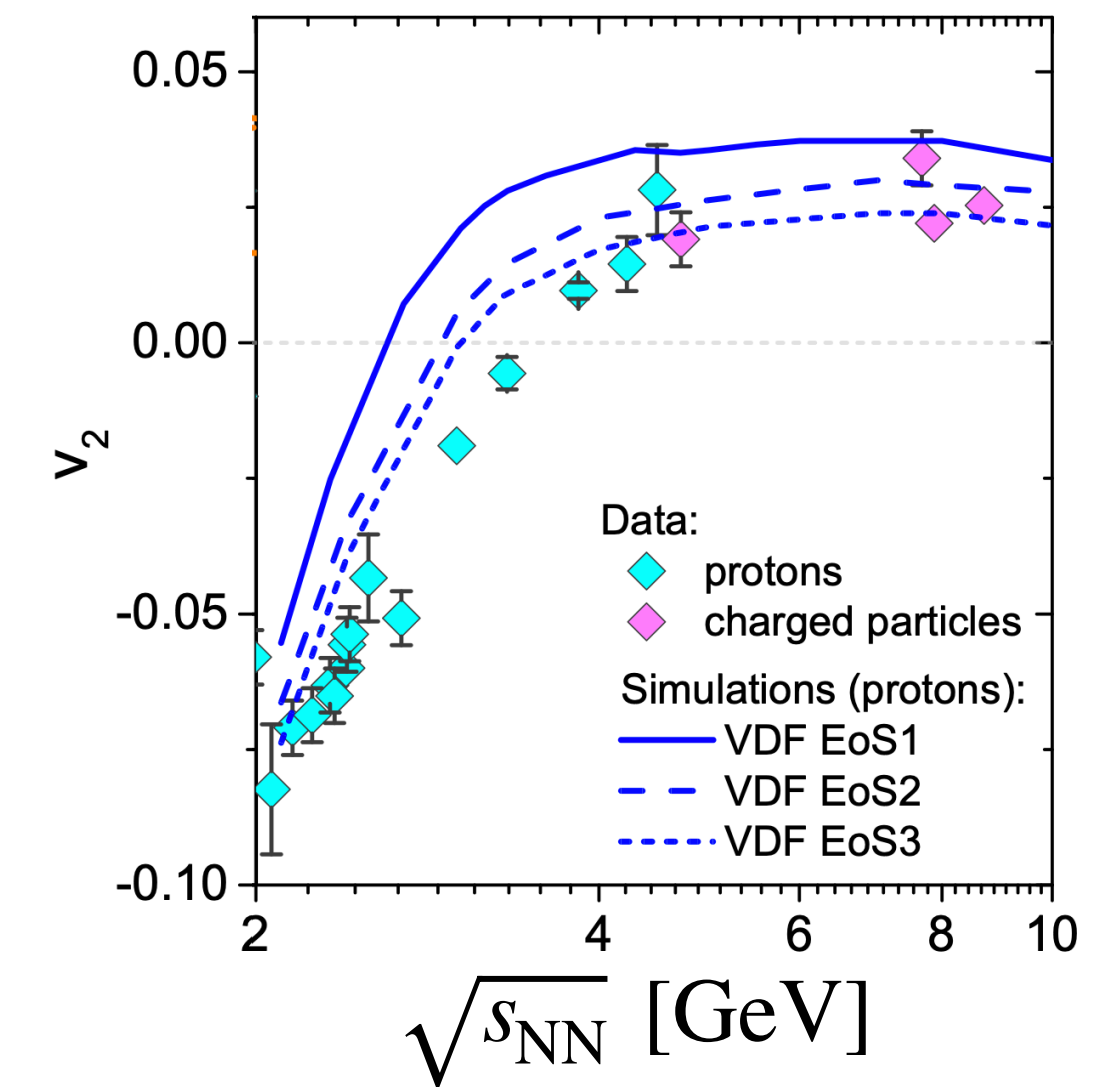
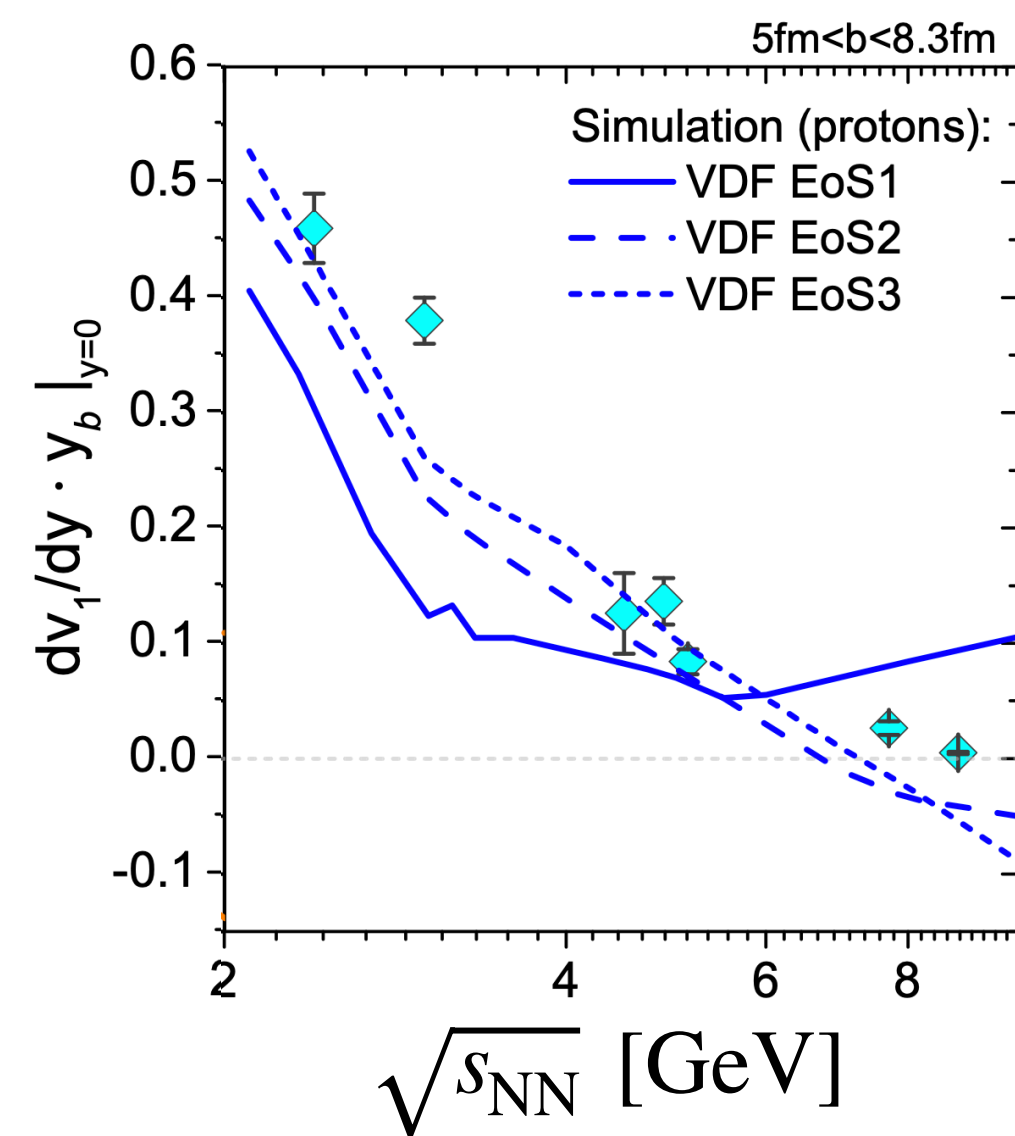
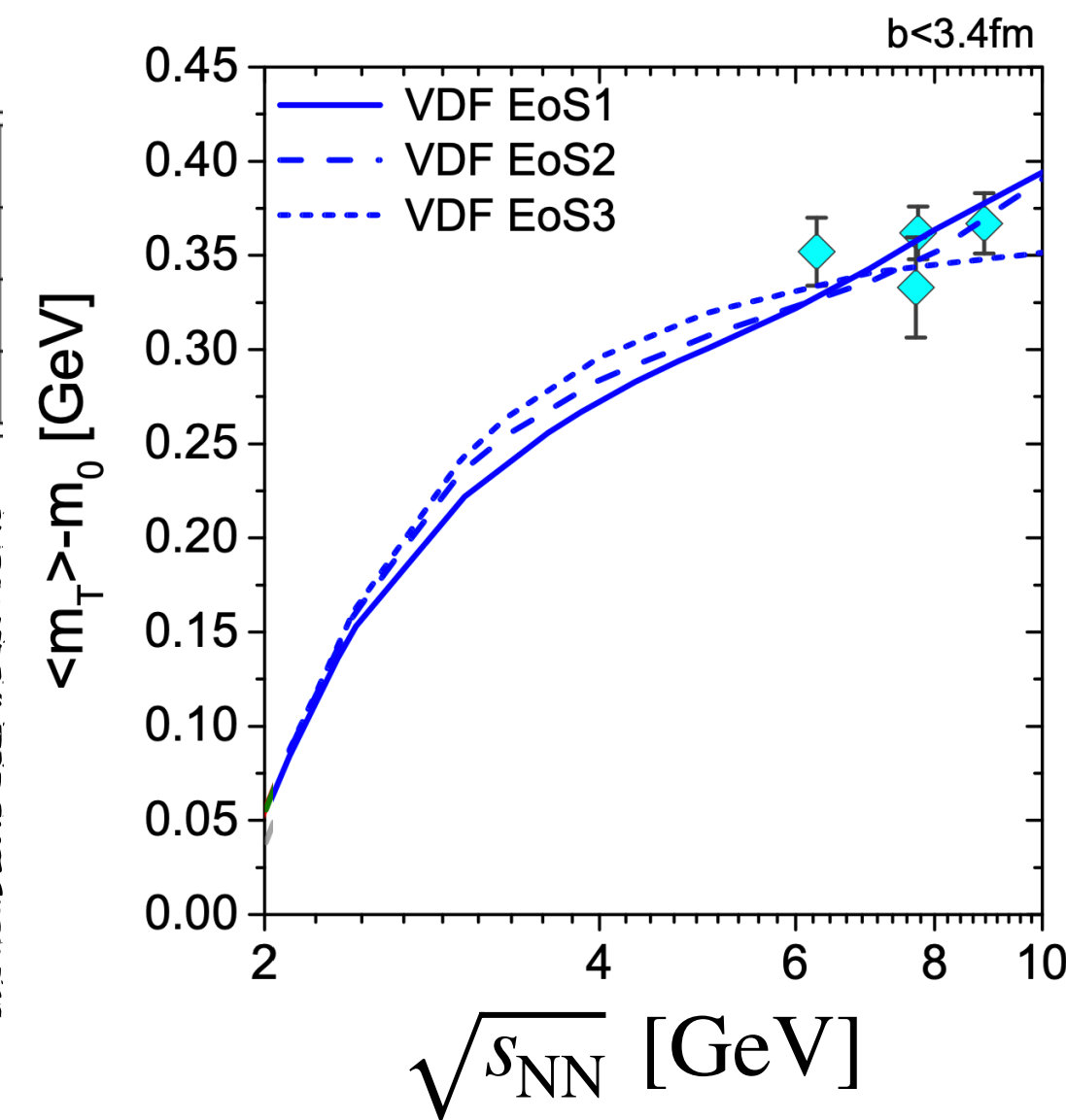
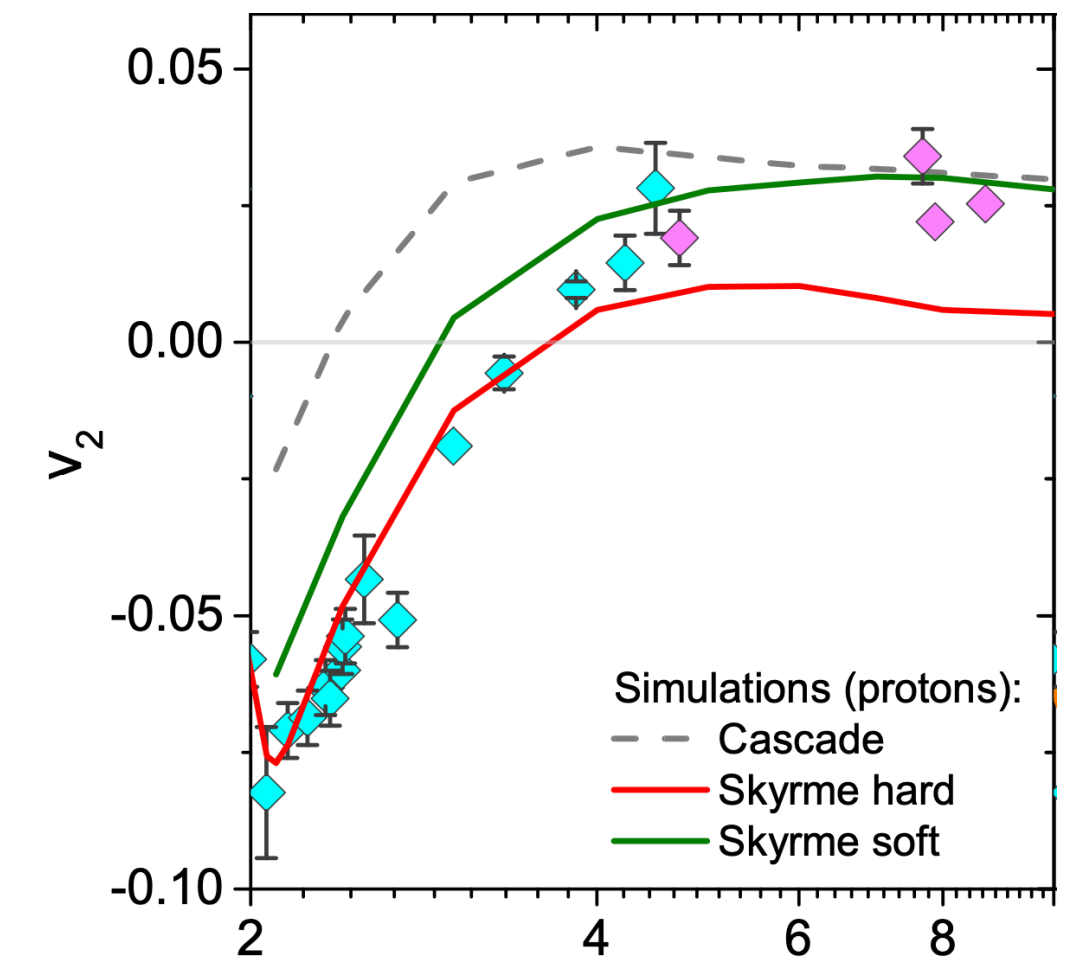
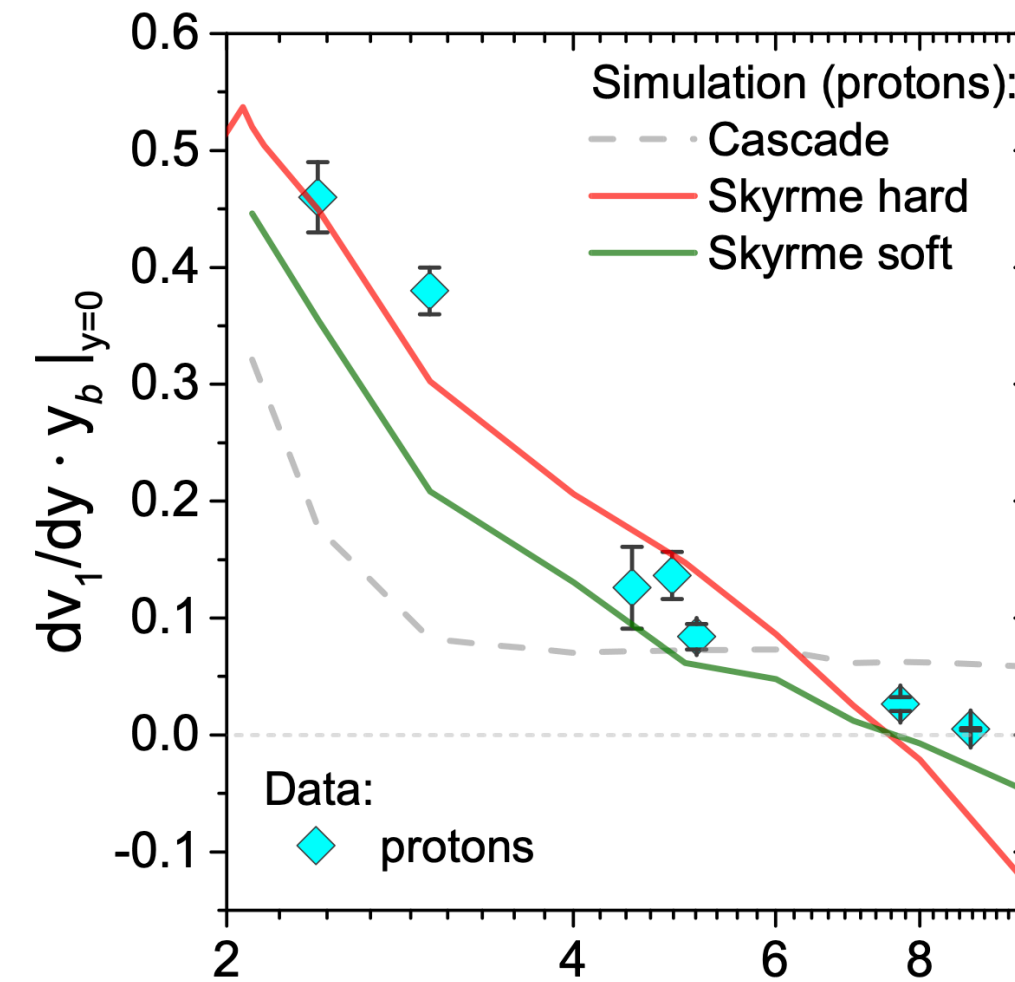
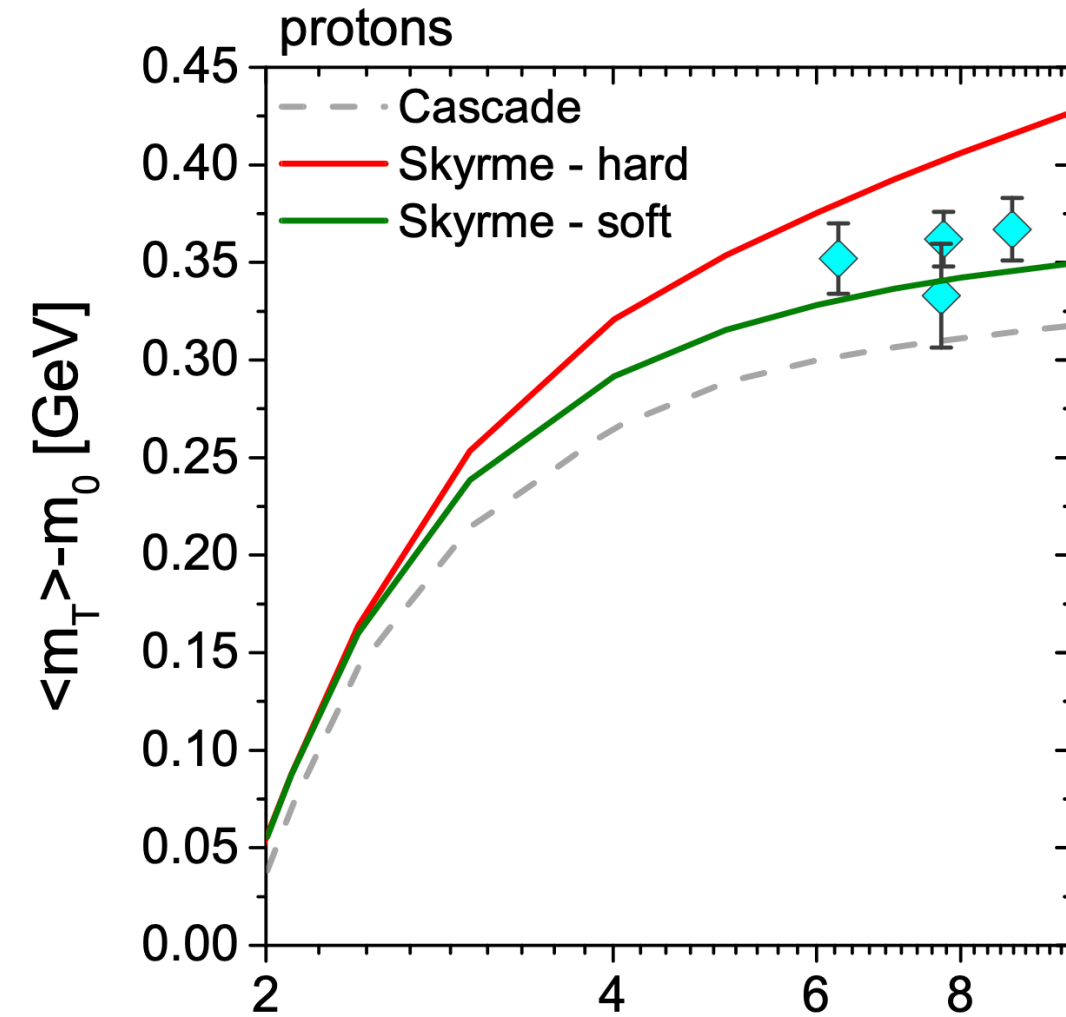
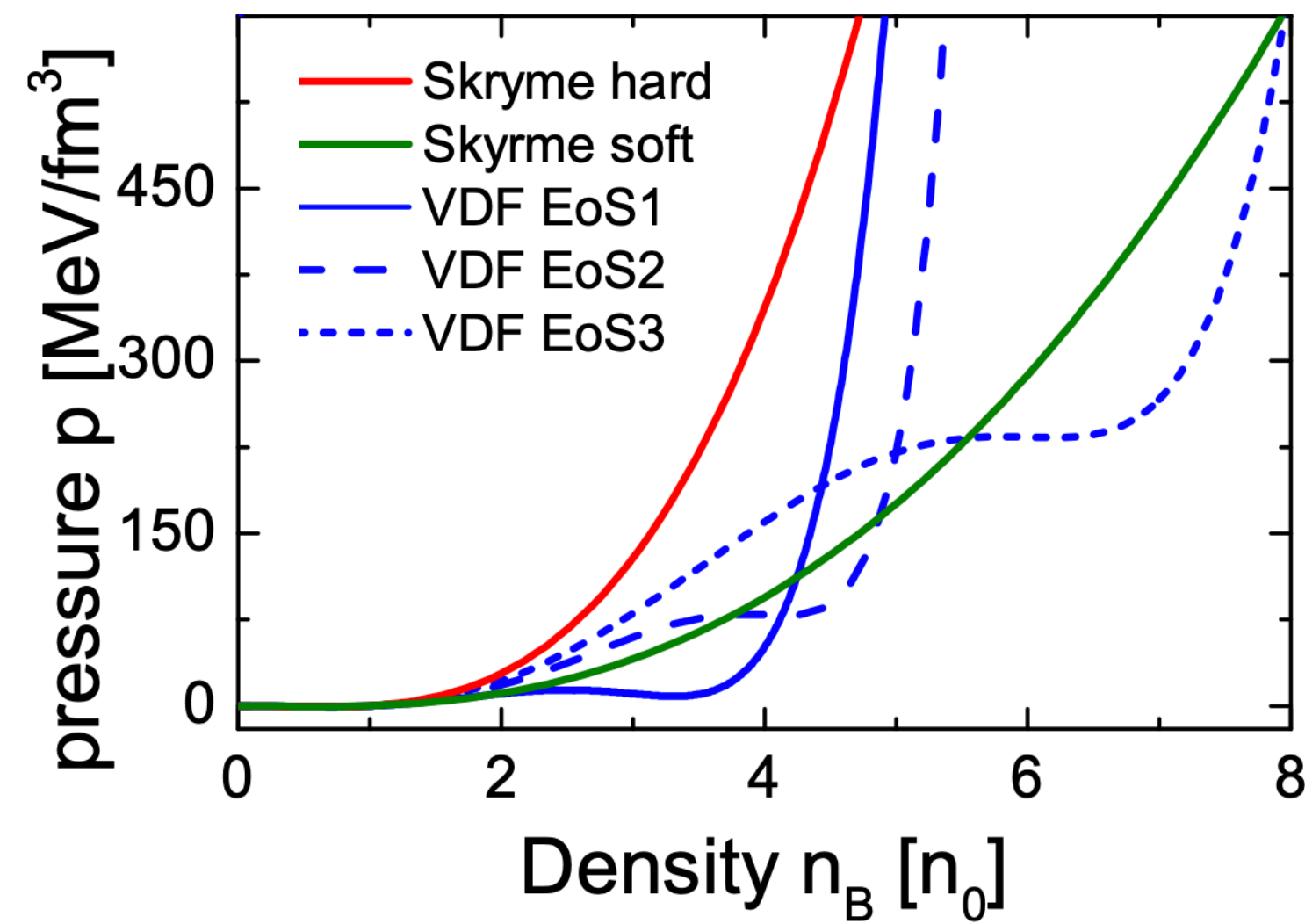
VDF model: relativistic potentials with **two** 1st order phase transitions

A. Sorensen, V. Koch, Phys. Rev. C **104** (2021) 3, 034904, arXiv:2011.06635



Results from UrQMD with (non-relativistic) VDF

J. Steinheimer, A. Motornenko, **A. Sorensen**, Y. Nara, V. Koch, M. Bleicher,
 Eur. Phys. J. C **82**, 10, 911 (2022) arXiv:2208.12091

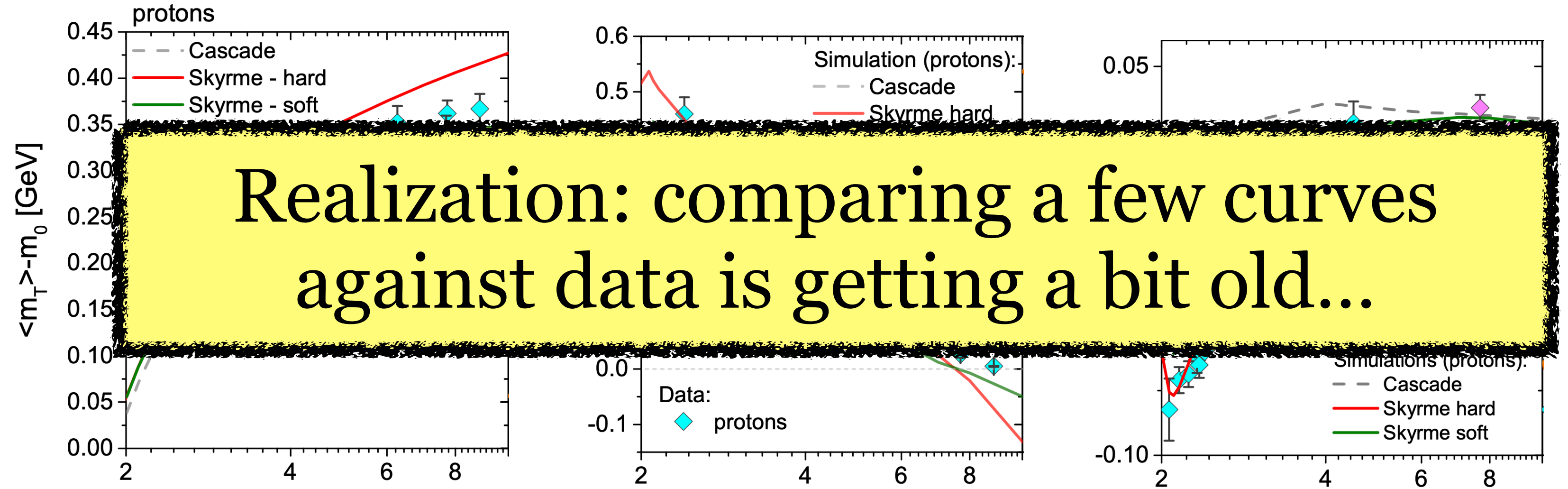
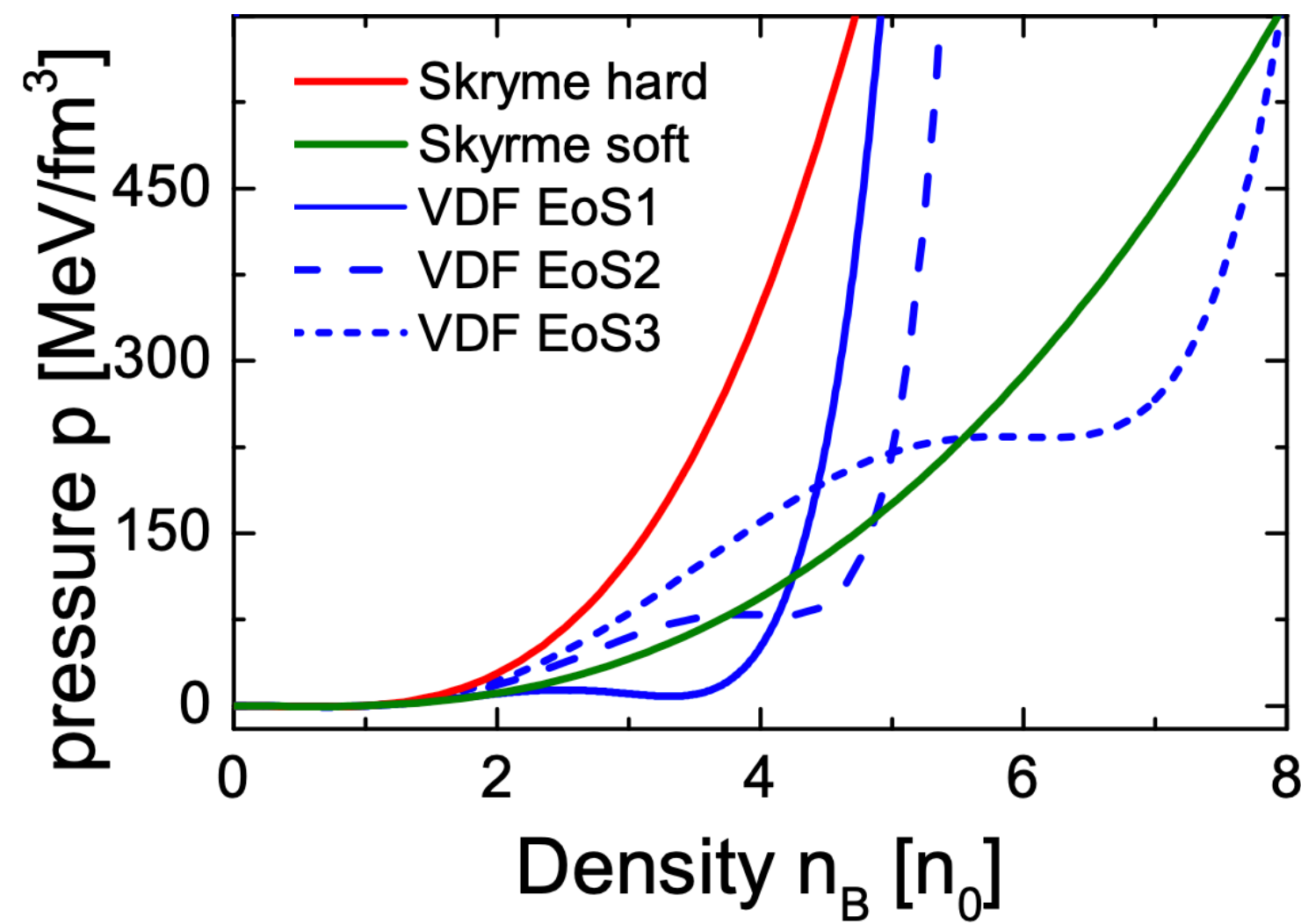


EoS	$T_c^{(N)}$ [MeV]	$n_c^{(Q)}$ [n_0]	$T_c^{(Q)}$ [MeV]	K_0 [MeV]
VDF1	18	3.0	100	261
VDF2	18	4.0	50	279
VDF3	22	6.0	50	356

Very soft EOS at $n_B \in (2,3)n_0$
 not supported in VDF+UrQMD

Results from UrQMD with (non-relativistic) VDF

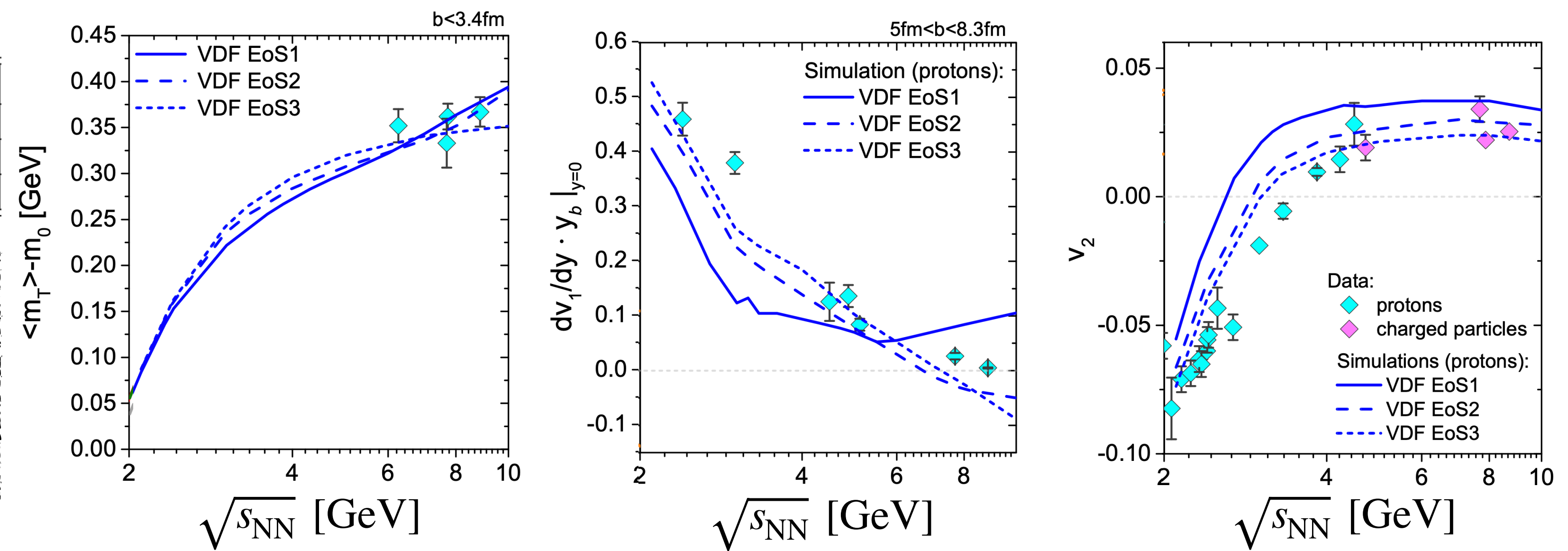
J. Steinheimer, A. Motornenko, **A. Sorensen**, Y. Nara, V. Koch, M. Bleicher,
 Eur. Phys. J. C **82**, 10, 911 (2022) arXiv:2208.12091



Realization: comparing a few curves against data is getting a bit old...

EoS	$T_c^{(N)}$ [MeV]	$n_c^{(Q)}$ [n_0]	$T_c^{(Q)}$ [MeV]	K_0 [MeV]
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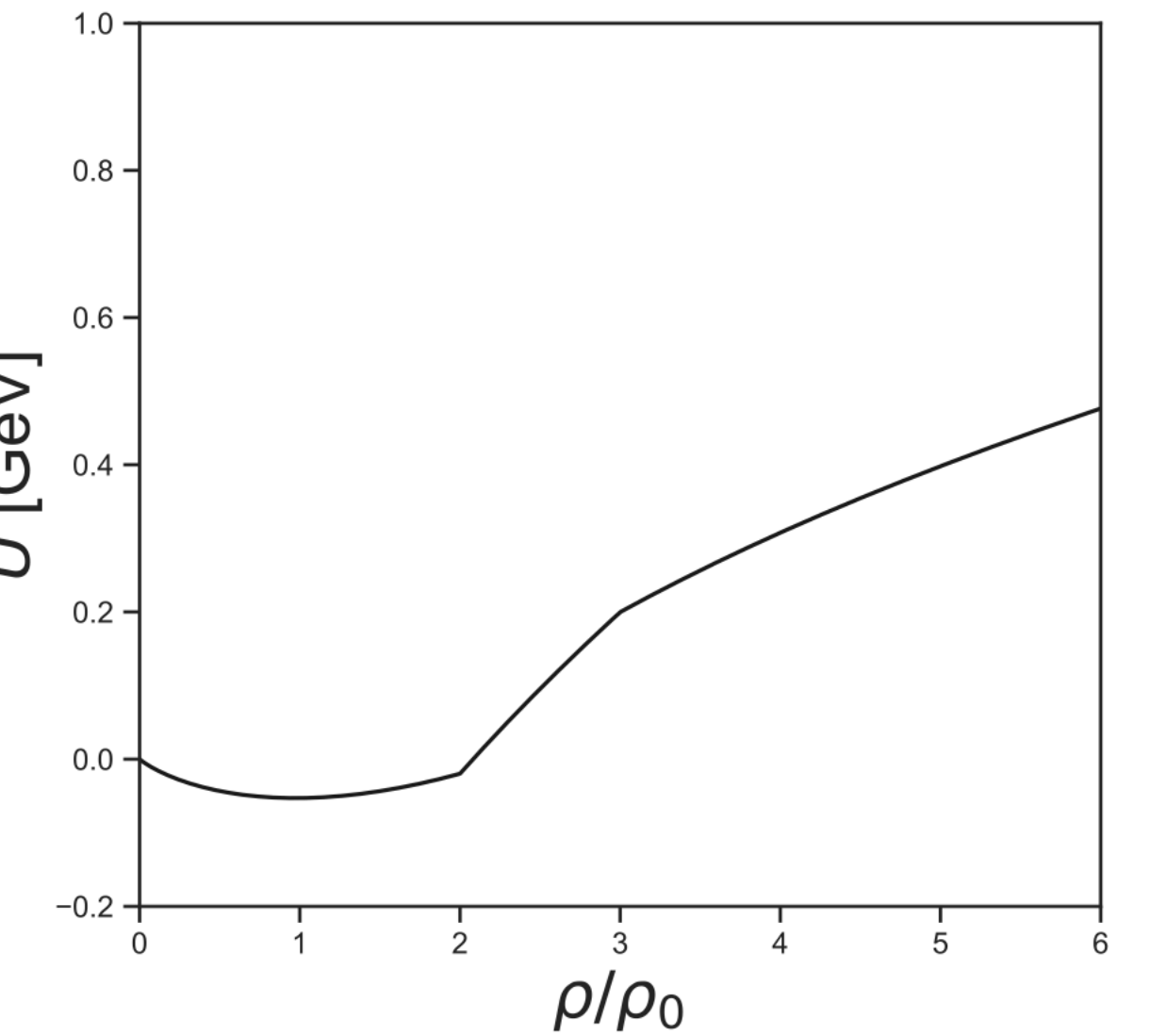
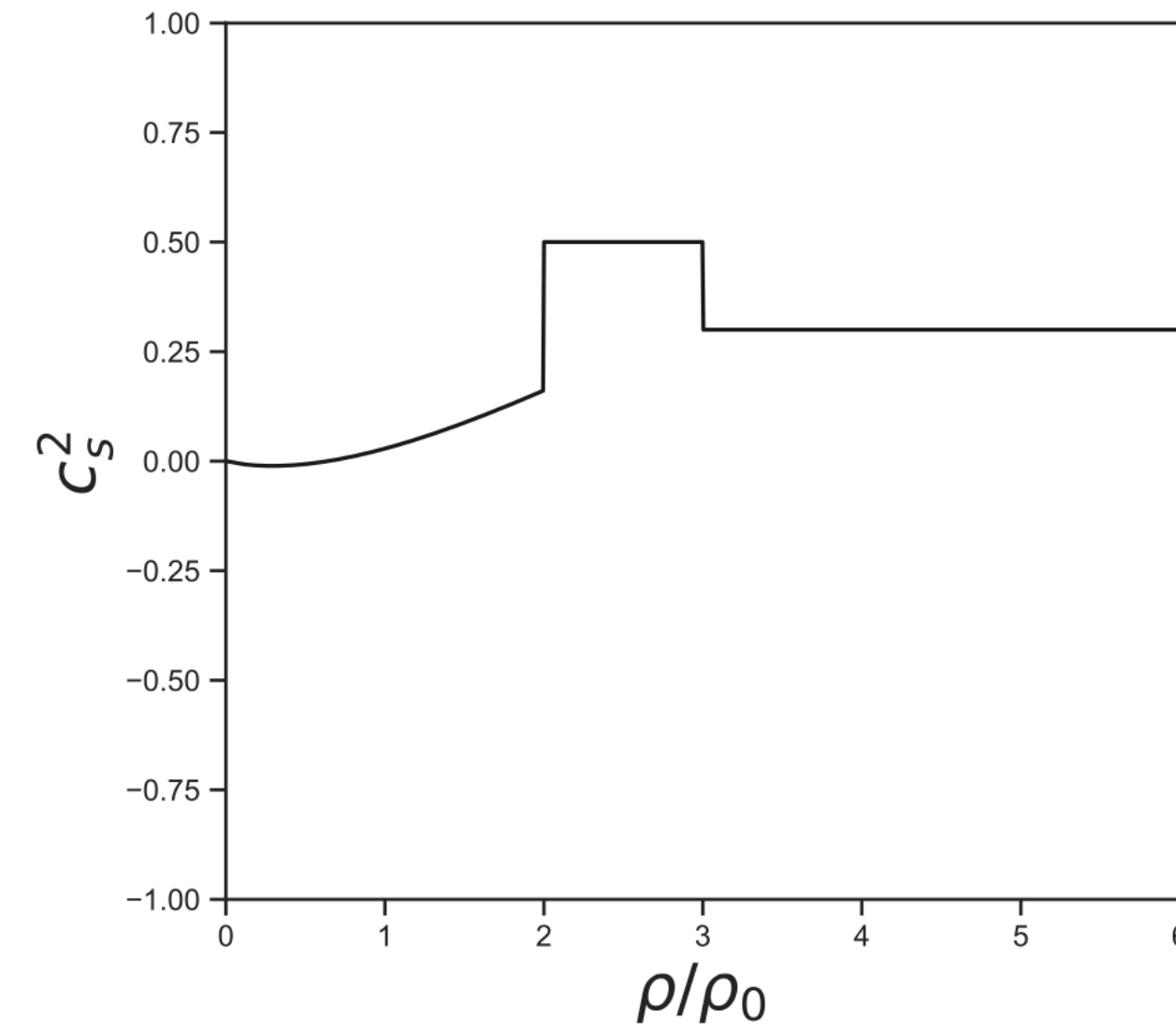
Very soft EOS at $n_B \in (2,3)n_0$ not supported in VDF+UrQMD



Bayesian analysis: piecewise parametrization of c_s^2

Piecewise parametrization of $c_s^2(n_B)$:

$$c_s^2(n_B) = \begin{cases} c_s^2(\text{Skyrme}), & n_B < n_1 = 2n_0 \\ c_1^2, & n_1 < n_B < n_2 \\ c_2^2, & n_2 < n_B < n_3 \\ \dots & \\ c_m^2, & n_m < n_B \end{cases}$$



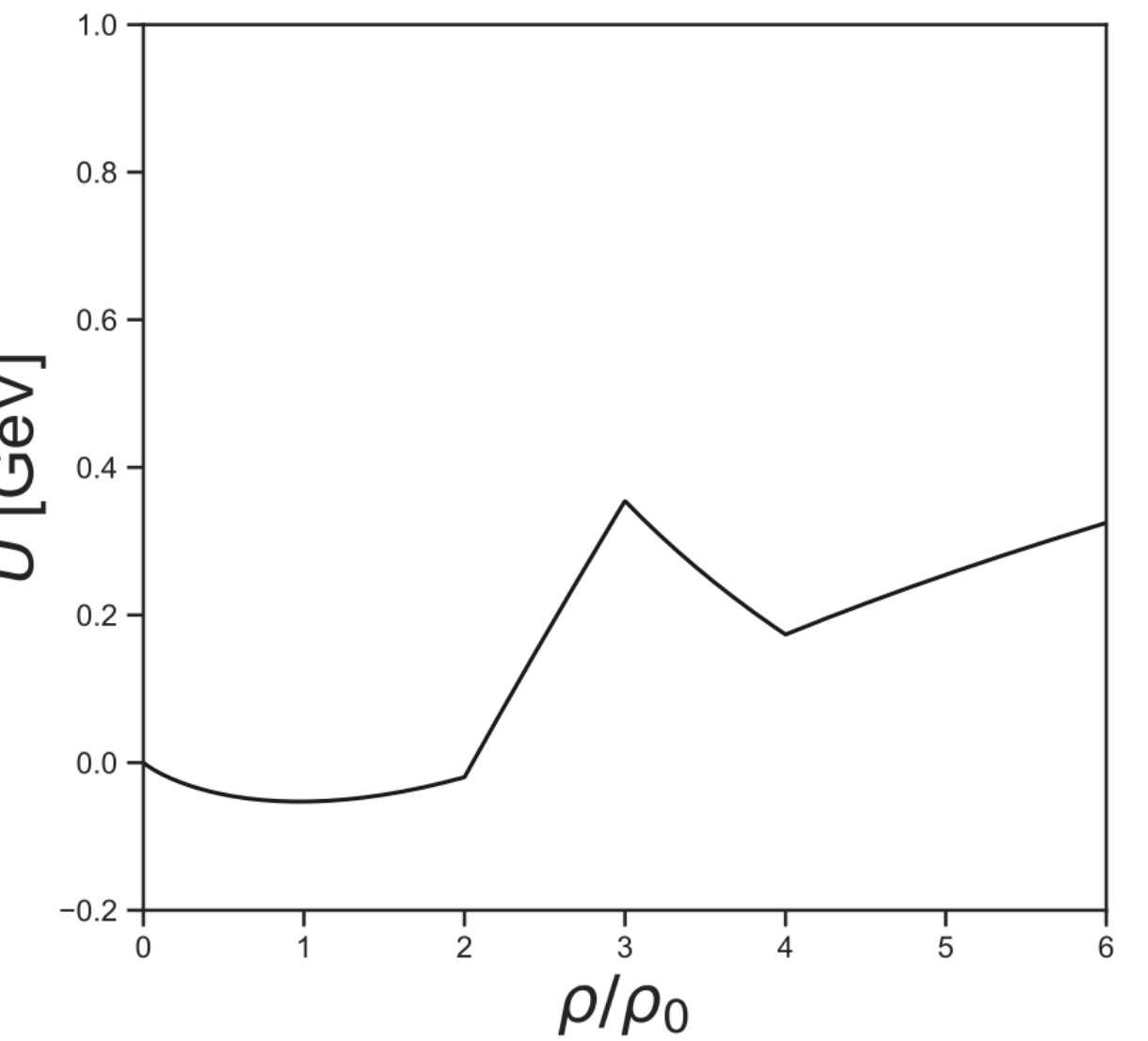
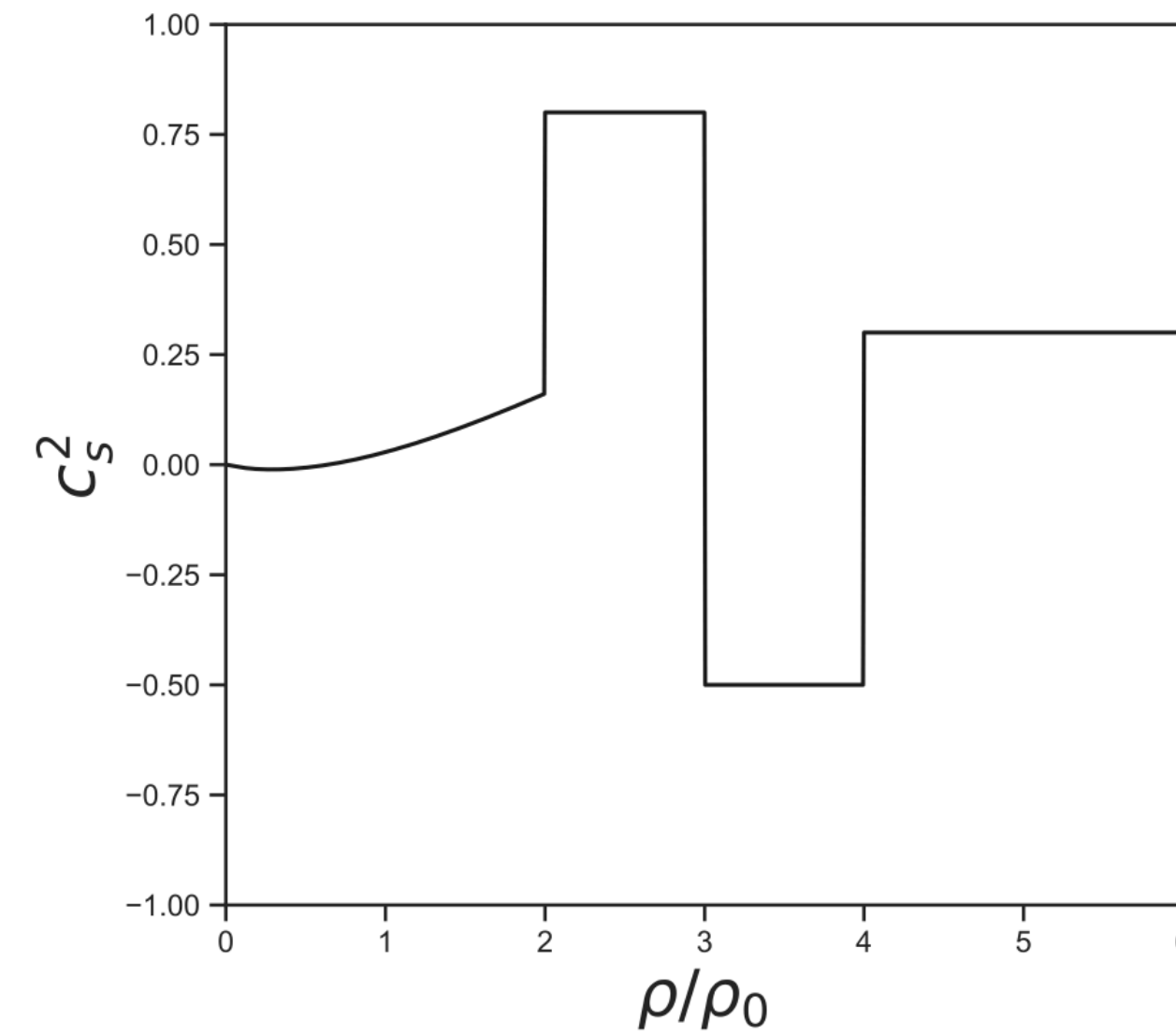
1-to-1 relation to the single-particle potential $U(n_B)$:

$$U(n_B) = \begin{cases} U_{\text{Sk}}(n_B) & n_B < n_1 = 2n_0 \\ U_1(n_B) & n_1 < n_B < n_2 \\ \dots & \\ U_k(n_B) & n_k < n_B < n_{k+1} \end{cases}$$

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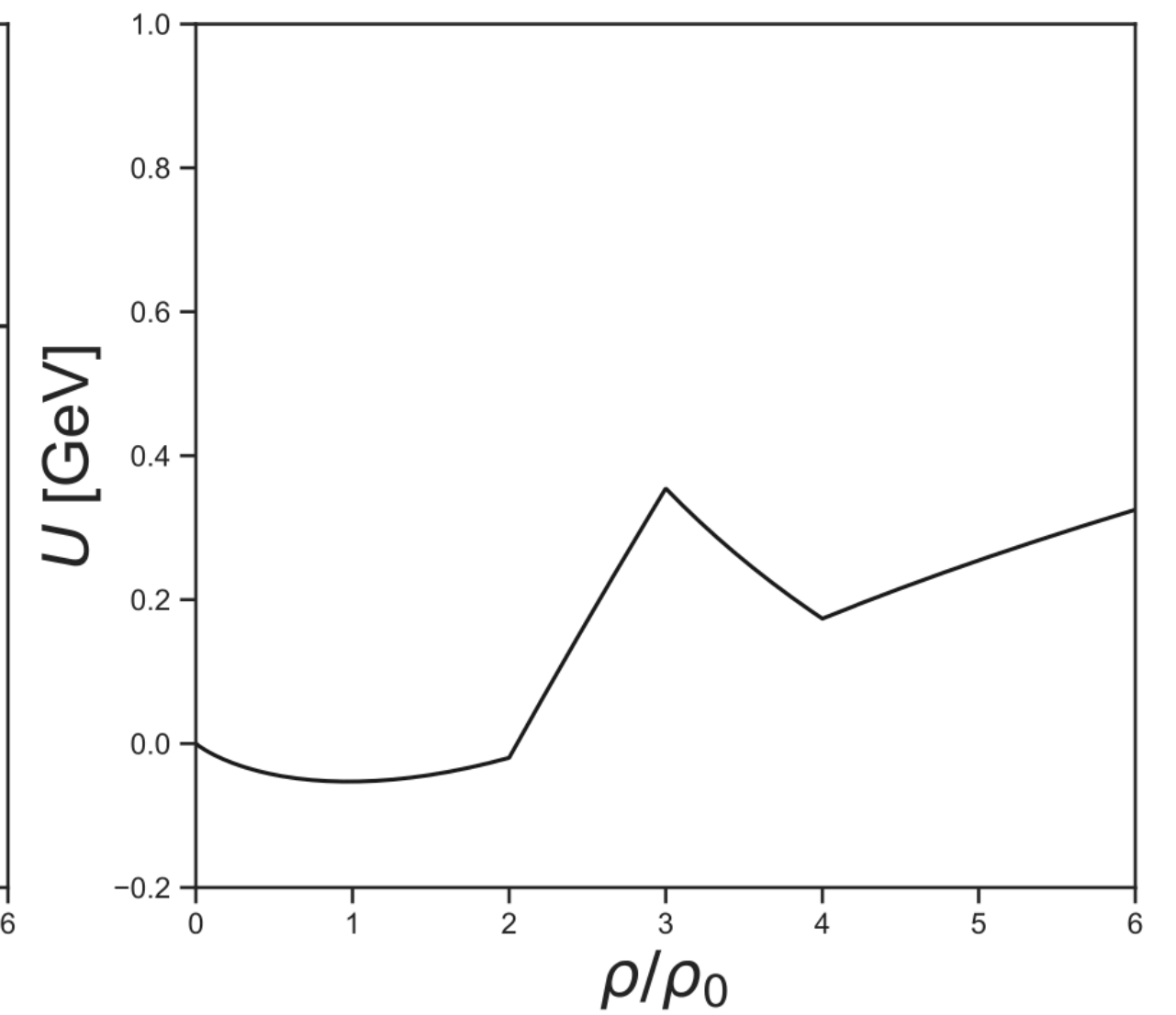
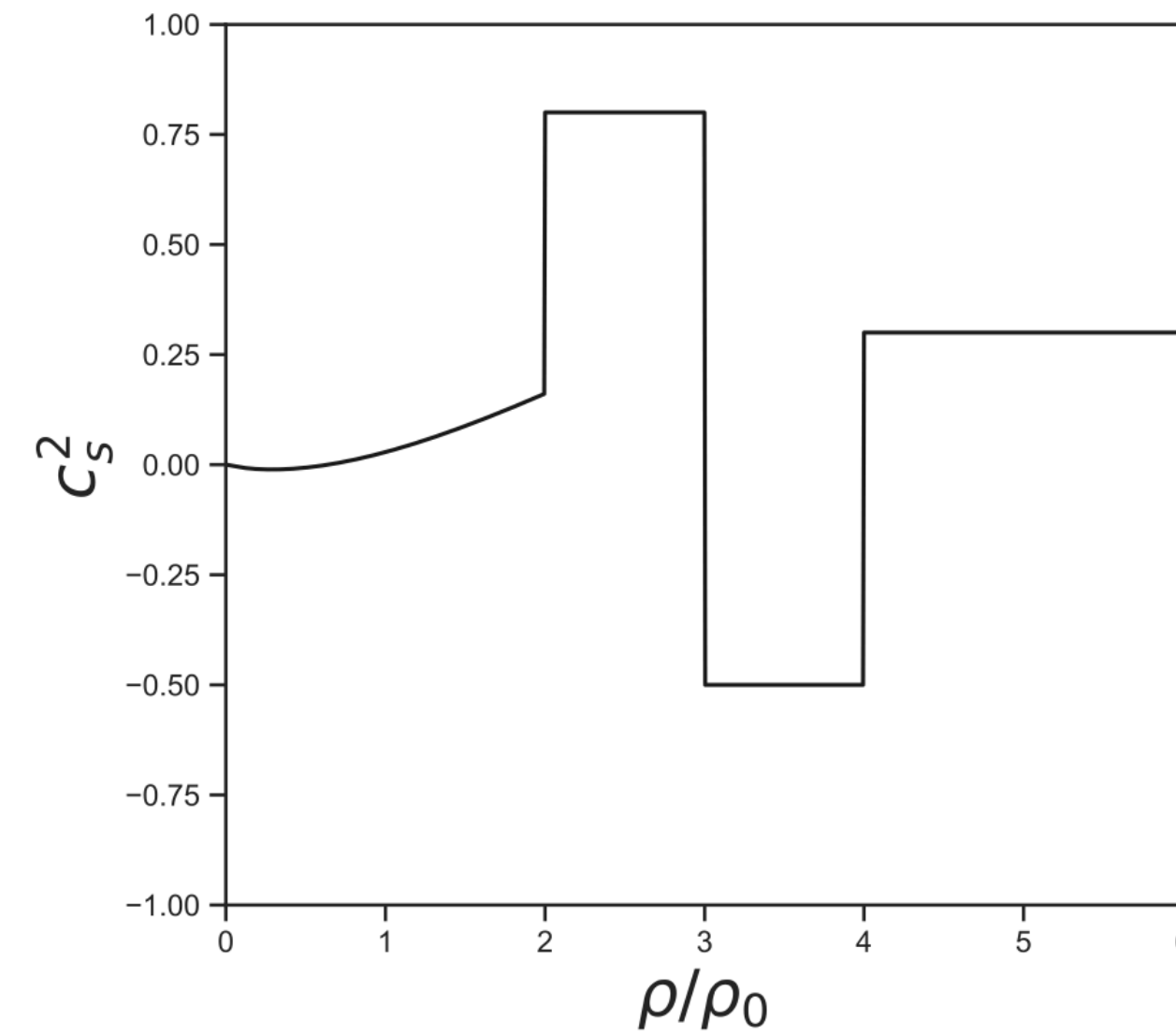
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Gradients of $U(n_B)$ enter the EOMs
= directly affect the evolution in simulations

Bayesian analysis: piecewise parametrization of c_s^2

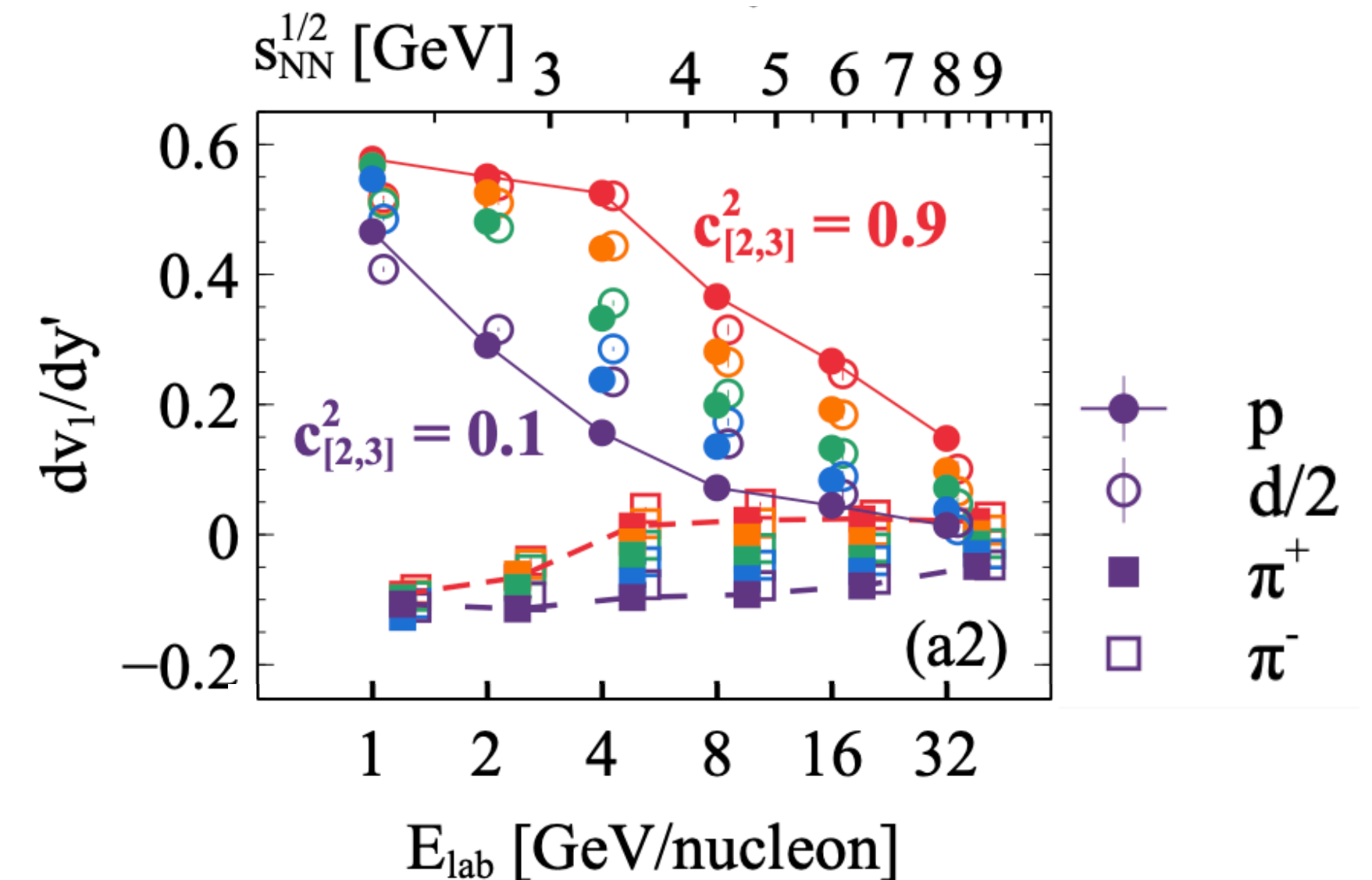
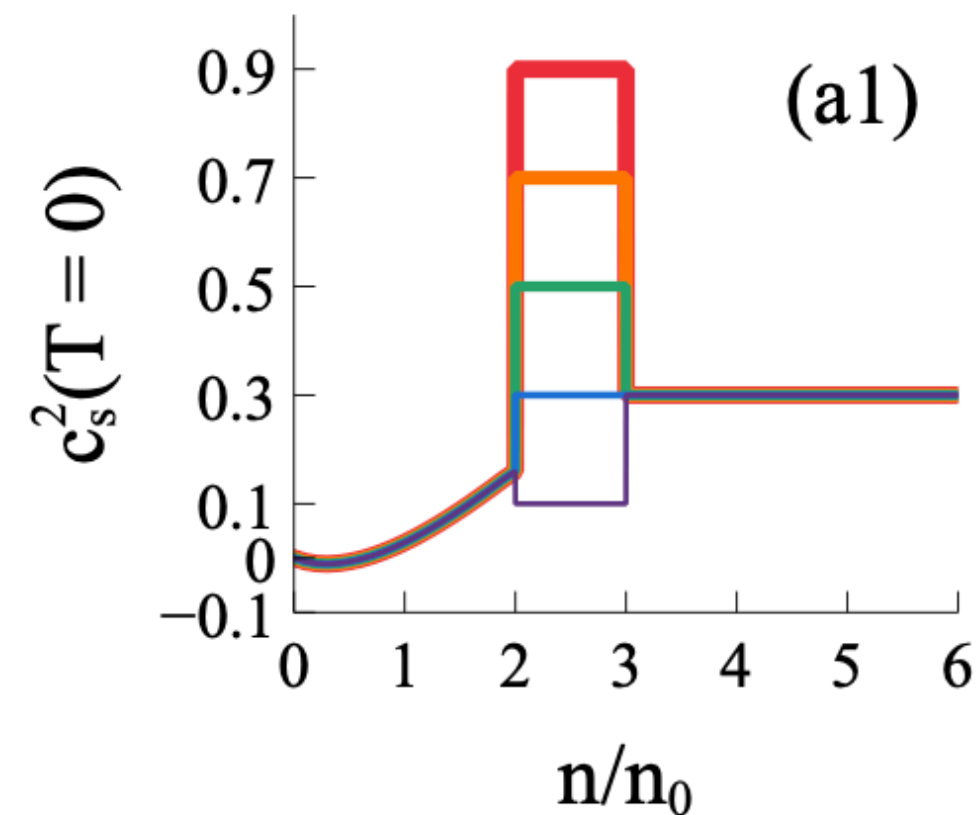
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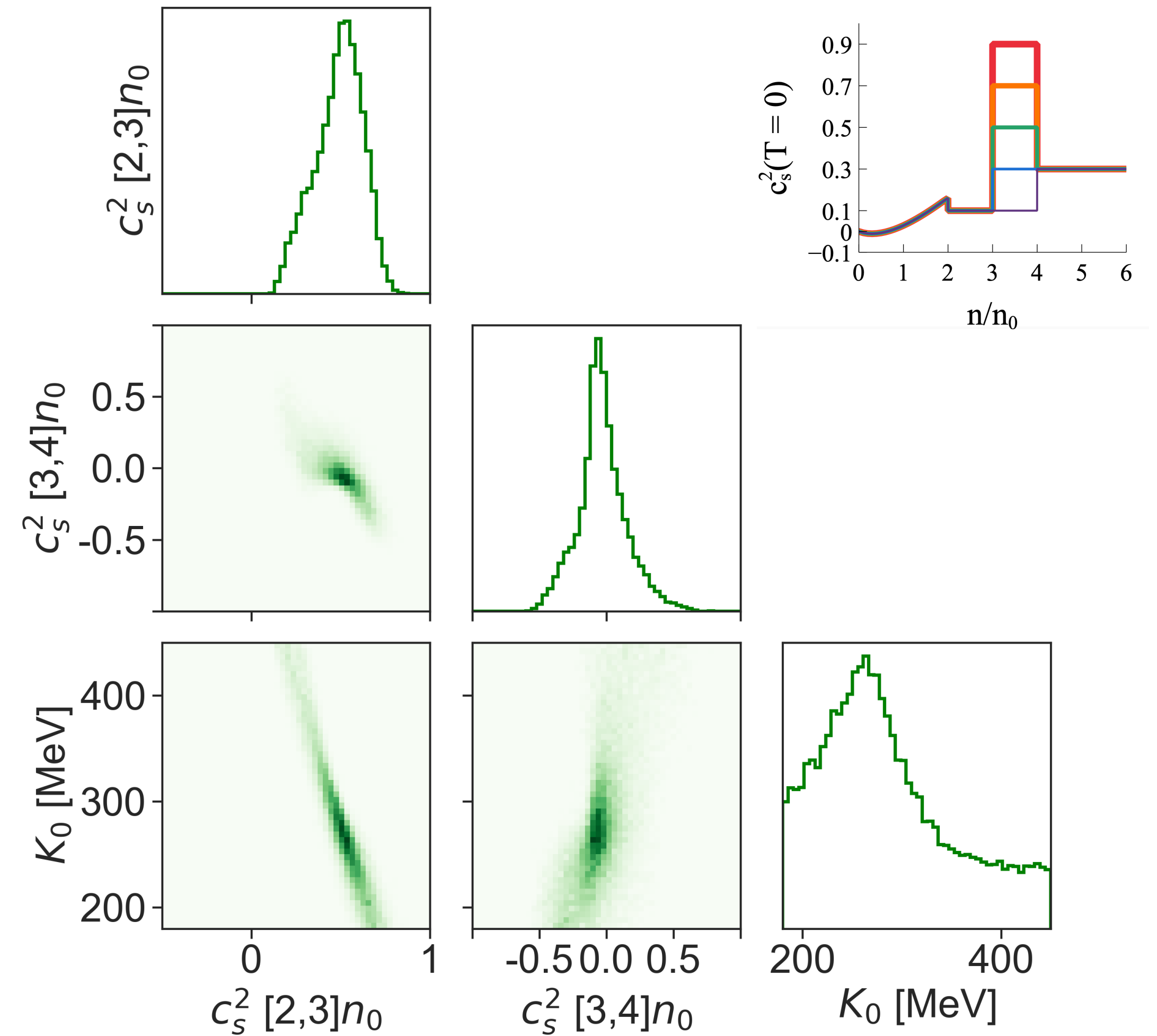
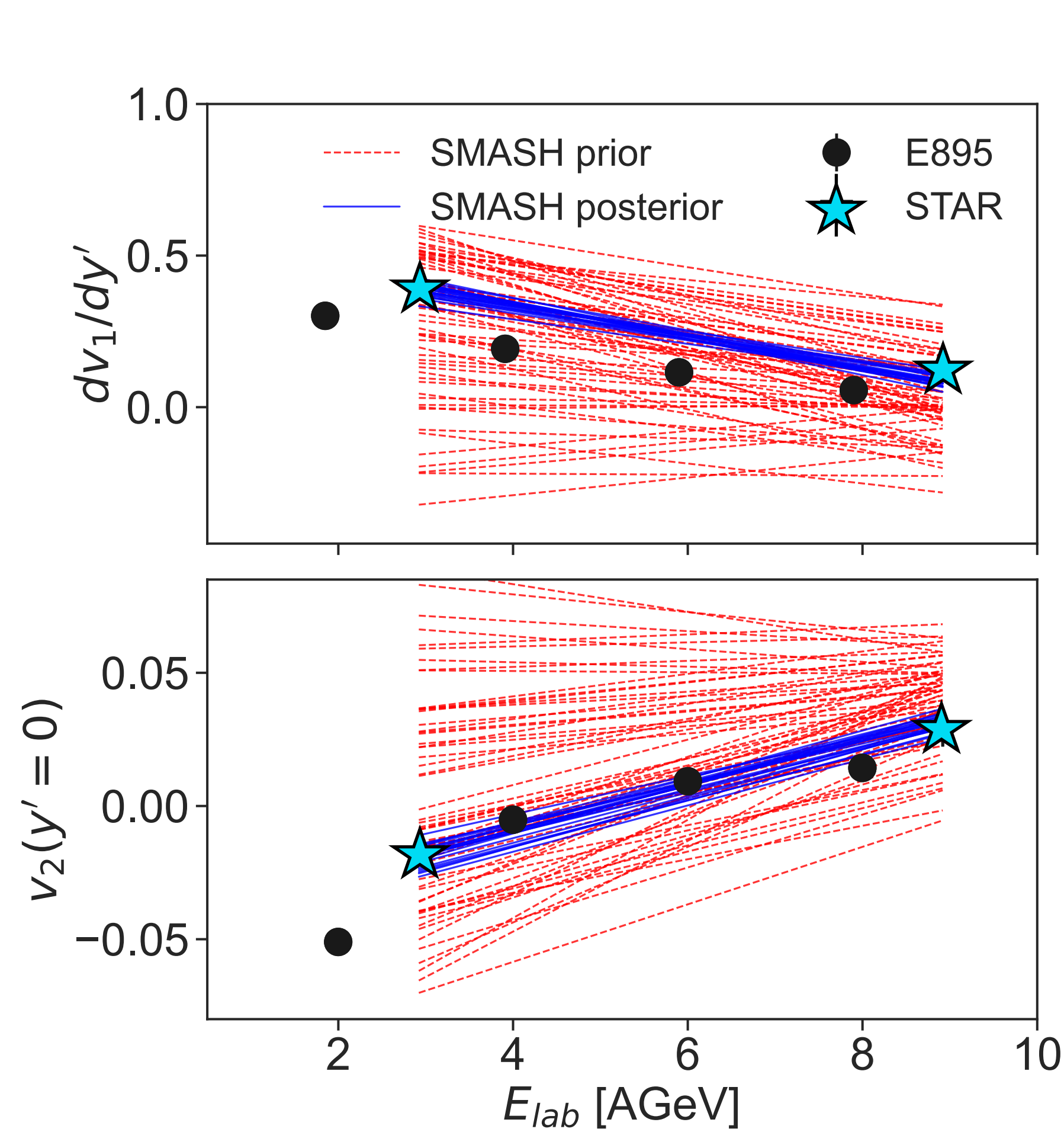
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D. Oliinychenko, A. Sorensen, V. Koch, L. McLerran,
Phys. Rev. C **108**, 3, 034908 (2023), arXiv:2208.11996

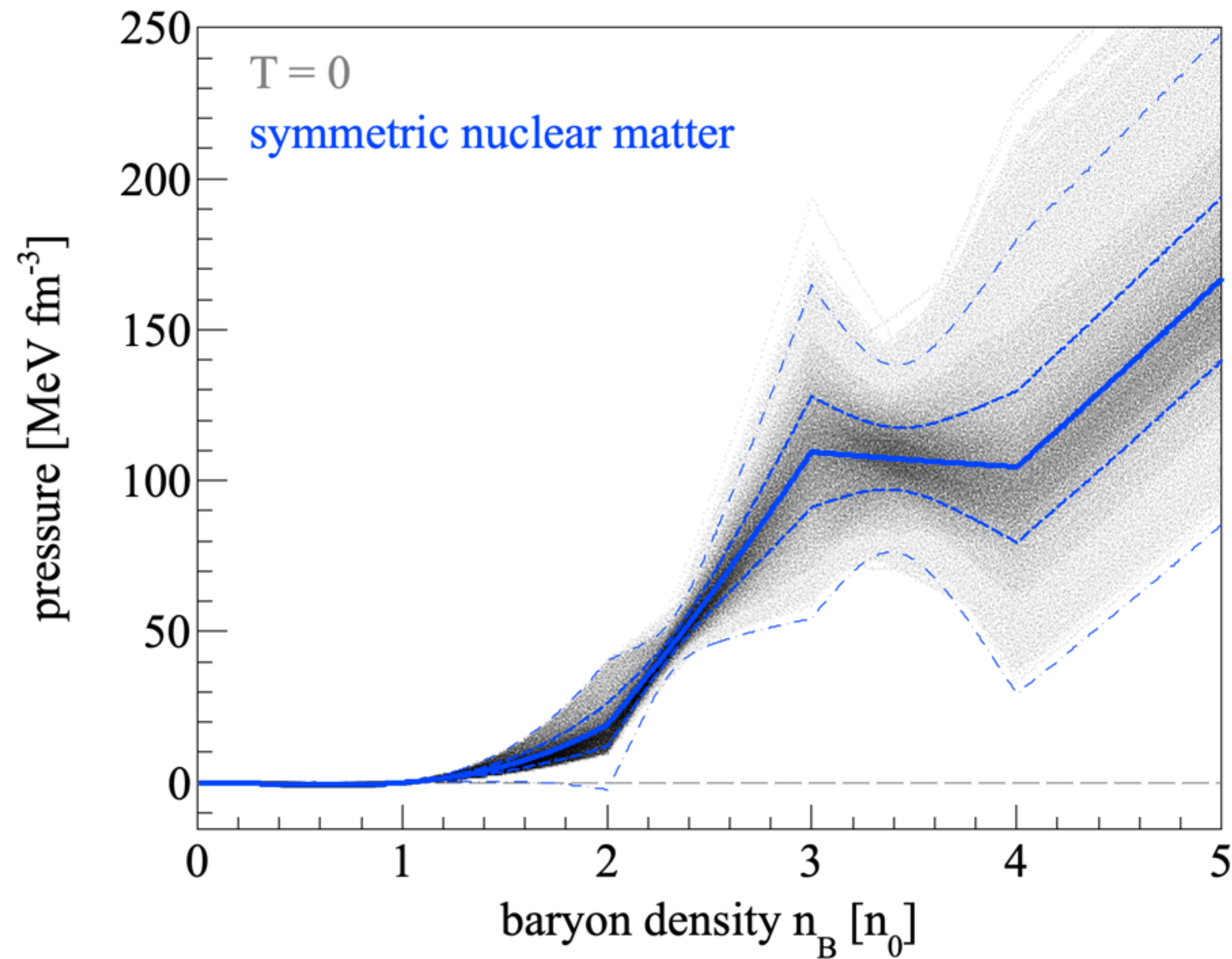
Bayesian analysis of BES flow in BUU with varying K_0 , $c_{[2,3]n_0}^2$, $c_{[3,4]n_0}^2$



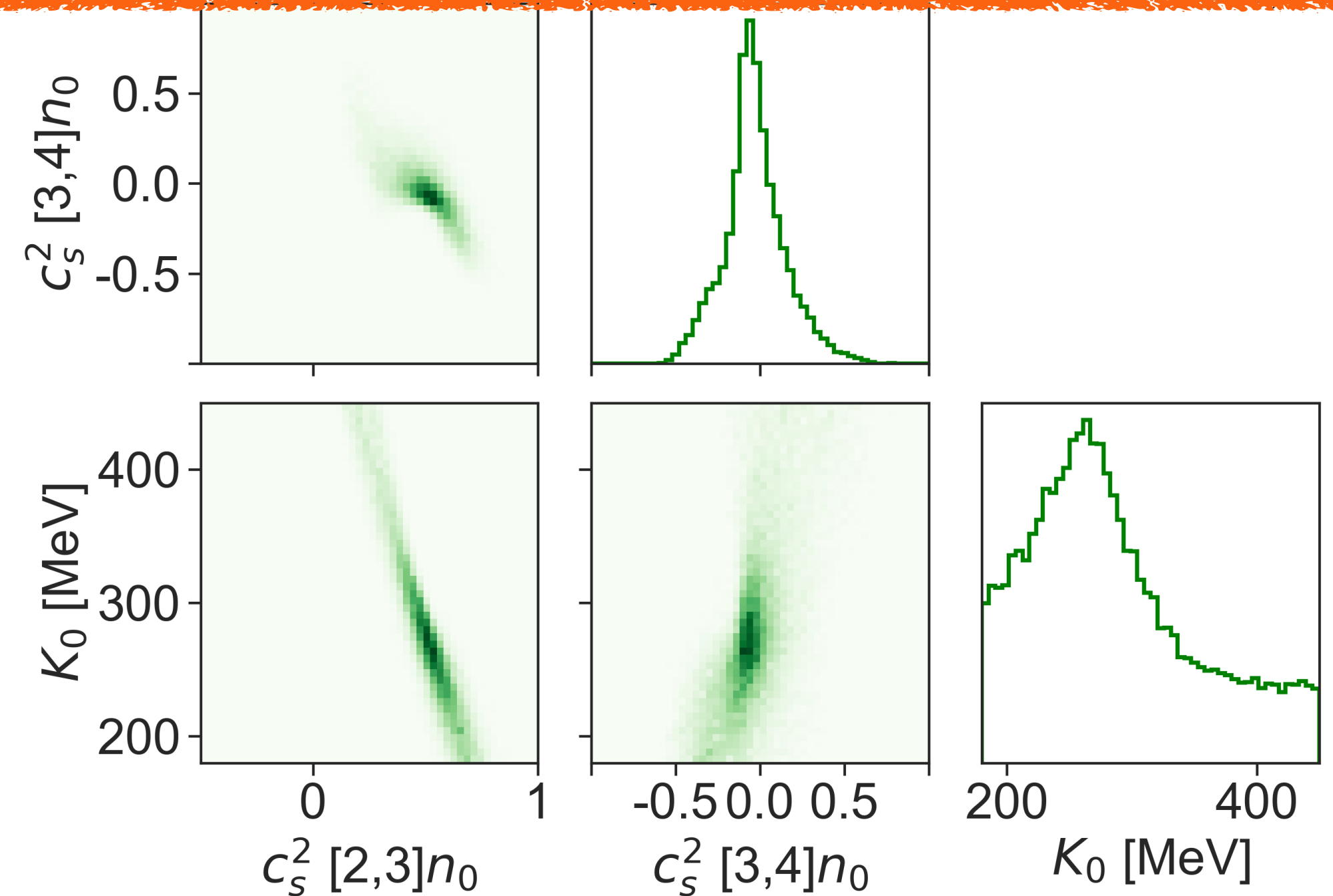
The maximum a posteriori probability (MAP) parameters are
 $K_0 = 285 \pm 67$ MeV, $c_{[2,3]n_0}^2 = 0.49 \pm 0.13$, $c_{[3,4]n_0}^2 = -0.03 \pm 0.15$

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The constrained EOS is very stiff at $n_B \in (2,3)n_0$ and very soft at $n_B \in (3,4)n_0$

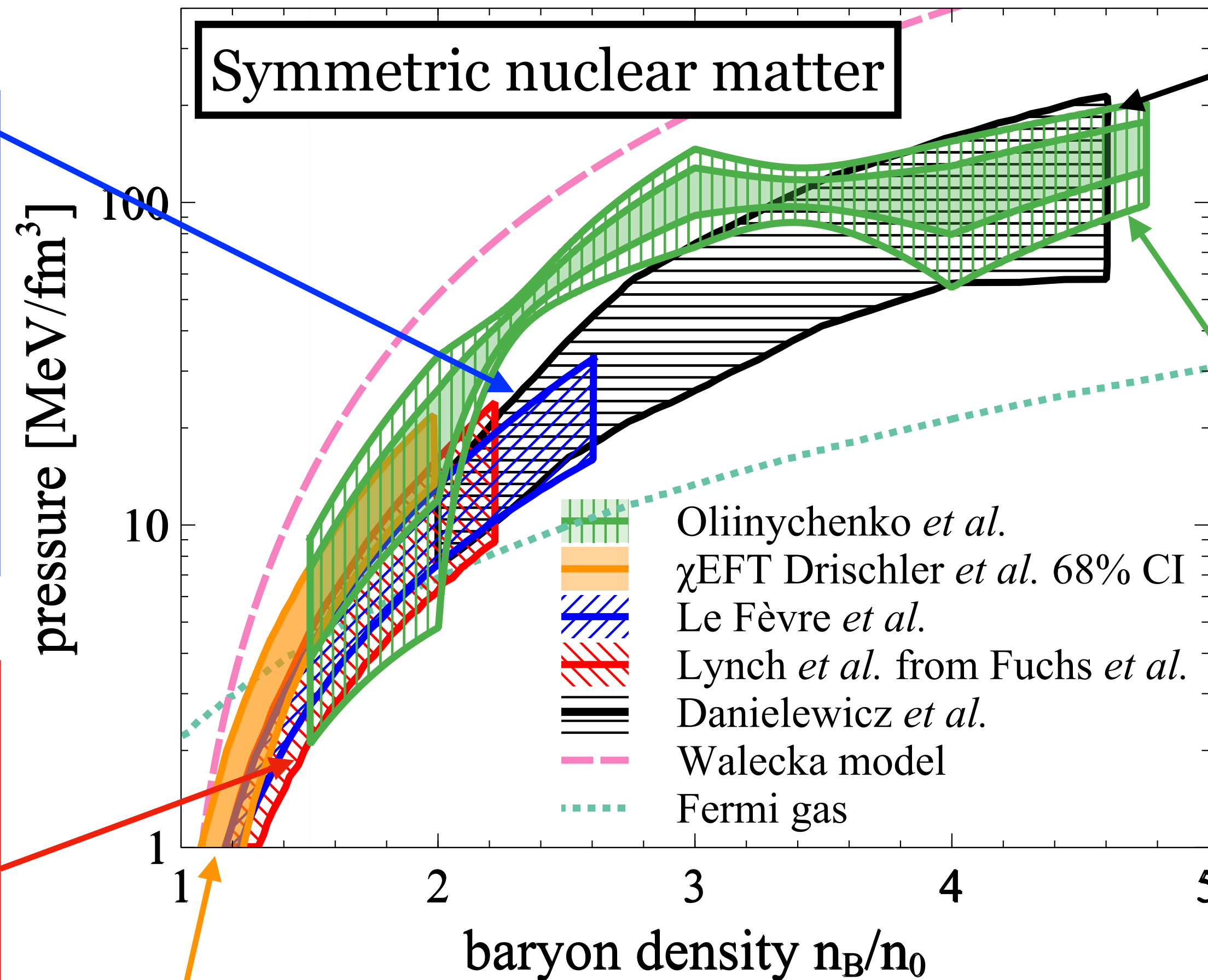


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D. Oliinychenko, **A. Sorensen**, V. Koch, L. McLerran,
 Phys. Rev. C **108**, 3, 034908 (2023), arXiv:2208.11996

EOS of symmetric nuclear matter: selected (*few*) results

Symmetric nuclear matter



197Au+197Au @ 0.15–10 GeV/u
 ($\sqrt{s_{NN}} = 1.95 - 4.72$ GeV)
 observables: proton flow
 (Plastic Ball, EOS, E877, E895)
 model used: **pBUU** w/ nucleons, Δ ,
 $N^*(1440)$, pions;
 EOS parametrized by K_0 ;
 momentum dependence
 P. Danielewicz, R. Lacey, W. G. Lynch,
 Science **298**,1592–1596 (2002)

197Au+197Au @ 2.9–9 GeV/u
 ($\sqrt{s_{NN}} = 3 - 4.5$ GeV)
 observables: proton flow (STAR)
 model used: SMASH w/ over 120 hadronic
 species, including deuterons;
 relativistic EOS parametrized independently in
 different density regions;
NO momentum dependence
 D. Oliinychenko, **A. Sorensen**, V. Koch,
 L. McLerran, Phys. Rev. C **108**, 3, 034908
 (2023), arXiv:2208.11996

197Au+197Au @ 0.4–1.5 GeV/u
 ($\sqrt{s_{NN}} = 2.07 - 2.52$ GeV)
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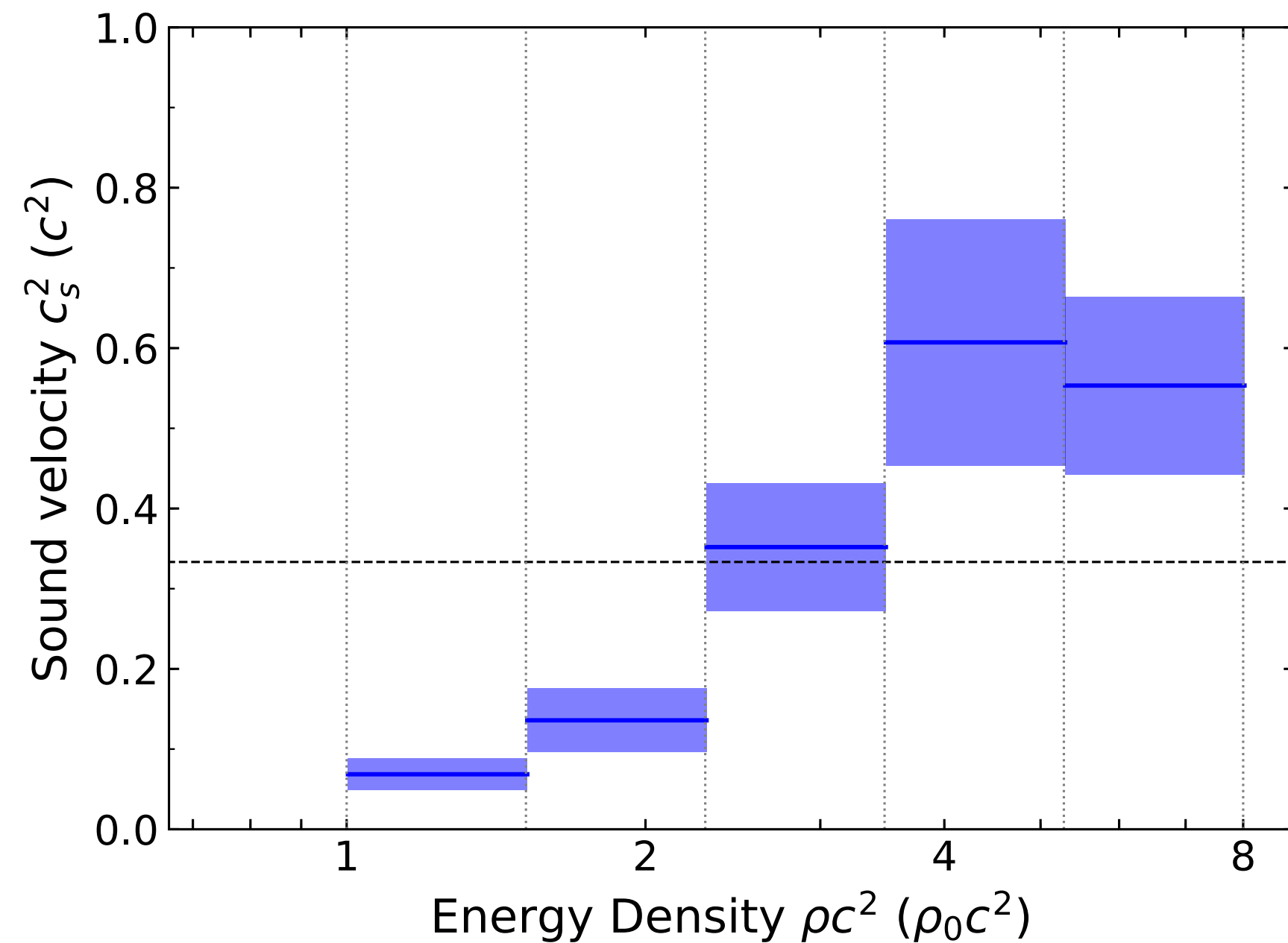
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A. Sorensen *et al.*, Prog. Part. Nucl. Phys. **134**, 104080 (2024)
 arXiv:2301.13253

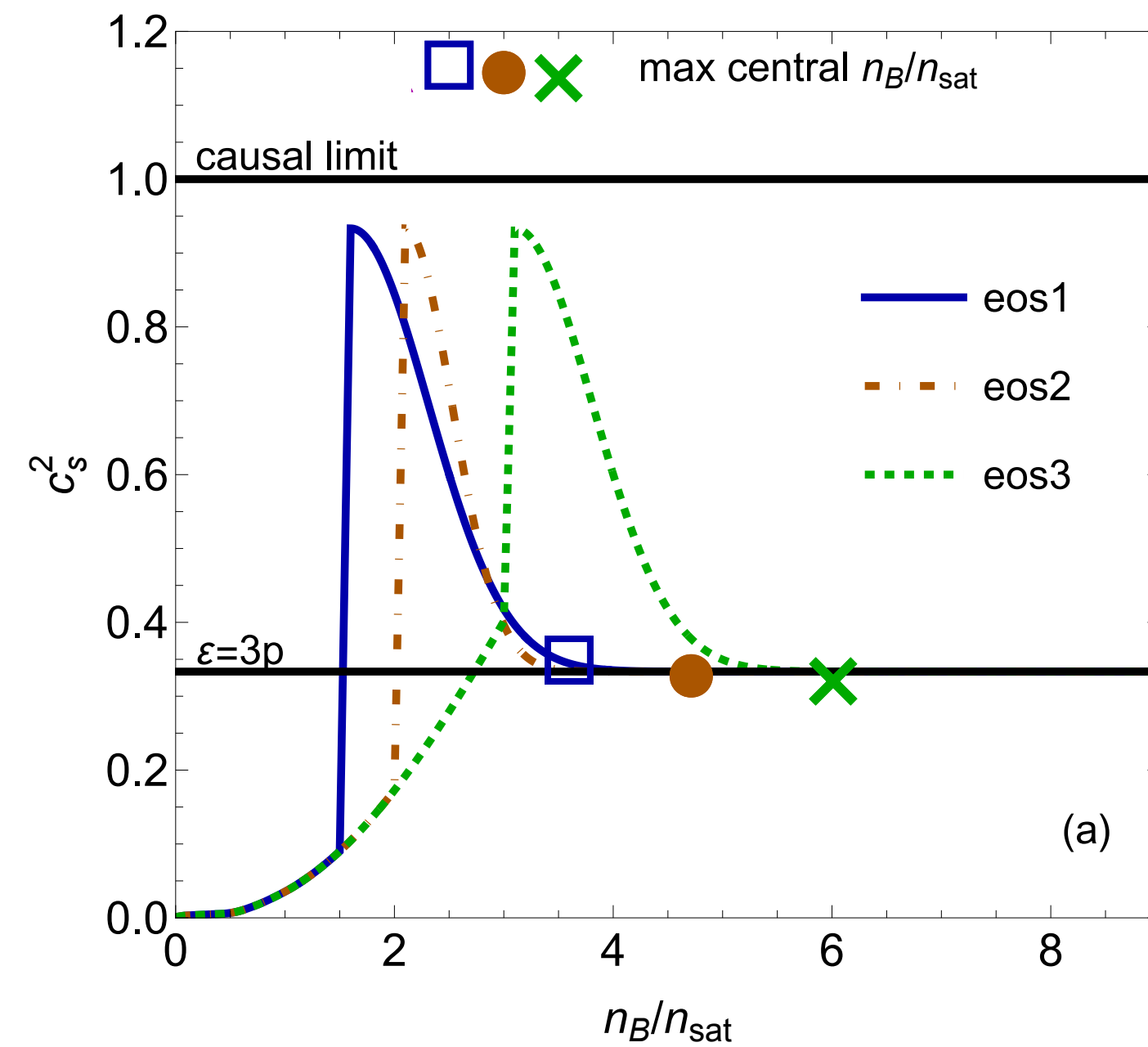
χ EFT
 C. Drischler *et al.*, Phys. Rev. C **102** 5, 054315 (2020)
 arXiv:2004.07805

Intriguing results from analyses of astrophysical observations

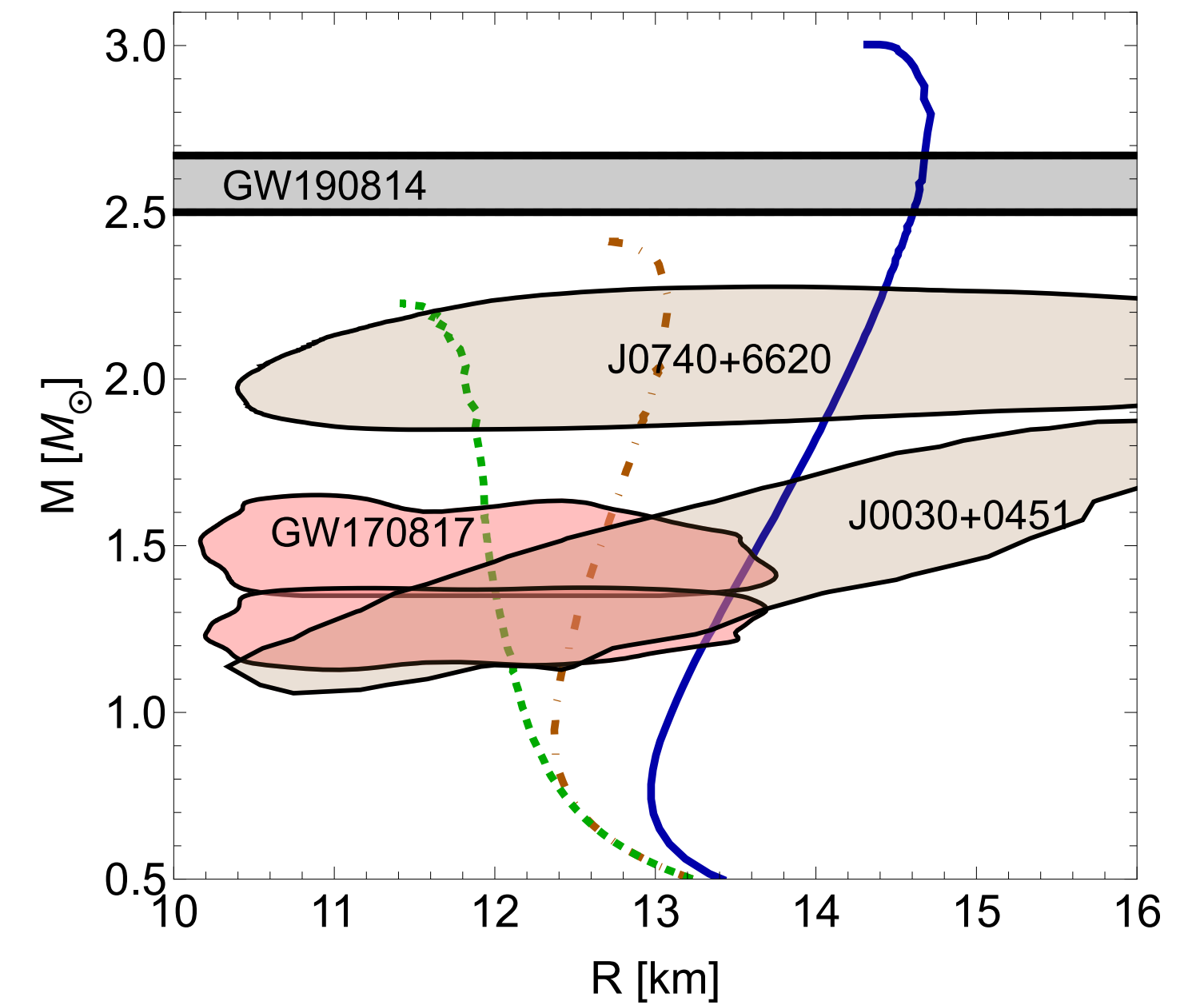
Recent astrophysical measurements suggest NS EOS may have a nontrivial density-dependence



Y. Fujimoto, K. Fukushima, K. Murase,
Phys. Rev. D **101**, 5, 054016 (2020), arXiv:1903.03400



N. Yao, **A. Sorensen**, V. Dexheimer, J. Noronha-Hostler, arXiv:2311.18819
H. Tan, T. Dore, V. Dexheimer, J. Noronha-Hostler, N. Yunes, Phys. Rev. D **105** 2, 023018 (2022)



Which of the allowed NS EOSs are compatible with heavy-ion collision measurements?

EOS from neutron stars to heavy-ion collisions

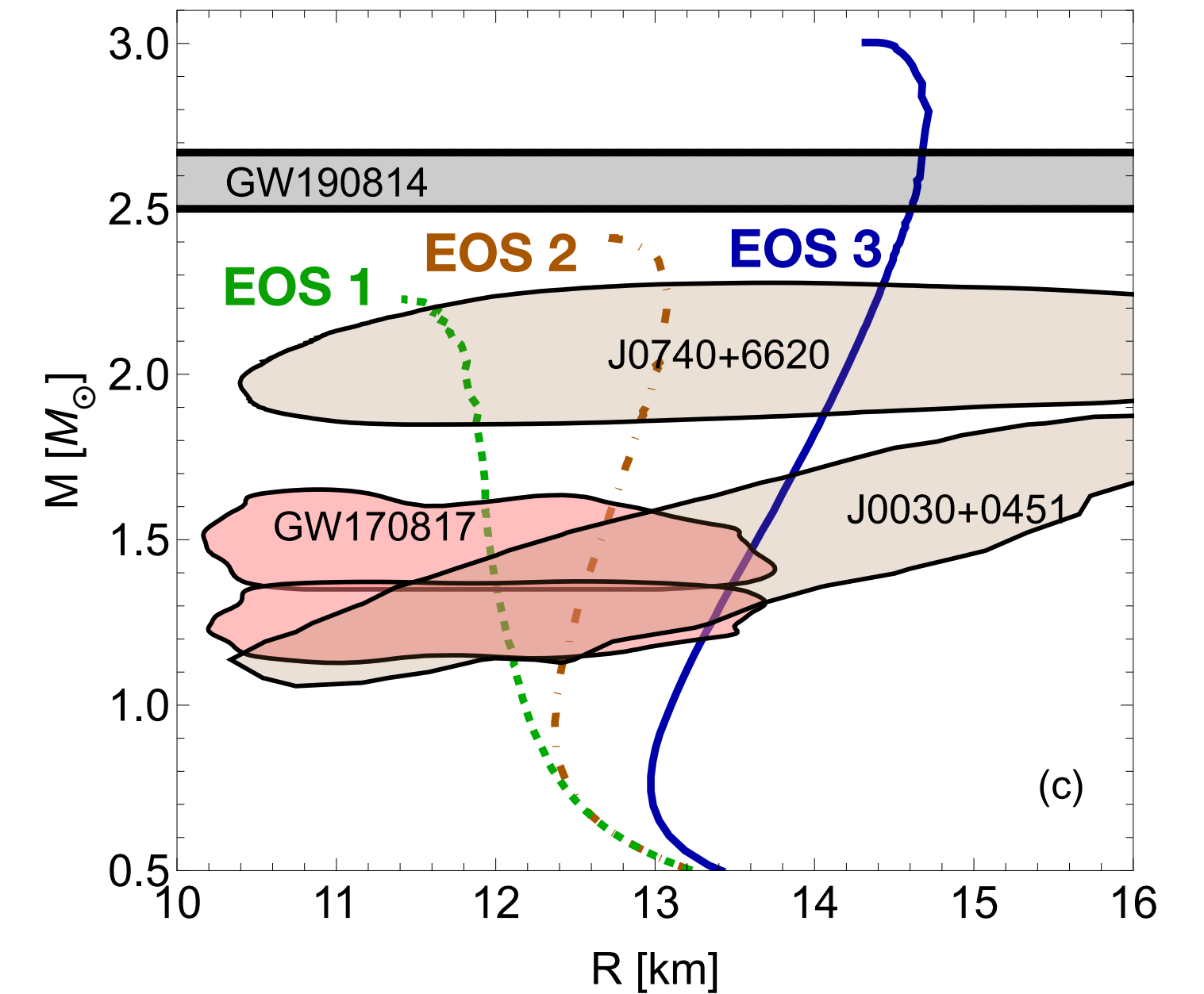
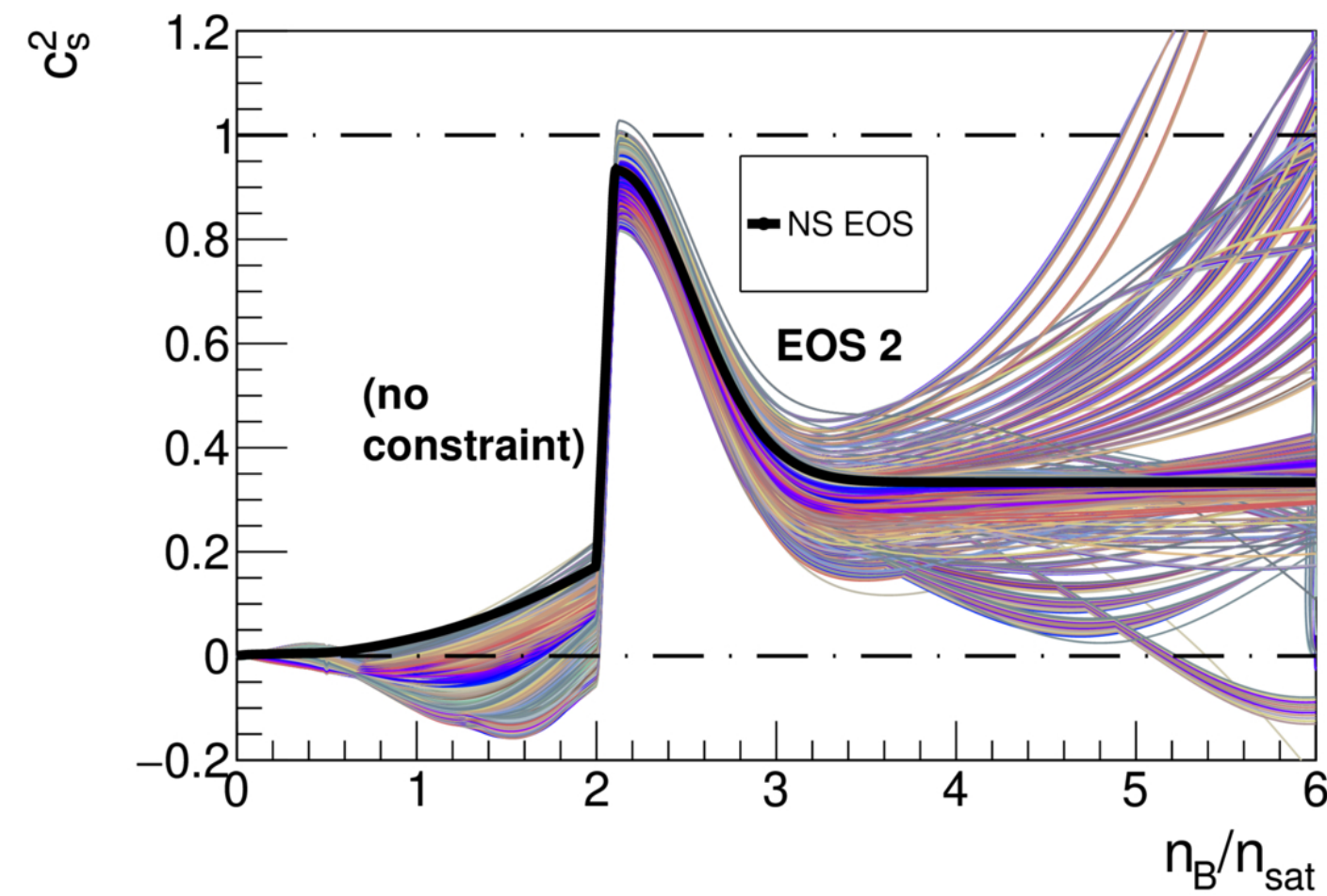
Use symmetry energy parameters exploring the allowed parameter space:

N. Yao, A. Sorensen, V. Dexheimer, J. Noronha-Hostler, arXiv:2311.18819

$$\mathcal{E}_{\text{HIC}} = \mathcal{E}_{\text{NS}} - n_B \left[E_{\text{sym}} + \frac{L_{\text{sym}}}{3} \left(\frac{n_B}{n_0} - 1 \right) + \frac{K_{\text{sym}}}{18} \left(\frac{n_B}{n_0} - 1 \right)^2 + \frac{J_{\text{sym}}}{162} \left(\frac{n_B}{n_0} - 1 \right)^3 \right] \delta^2, \quad \delta \equiv \frac{n_N - n_P}{n_N + n_P}$$

\approx symmetry energy $S(n_B)$

Coefficient	Range	Step size
$E_{\text{sym,sat}}$	27 – 40	1 MeV
$L_{\text{sym,sat}}$	30 – 130	10 MeV
$K_{\text{sym,sat}}$	-220 – 180	50 MeV
$J_{\text{sym,sat}}$	-200 – 800	100 MeV



H. Tan, T. Dore, V. Dexheimer, J. Noronha-Hostler, N. Yunes, Phys. Rev. D **105** 2, 023018 (2022)

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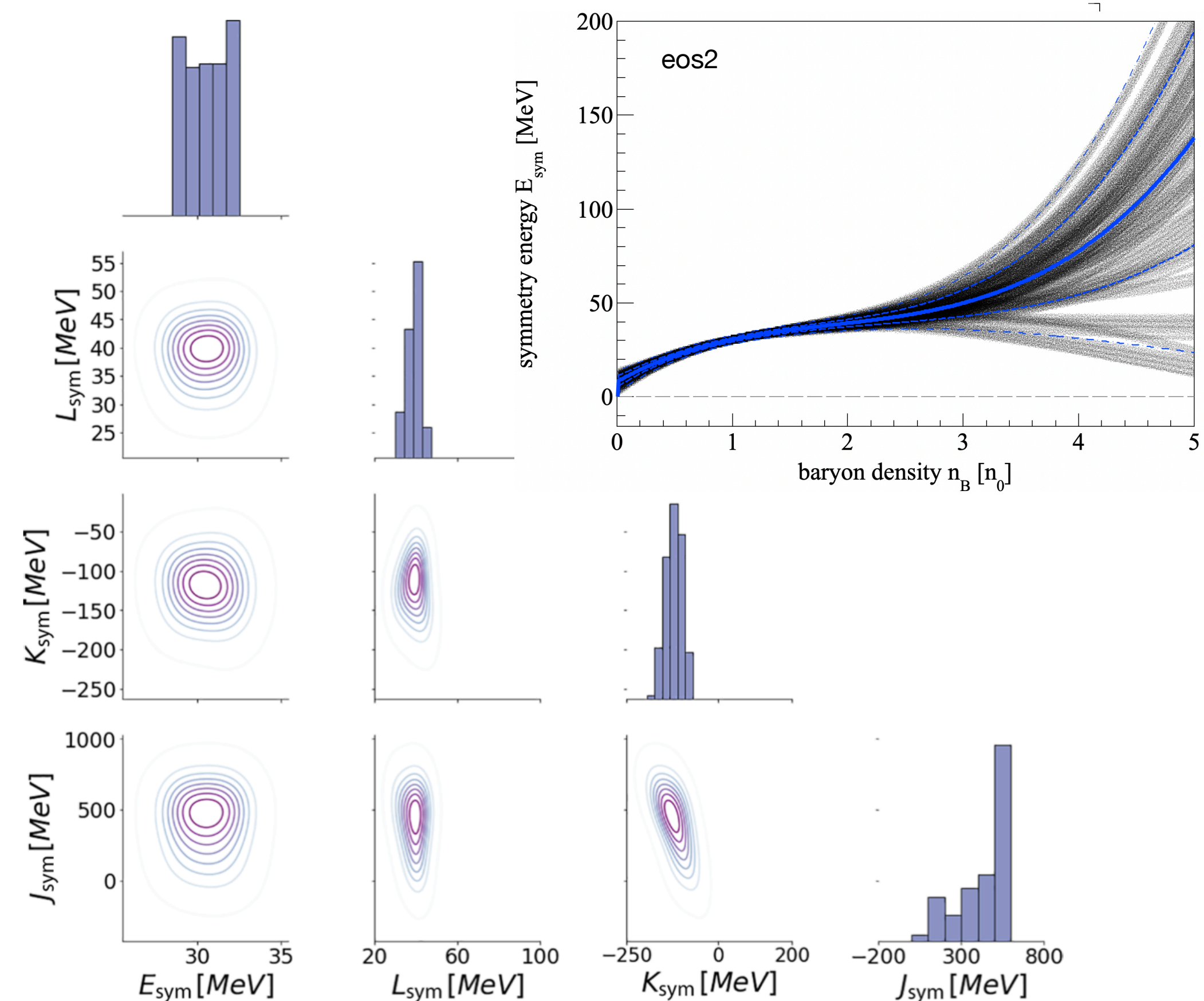
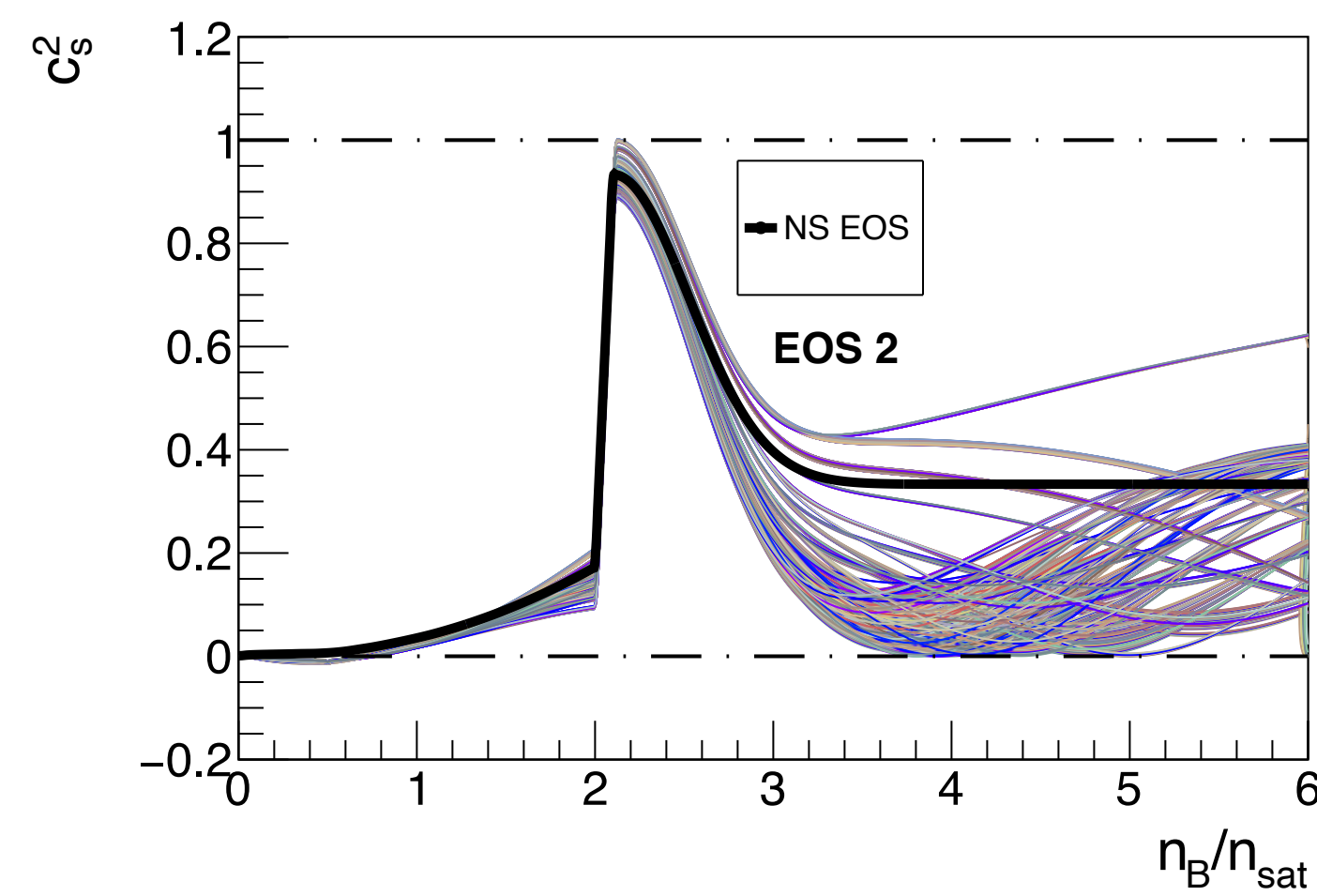
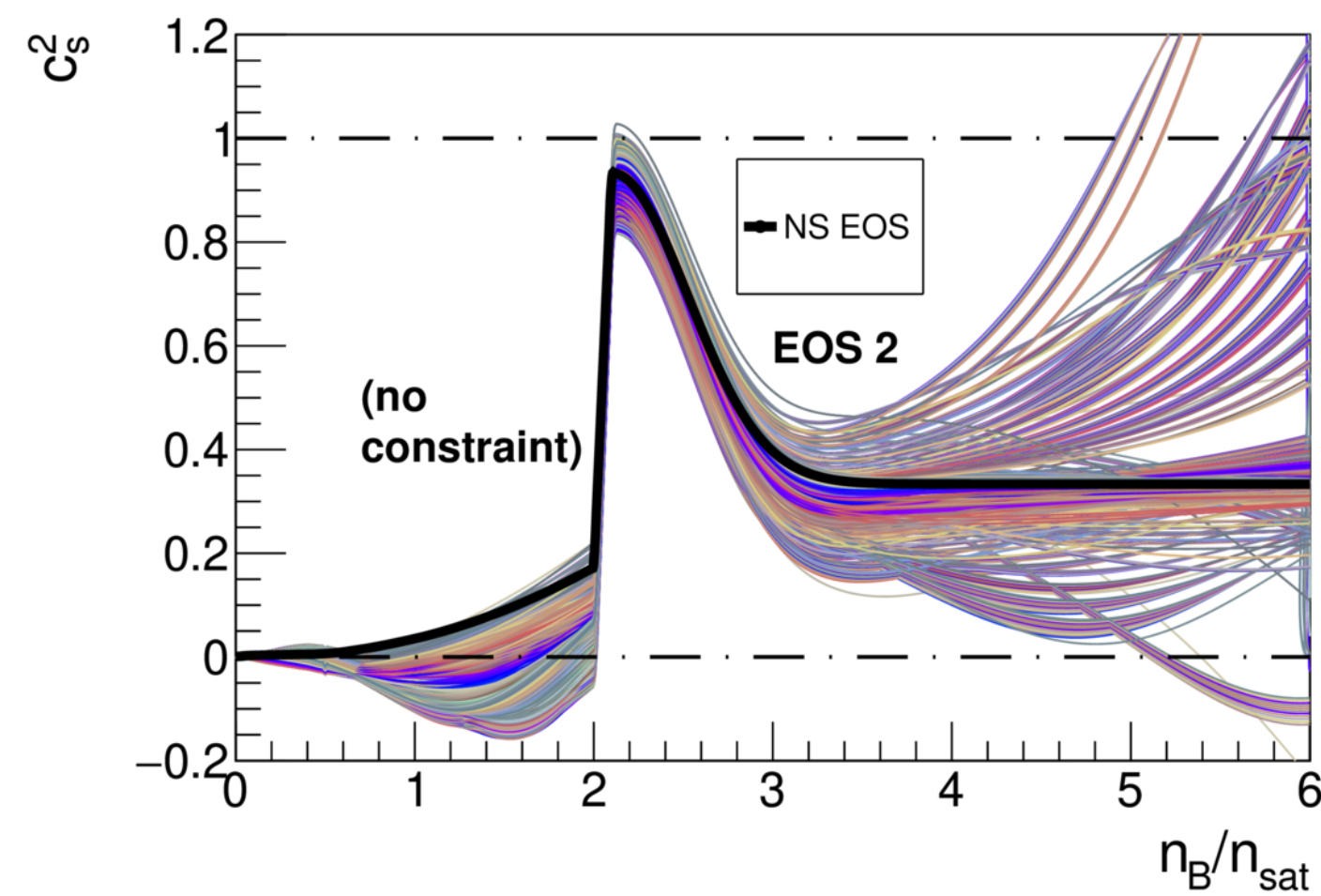
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N. Yao, A. Sorensen, V. Dexheimer, J. Noronha-Hostler, arXiv:2311.18819

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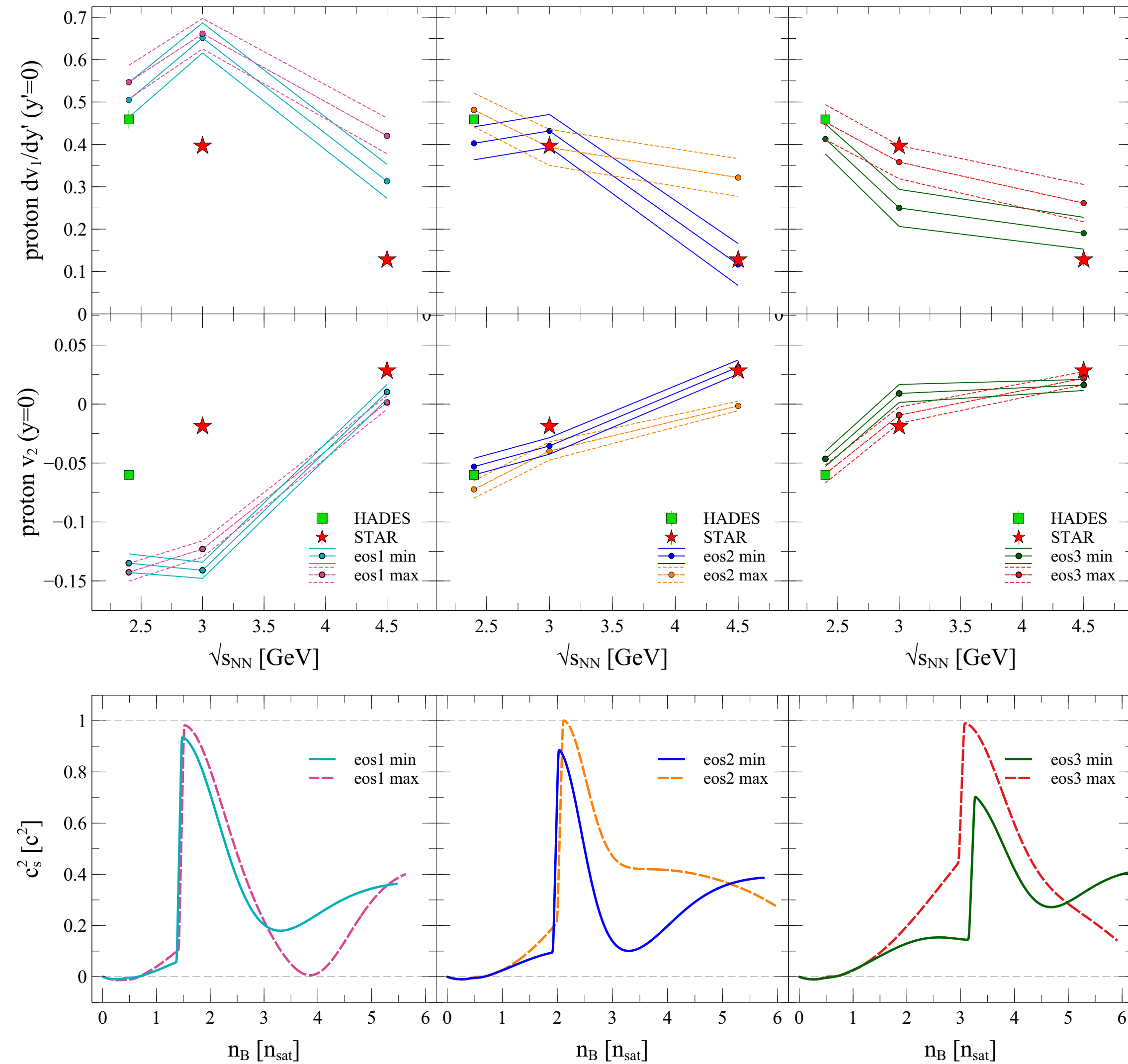
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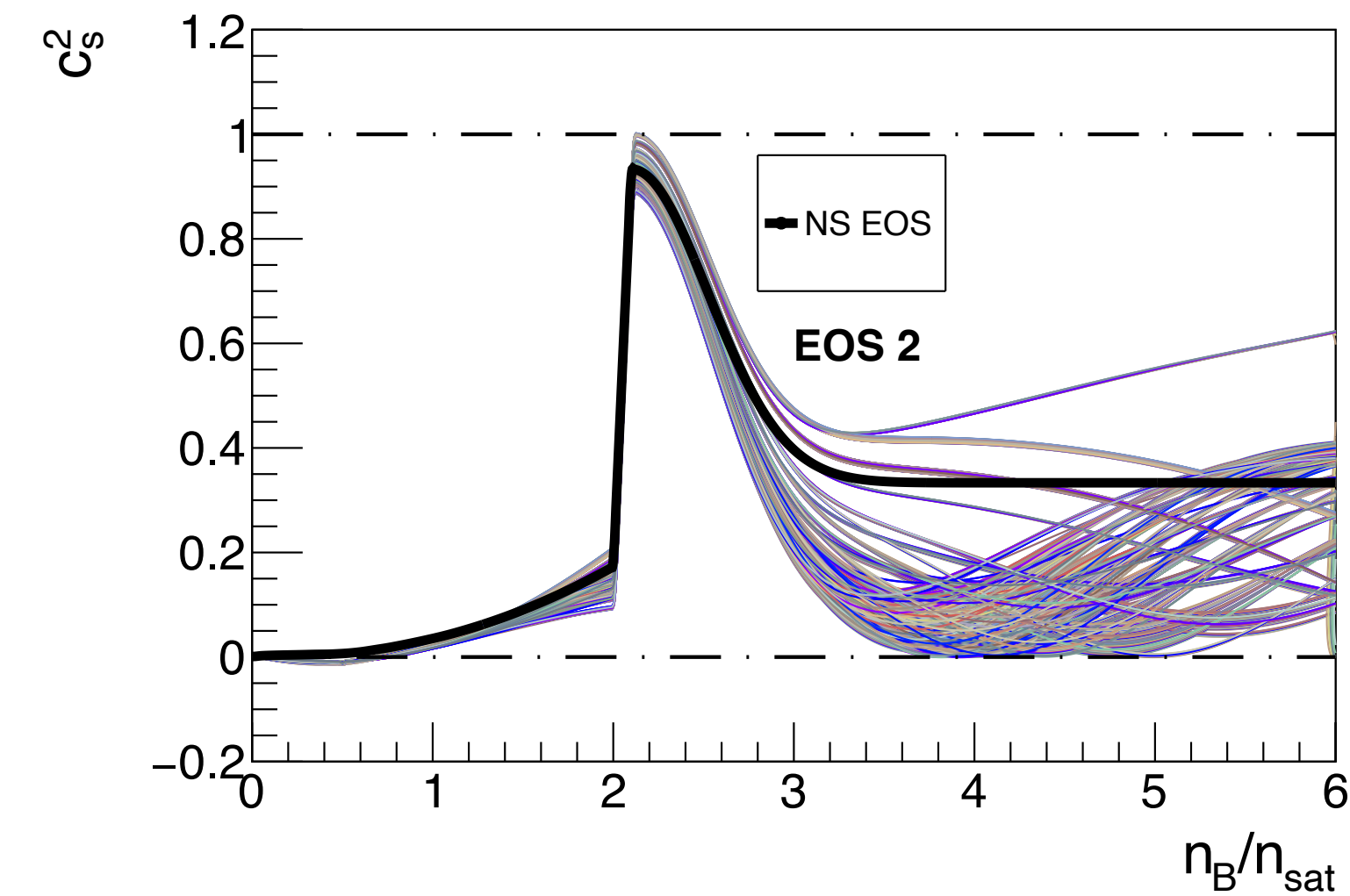
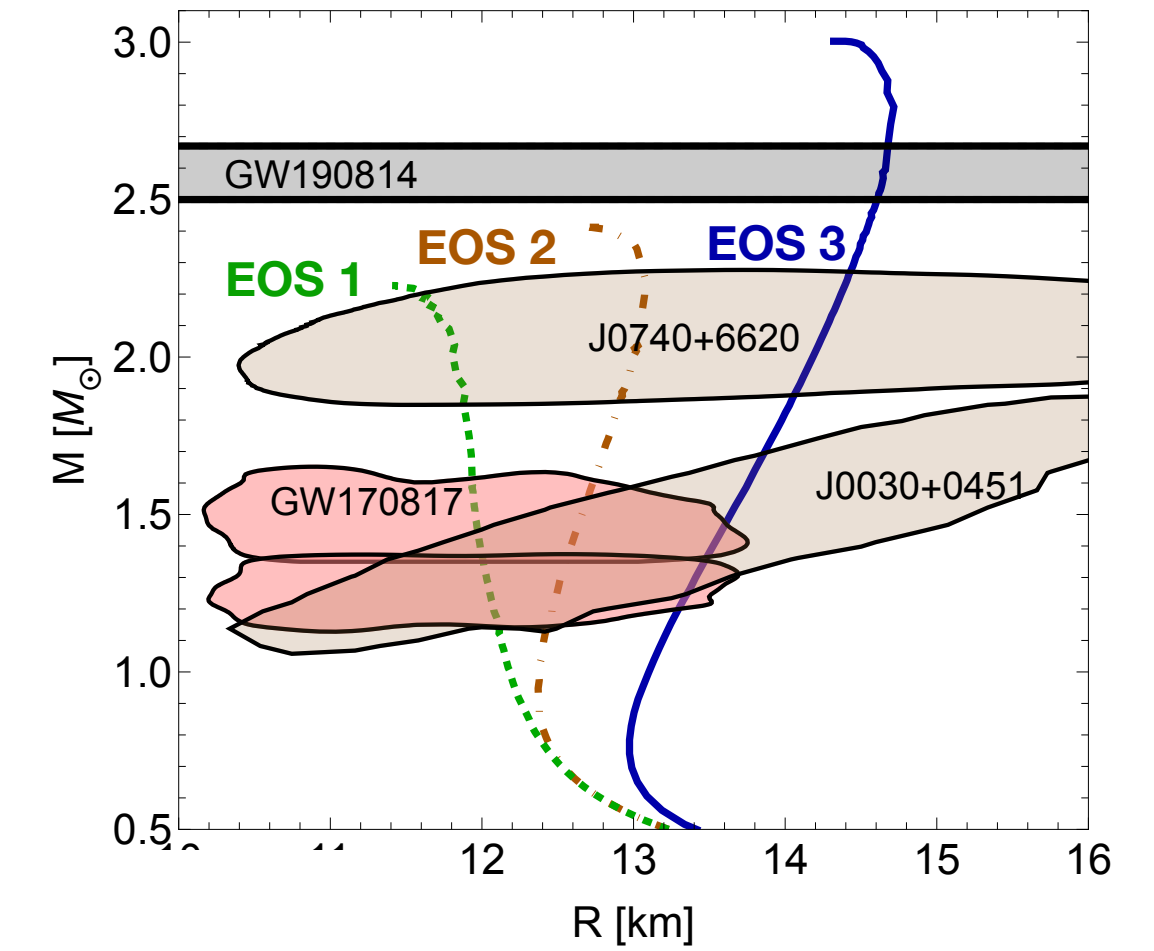


Enforcing nuclear matter properties, causality, and stability above $n_B \approx n_0$ leads to constraints on E_{sym} , L_{sym} , K_{sym} , J_{sym}

“Minimal” and “maximal” EOSs from each family tested against heavy-ion measurements:



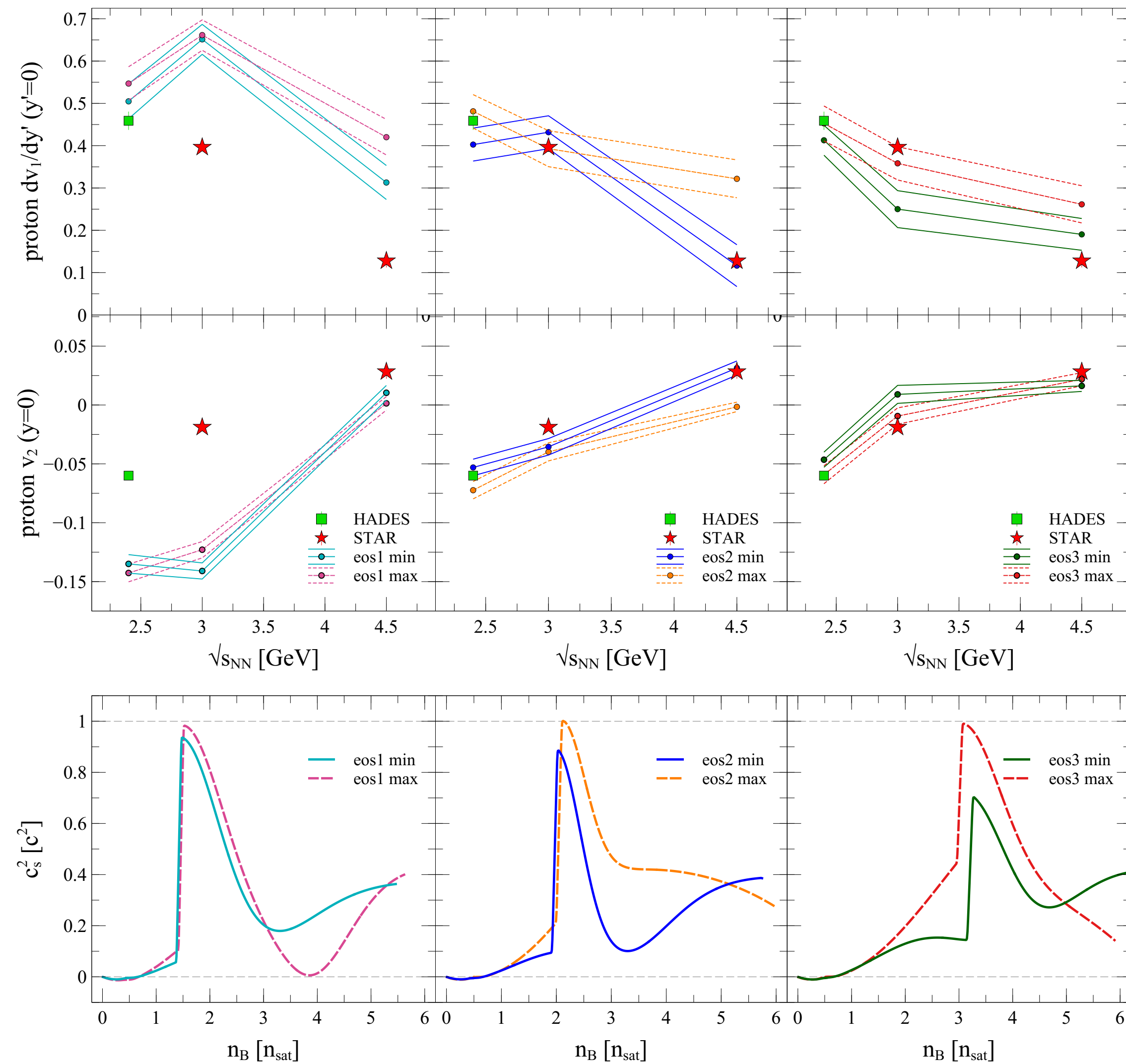
EOS	$n_{\text{sat}} [\text{fm}^{-3}]$	$B [\text{MeV}]$	$K_0 [\text{MeV}]$	$c_s^2(n_B = n_{\text{sat}})$
eos1 min	0.175	-14.6	200.5	0.024
eos1 max	0.171	-17.8	325.9	0.039
eos2 min	0.167	-14.6	206.7	0.025
eos2 max	0.161	-16.9	214.8	0.026
eos3 min	0.153	-14.8	220.2	0.027
eos3 max	0.162	-16.5	201.7	0.024



N. Yao, A. Sorensen, V. Dexheimer,
J. Noronha-Hostler, arXiv:2311.18819

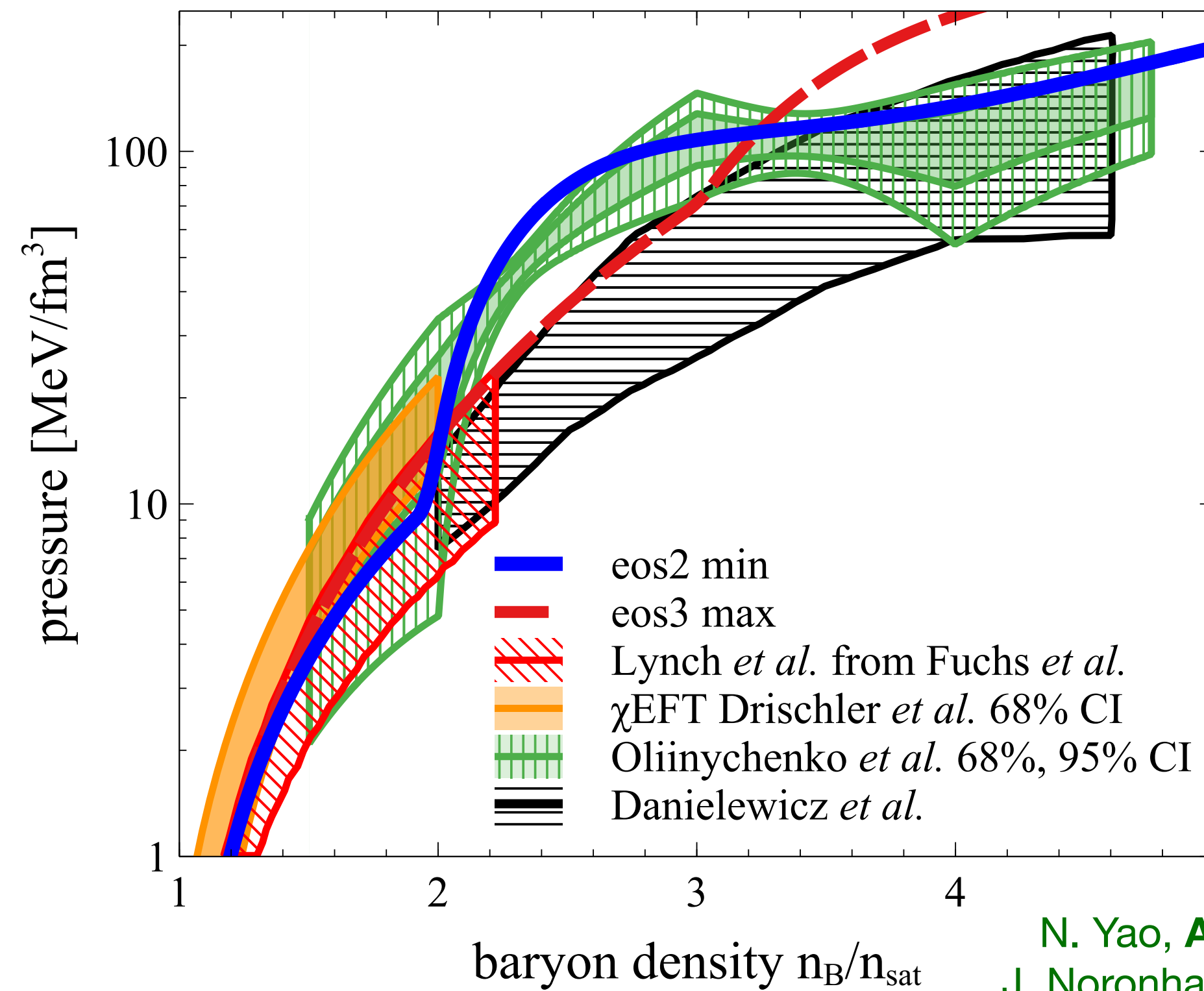
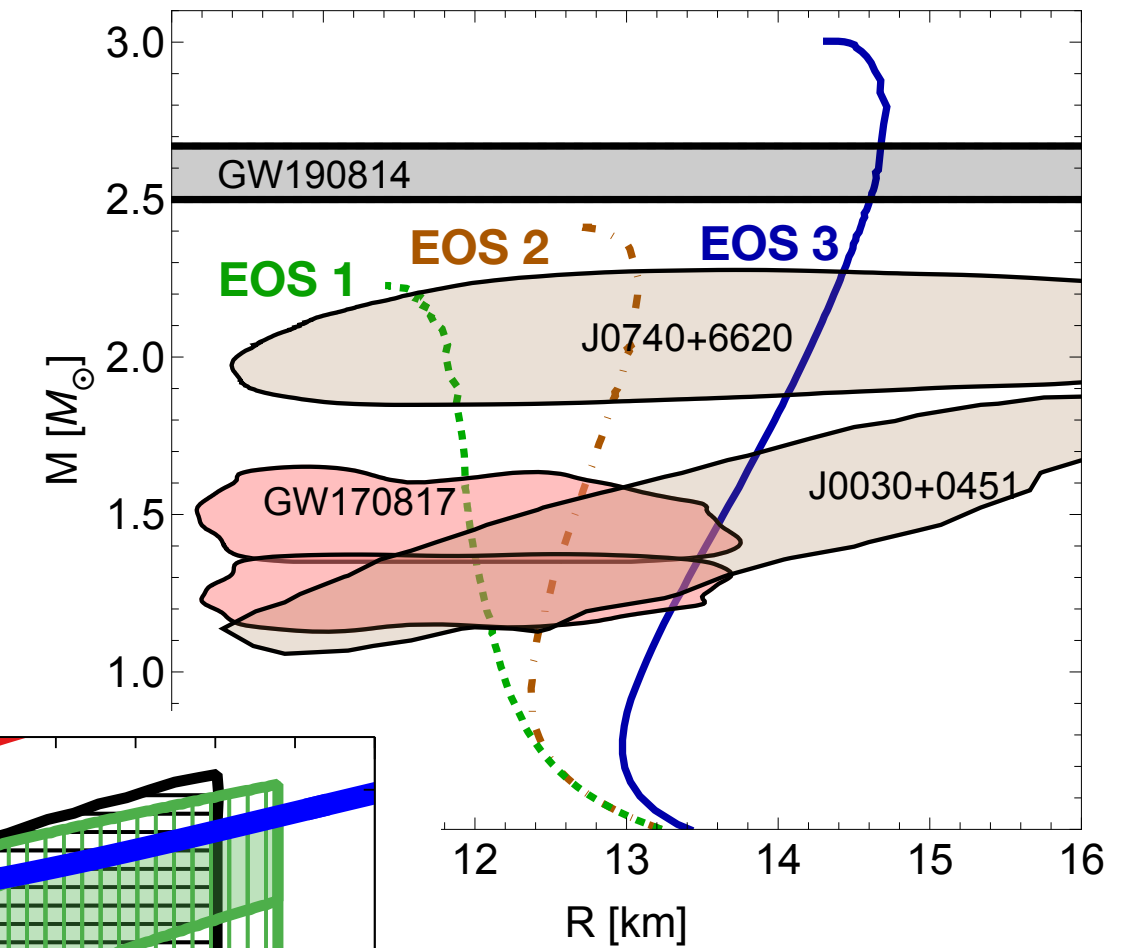
EOS from neutron stars to heavy-ion collisions

“Minimal” and “maximal” EOSs from each family tested against heavy-ion measurements:



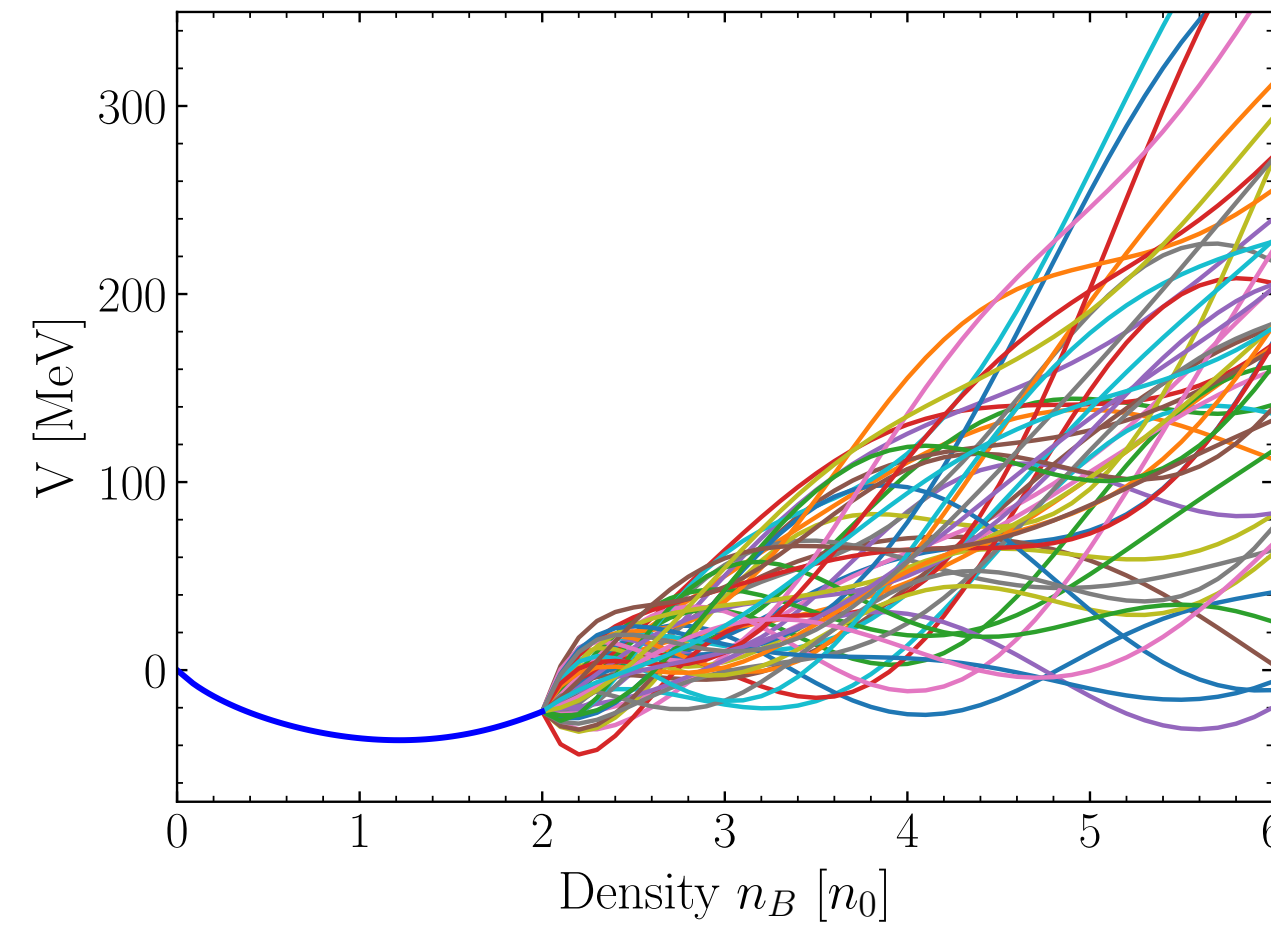
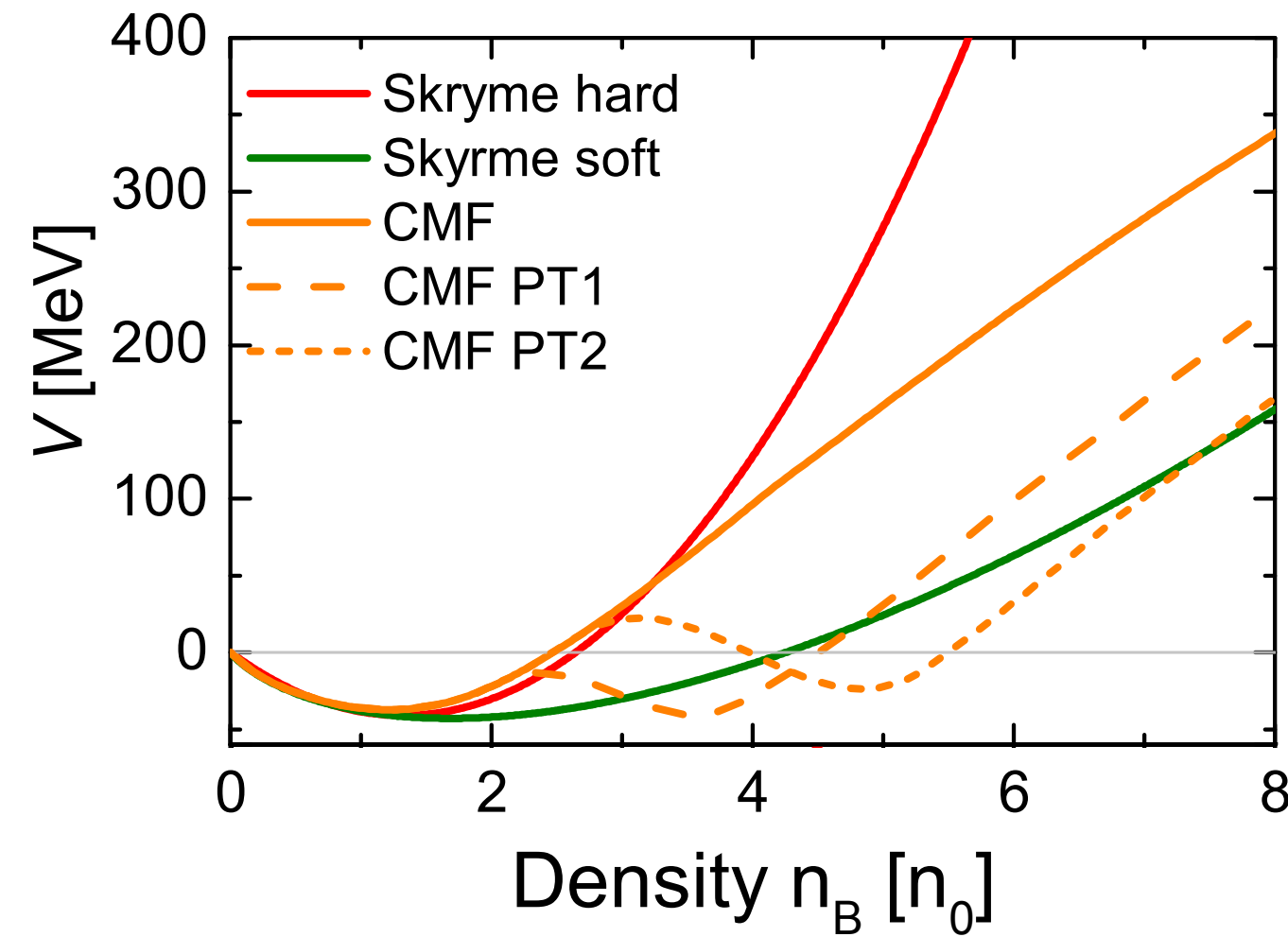
EOS	$n_{\text{sat}} [\text{fm}^{-3}]$	$B [\text{MeV}]$	$K_0 [\text{MeV}]$	$c_s^2(n_B = n_{\text{sat}})$
eos1				
eos2				
eos3				
eos3 max	0.162	-16.5	201.7	0.024

Massive neutron stars and heavy-ion collision data can be described with one EOS



N. Yao, A. Sorensen, V. Dexheimer, J. Noronha-Hostler, arXiv:2311.18819

Bayesian analysis of flow data in UrQMD



M. Omana Kuttan, J. Steinheimer, K. Zhou, H. Stoecker,
 Phys. Rev. Lett. **131** 20, 202303 (2023)
 arXiv:2211.11670

$$V(n_B) = \begin{cases} V_{\text{CMF}} & n_B \leq 2n_0 \\ \sum_{i=1}^7 \theta_i \left(\frac{n_B}{n_0} - 1\right)^i + C & n_B > 2n_0 \end{cases}$$

proton mean transverse kinetic energy $\langle m_T \rangle - m_0$:

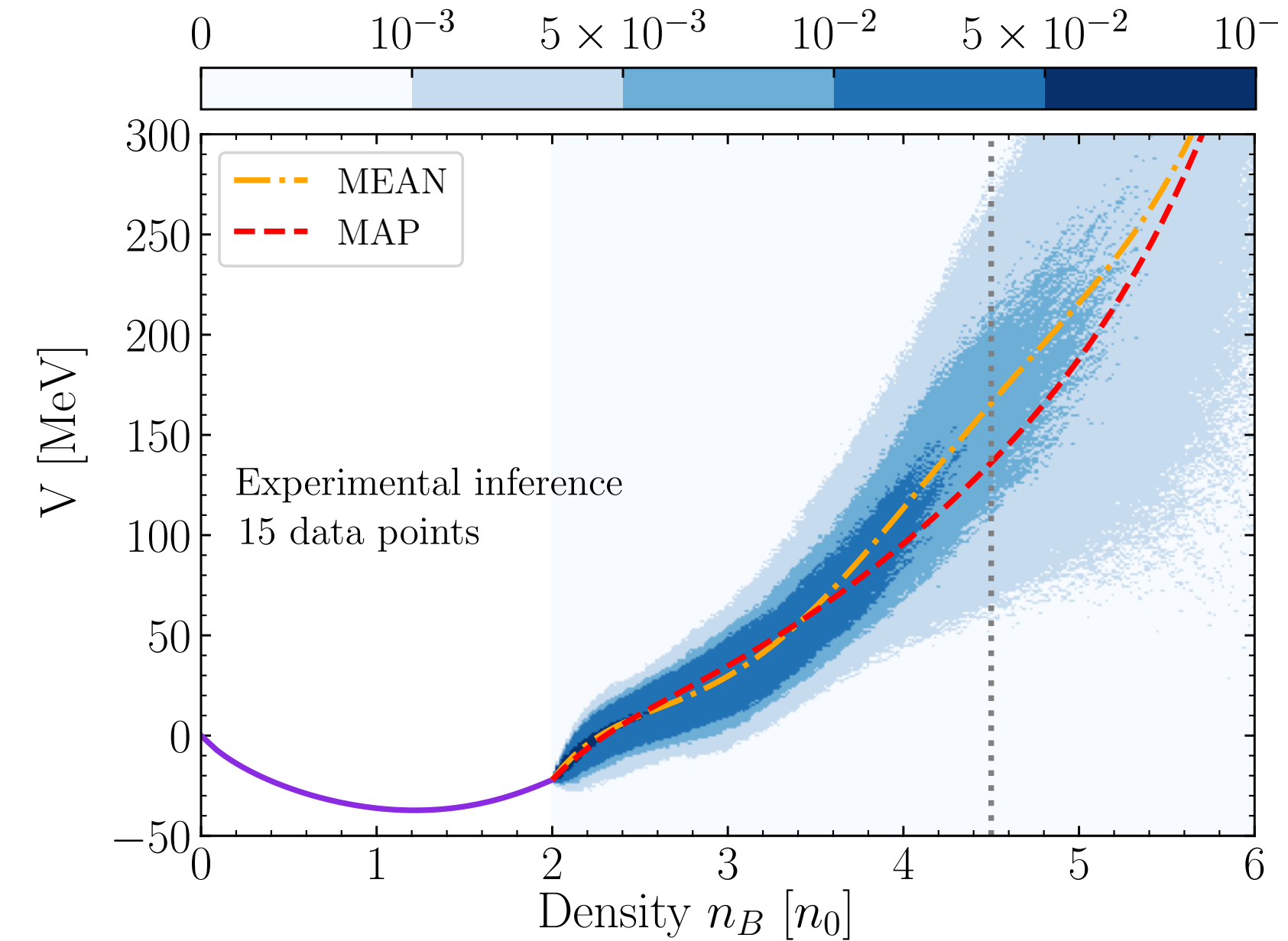
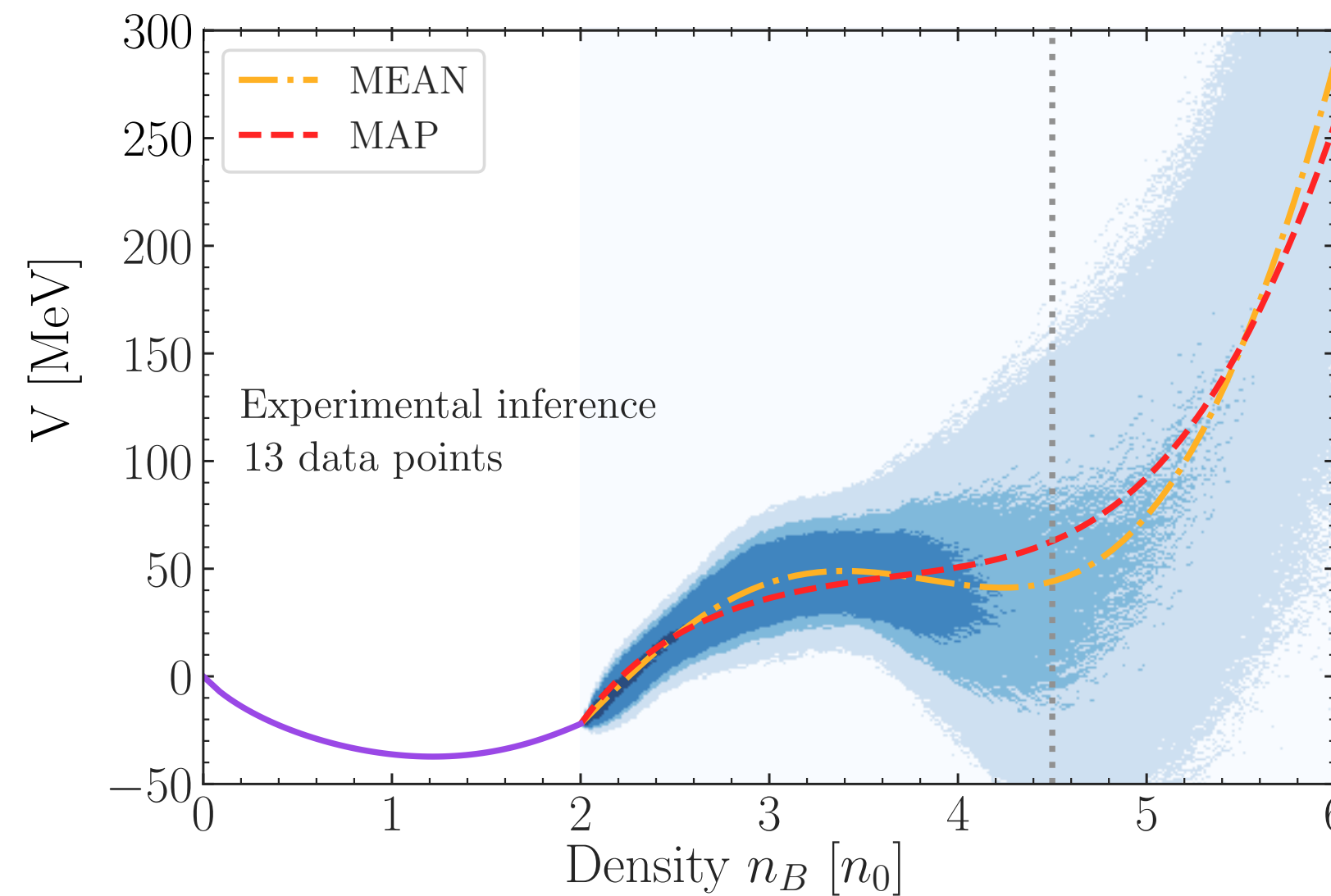
$$\sqrt{s_{\text{NN}}} \in [3.83, 8.86] \text{ GeV}$$

proton elliptic flow v_2 at midrapidity:

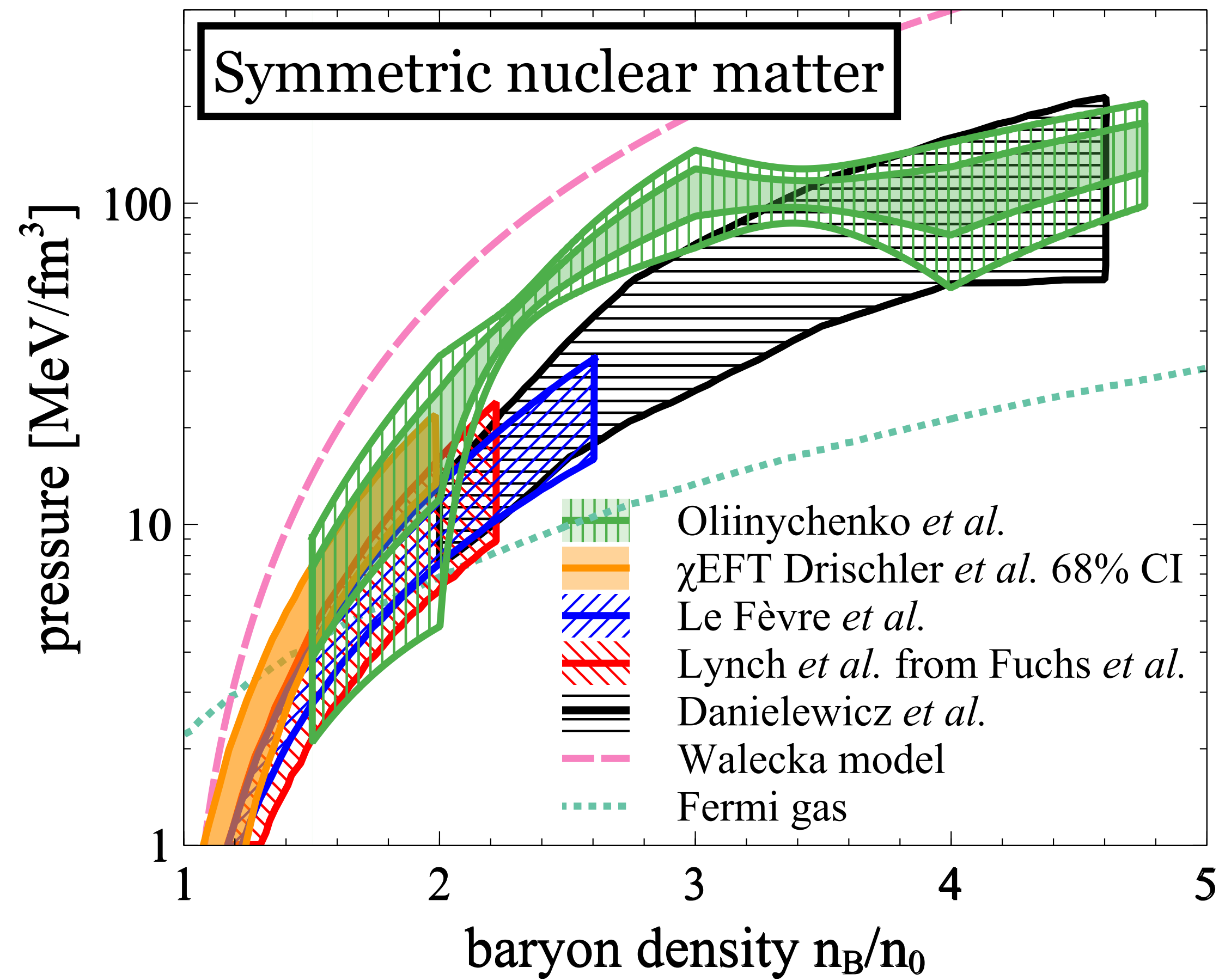
$$\sqrt{s_{\text{NN}}} \in [2.24, 4.72] \text{ GeV}$$

13 points = excluding $\langle m_T \rangle - m_0$ at the two lowest collision energies

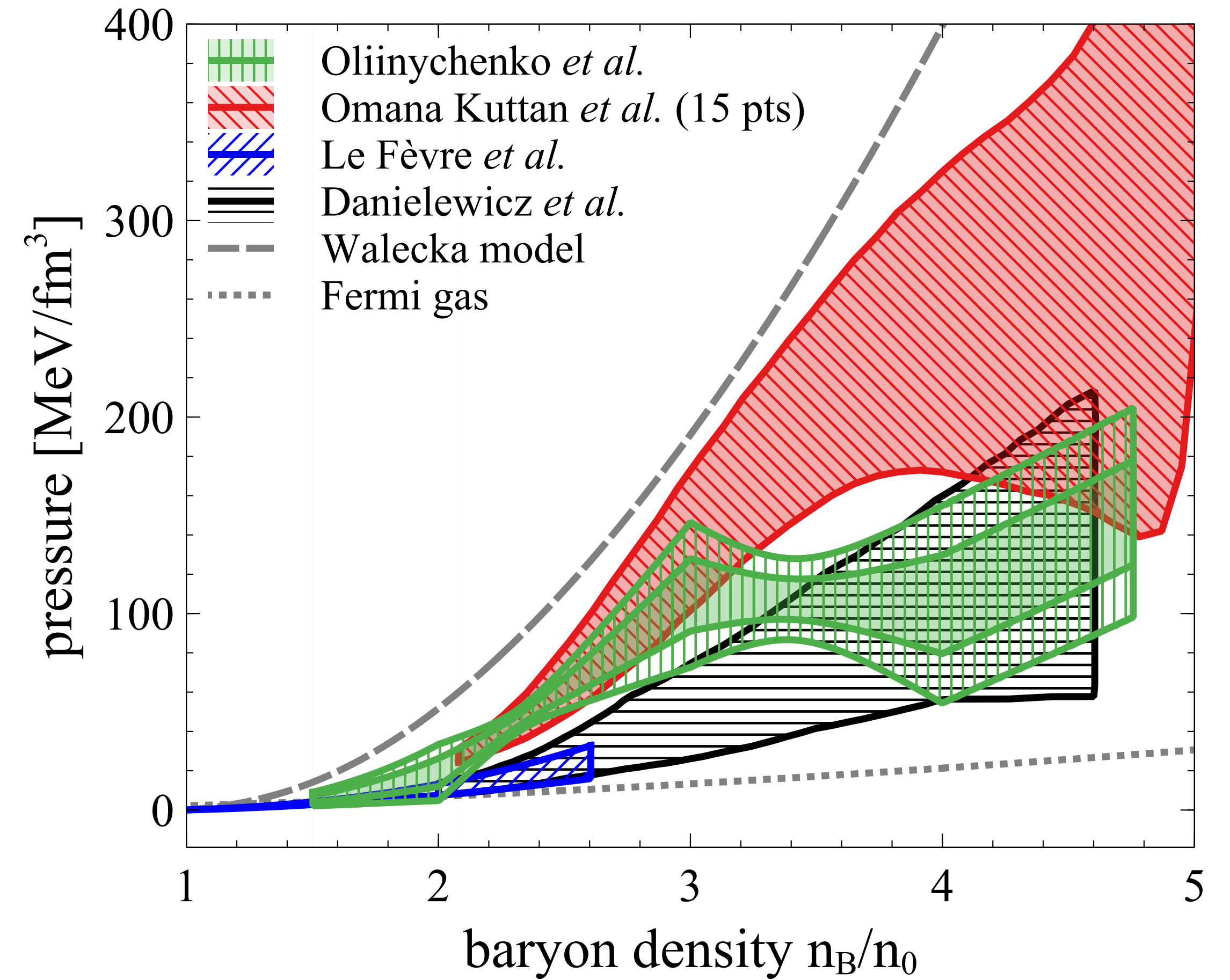
$$\sqrt{s_{\text{NN}}} = 3.83, 4.29 \text{ GeV}$$



EOS of symmetric nuclear matter: selected (*few*) results



A. Sorensen *et al.*, Prog. Part. Nucl. Phys. **134**, 104080 (2024)
arXiv:2301.13253



L. Du, A. Sorensen, M. Stephanov, Int. J. Mod. Phys. E
(available online), arXiv: 2402.10183

Bayesian analysis of heavy-ion collision and astronomical data

Determination of the equation of state from nuclear experiments and neutron star observations

Chun Yuen Tsang, ManYee Betty Tsang , William G. Lynch, Rohit Kumar & Charles J. Horowitz

Symmetric matter

Constraints

HIC(DLL)

HIC(FOPI)

GMR

Asymmetric matter

Constraints

Nuclear structure

α_D

PREX-II

Nuclear masses

Mass(Skyrme)

Mass(DFT)

IAS

Heavy-ion collisions

HIC(Isodiff)

HIC(n/p ratio)

HIC(π)

HIC(n/p flow)

Astronomical

Constraints

LIGO

*Riley PSR J0030+0451

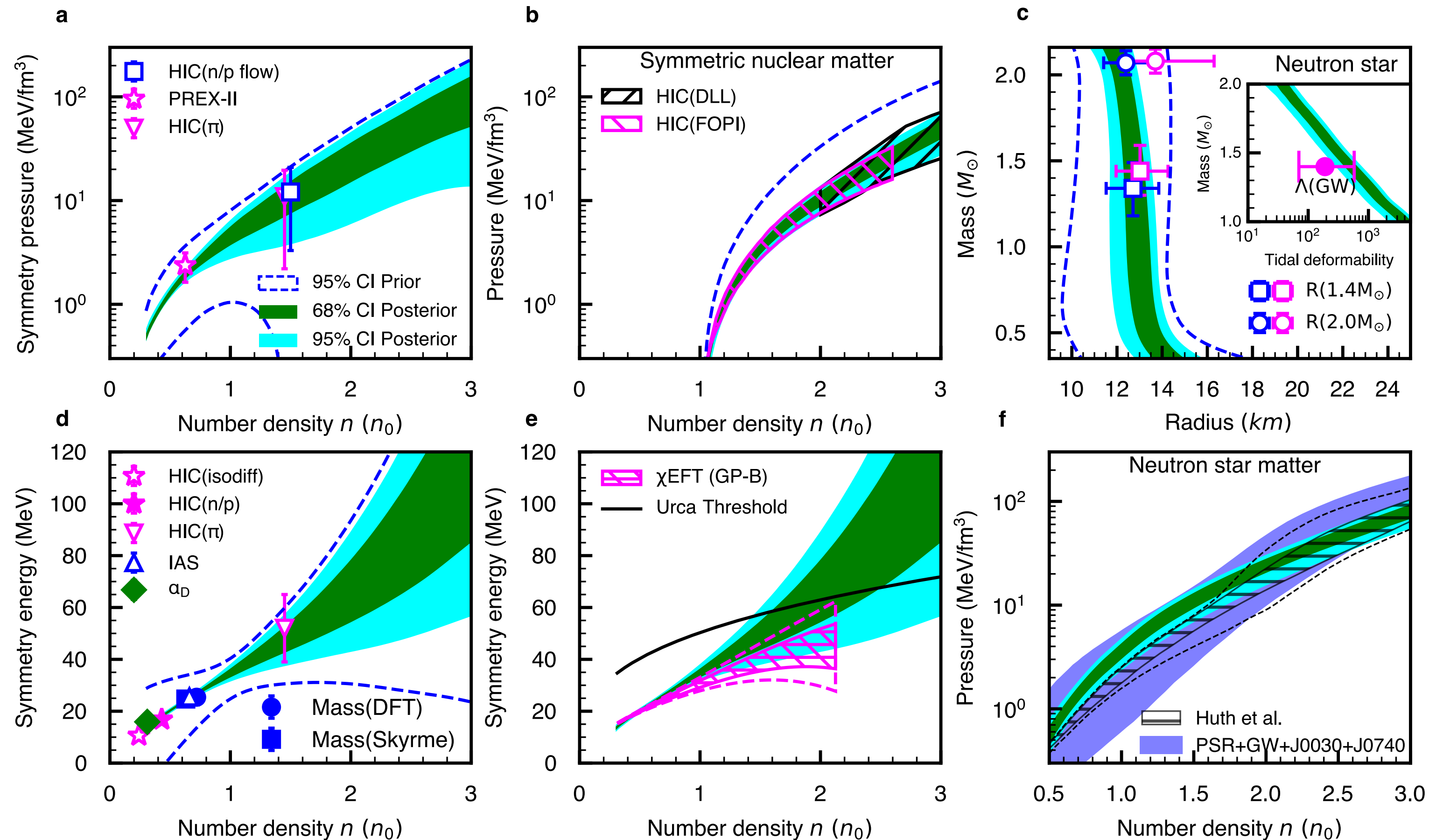
*Miller PSR J0030+0451

*Riley PSR J0740+6620

*Miller PSR J0740+6620

see talk by B. Tsang
Tue 12:00 pm plenary

C.-Y. Tsang, M. B. Tsang, W. G. Lynch, R. Kumar, C. J. Horowitz,
Nature Astron. **8** 3, 328-336 (2024) arXiv:2310.11588



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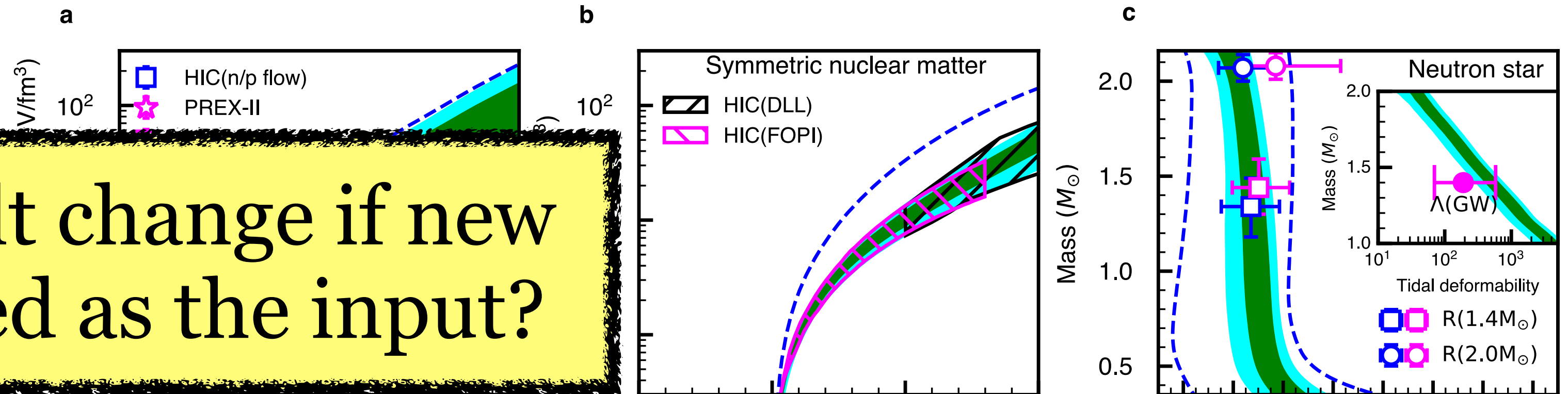
LIGO

*Riley PSR J0030+0451

*Miller PSR J0030+0451

*Riley PSR J0740+6620

*Miller PSR J0740+6620



How would the result change if new constraints were used as the input?

DLL is still the state-of-the-art result...

- momentum-dependence
- in-medium cross-sections
- also:
 - cluster production
 - meson potentials
 - initial state (e.g., short-range correlations)
 - ...

... for now :)

Summary

- Heavy-ion collisions at low energies probe multiple **fundamental properties of nuclear matter**
 - **density**, isospin, and momentum dependence of nuclear interactions
 - in-medium cross sections
 - cluster production mechanisms
 - ...
- Transport models needed for interpreting experiments (**STAR FXT, HADES, FRIB, CBM, FRIB400**)
- How to reconcile various effects affecting the extraction of the EOS?
(momentum dependence, in-medium cross sections, cluster production, ...)
Can we move away from phenomenology and toward guidance from theory?

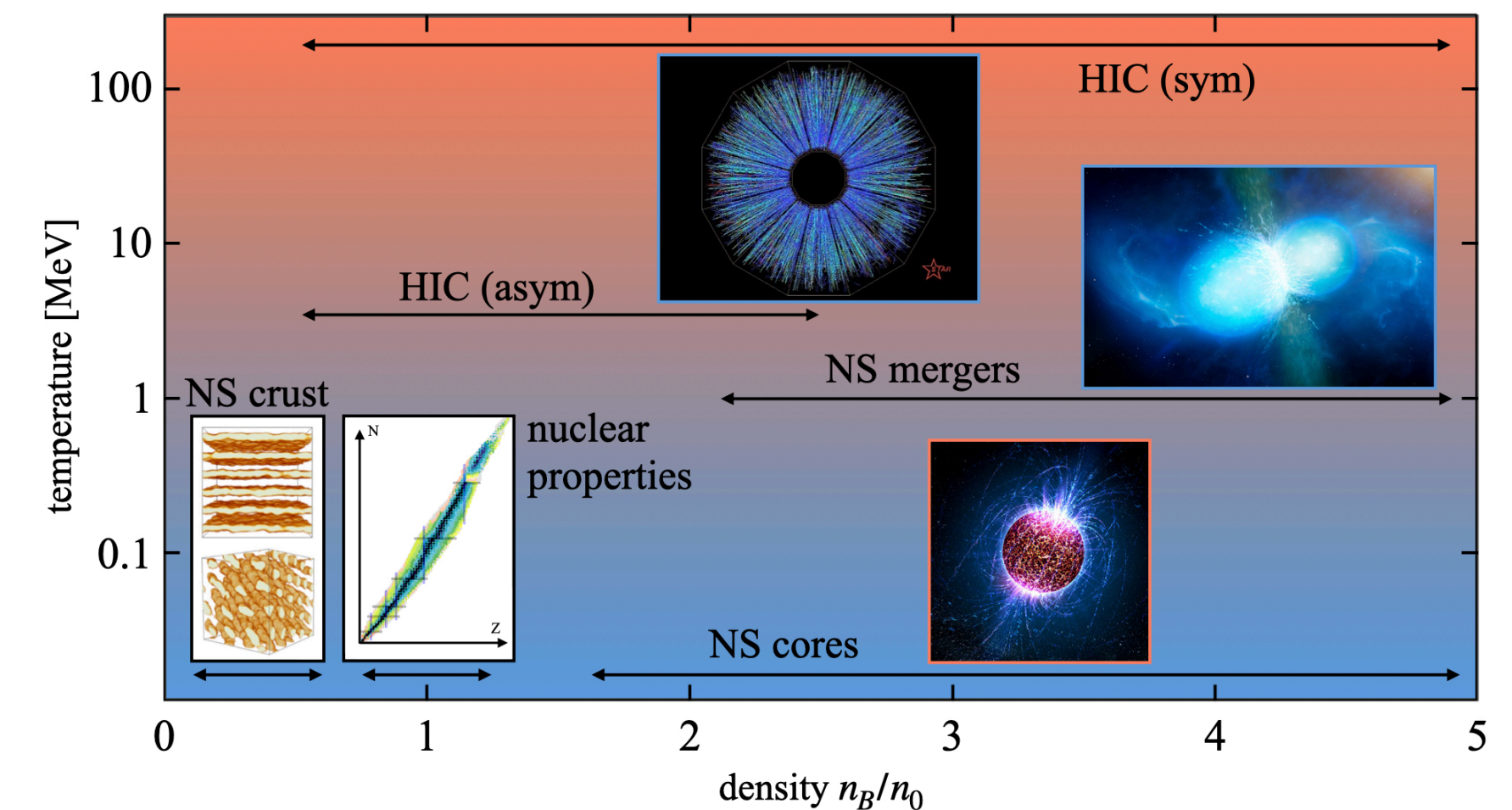
(needed *also* to ease the computational cost)

Some ideas in

A. Sorensen *et al.*,
Prog. Part. Nucl. Phys.
134, 104080 (2024)
arXiv:2301.13253

Dense Nuclear Matter Equation of State from Heavy-Ion Collisions *

Agnieszka Sorensen¹, Kshitij Agarwal², Kyle W. Brown^{3,4}, Zbigniew Chajecki⁵,
Paweł Danielewicz^{3,6}, Christian Drischler⁷, Stefano Gandolfi⁸, Jeremy W. Holt^{9,10},
Matthias Kaminski¹¹, Che-Ming Ko^{9,10}, Rohit Kumar³, Bao-An Li¹², William G. Lynch^{3,6},
Alan B. McIntosh¹⁰, William G. Newton¹², Scott Pratt^{3,6}, Oleh Savchuk^{3,13}, Maria Stefaniak¹⁴,
Ingo Tews⁸, ManYee Betty Tsang^{3,6}, Ramona Vogt^{15,16}, Hermann Wolter¹⁷, Hanna Zbroszczyk¹⁸



- Besides the extraction of the EOS, **how can low-energy heavy-ion collisions inform other sub-fields in nuclear physics?** (e.g., strangeness interactions important for physics of neutron stars)

Thank you for your attention!

